

Hands-on session on the R-matrix method:

Excel program talent_rmatrix_calc.xlsm

1. R-matrix formulae
2. Calculable R-matrix : general definitions
3. Writeup of the program rmatrix_calc.xlsm
4. Questions about the calculable R-matrix

1. R-matrix formulae

Note: all quantities depend on the angular momentum ℓ (not written)

$$\textbf{Collision (or scattering) matrix} : U(E) = \frac{I(ka) 1-L^*R(E)}{O(ka) 1-LR(E)} = \exp(2i\delta) = \exp(2i\delta_{HS}) \exp(2i\delta_R)$$

$$\textbf{Hard-sphere phase shift} : \delta_{HS} = -\text{atan} \frac{F(ka)}{G(ka)}$$

$$\textbf{R-matrix phase shift} : \delta_R = \text{atan} \frac{P(E)R(E)}{1-S(E)R(E)}$$

Incoming, outgoing Coulomb functions $I(x) = F(x) - G(x)$, $O(x) = F(x) + iG(x)$, where $F(x)$ and $G(x)$ are the Coulomb functions

$$\textbf{Penetration factor, shift factor} \text{ defined from } L(E) = ka \frac{O'(ka)}{O(ka)} = S(E) + iP(E)$$

Resonance energy E_r

In the R-matrix theory, the resonance energy E_r is a solution of $1 - S(E_r)R(E_r) = 0$

In the literature, the resonance energy E_r is also defined in different ways:

- 1) $\delta_r(E_r) = 90^\circ$ (corresponds to the definition above)
- 2) $\delta(E_r) = 90^\circ$
- 3) $d \sin \delta / dE |_{E=E_r} = 0$

For narrow resonances, these definitions provide similar results.

2. Calculable R-matrix : general definitions

The program solves the radial Schrödinger equation (positive or negative energies)

$$-\frac{\hbar^2}{2\mu} \left(\frac{d^2}{dr^2} - \frac{\ell(\ell+1)}{r^2} \right) u_\ell + (V_C(r) + V_N(r))u_\ell = E u_\ell$$

with the R matrix method. The coulomb $V_C(r)$ and nuclear $V_N(r)$ potentials are defined as (units are fm and MeV):

$$V_C(r) = \frac{Z_1 Z_2 e^2}{r}$$

$$V_N(r) = V_0 \exp(-(r/r_0)^2)$$

The wave function $u_\ell(r)$ is expanded over a set of ng basis functions as

$$u_\ell(r) = \sum_{i=1}^{ng} c_{\ell,i} \phi_i(r),$$

where $c_{\ell,i}$ are linear coefficients. Here we use Gaussian functions

$$\phi_i(r) = r^\ell \exp(-(r/a_i)^2),$$

with parameters $a_i = x_0 \times a_0^{i-1}$. Typical values are: $x_0=0.6$ fm, $a_0=1.4$. The matrix elements of the Hamiltonian are computed analytically. Parameters (x_0, a_0) can be changed by the user (cells F7 and F8). The phase shifts are compared with the exact solutions.

The R-matrix is defined by

$$\begin{aligned} R(E) &= \frac{\hbar^2 a}{2\mu} \sum_{ij} \phi_i(a) (C^{-1})_{ij} \phi_j(a) \\ C_{ij}(E) &= \langle \phi_j | H - E + \mathcal{L}(0) | \phi_i \rangle_{int} \\ &= \int_0^a \phi_j(r) \left[-\frac{\hbar^2}{2\mu} \left(\frac{d^2}{dr^2} - \frac{\ell(\ell+1)}{r^2} \right) + V(r) + \mathcal{L}(0) - E \right] \phi_i(r) dr, \end{aligned}$$

where $V(r)$ is the total potential (nuclear + Coulomb) and $\mathcal{L}(0)$ is the Bloch operator.

