

ISOMERS Pb- Po- Rn- Fr...

- 1] Technicalities: lifetimes/widths/strengths..
 - 2] Building blocks...
 - 3] Signatures of structure: seniority isomers, K- isomers
 - 4] Triple shape coexistence - ^{188}Pb example
-

Four Case Studies

Case I: ^{209}Fr - maybe

Case II: ^{212}Rn - E3 transitions; signatures

- Empirical Shell Model View
- Deformed Independent Particle Model View

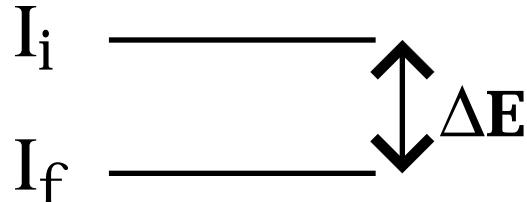
Case III: $^{190,192,194,194}\text{Pb}$ - 11^- isomers; E3 strengths

- oblate deformation

Case IV: ^{189}Pb - $\nu i_{13/2}^- 3^-$ isomer ?

- residual interactions/ non- maximal coupling

τ ... and the meaning of life...

$$\lambda \geq I_i - I_f$$

$$\tau \sim 1 / \{ | \langle f | T_\lambda | i \rangle |^2 \times \Delta E^{(2\lambda+1)} \}$$

**von Weizsacker 1936

structure dependent

partly structure dependent
partly random

difficult to calculate for
forbidden transitions

difficult to predict ΔE
accurately

low energy
small matrix element
high multipolarity...

nowhere to go...

Play School

τ (not $T_{1/2}$) - the "natural" unit

$$\Gamma \cdot \tau = 0.658 \times 10^{-15} \text{ eV.s}$$

$$\begin{aligned}\Gamma &= \Gamma_{\gamma_1} + \Gamma_{e_1} + \Gamma_{\gamma_2} + \Gamma_{e_2} + \dots \\ &= (1 + \alpha_1^T) \Gamma_{\gamma_1} + (1 + \alpha_2^T) \Gamma_{\gamma_2} + \dots\end{aligned}$$

Weisskopf unit (eV)

$$\Gamma_W(E1) = 6.748 \times 10^{-2} A^{2/3} E_\gamma^3$$

$$\Gamma_W(E2) = 4.792 \times 10^{-8} A^{4/3} E_\gamma^5$$

$$\Gamma_W(M1) = 2.072 \times 10^{-2} E_\gamma^3$$

$$\Gamma_