3. Writeup of the program rmatrix_calc.xlsm

Inputs

- Masses (B1,D1), charges (B2,D2)
- Parameters of the potential V_0 (B5), r_0 (D5), ℓ (F5)
- Number of basis functions ng (B6)
- Channel radius rmax (B7)
- Number of energies ne (B8), initial energy E0 (B9), energy step EP (B10)

Outputs

- *Phase shifts*
column A: energy
column B: “exact” phase shift (computed with the Numerov algorithm)
column C: R-matrix phase shift
column D: R matrix
column E: exact derivative of the wave function $u_\ell(r)$ at $r=r_{\max}$
column F: left derivative of the R-matrix wave function $u_\ell(r)$ at $r=r_{\max}$
column G: right derivative of the R-matrix wave function $u_\ell(r)$ at $r=r_{\max}$
column H: $1/S_\ell(E)$, where $S_\ell(E)$ =shift function
column I: hard-sphere phase shift
column J: penetration factor $P_\ell(E)$
- *Wave functions*
Computed at the energy given in cell M11, and displayed in columns:
column L: exact wave function
column M: internal R-matrix wave function
column N: external R-matrix wave function
- *Basis wave functions*
Displayed in sheet “Wave func.”
- *Potential*
Displayed in sheet “Potential”

4. Questions about the calculable R-matrix

We consider the $^{12}\text{C}+p$ system (see Fig.1)

$\ell = 0$: $V_0 = -73.8$, $r_0 = 2.7$ (reproduce the $^{12}\text{C}+p$ resonance ($1/2^+$) energy 0.42 MeV)

$\ell = 1$: $V_0 = -55.3$ (reproduce the ^{13}N ground state ($1/2^-$) energy -1.94 MeV)

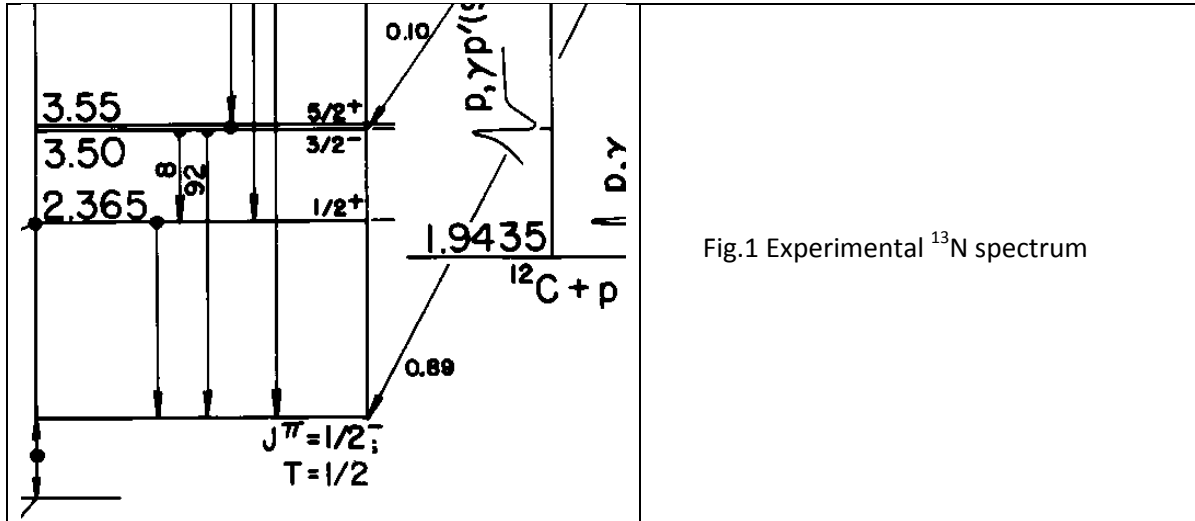


Fig.1 Experimental ^{13}N spectrum

- Take $x_0 = 0.6$ fm and $a_0 = 1.4$, and choose a realistic R-matrix radius r_{max} (tabulate the potential).
- Compute the phase shifts up to 2 MeV for different numbers of basis functions (n_g). Check out the left and right derivatives.
- Increase and decrease r_{max} . Adapt n_g accordingly.
- Choose an energy, and compute (and plot) the wave function in appropriate (n_g large enough) and poor (n_g too small) conditions.
- Verify the Thomas approximation (shift function $S(E)$ linear near the resonance energy)
- Determine approximately (by graphic) the resonance energy E_r by using the R-matrix definition above (plot $R(E)$ and $1/S(E)$)
- Choose an appropriate set (r_{max} , n_g) and compute the phase shifts with the pole expansion. Verify that the phase shifts are identical.
- With the pole energies E_λ and reduced widths γ_λ^2 , determine (plot) both terms of

$$R(E) = \frac{\gamma_0^2}{E_0 - E} + R_0(E)$$

and verify that the background term $R_0(E)$ is approximately constant.

Perform the same analysis for another R-matrix radius, and verify that both terms are sensitive to r_{max} (although the phase shift is not).

- Compute the phase shifts with the first term only (use $P(E)$ and $S(E)$ given in columns H and J). Evaluate the importance of the background term.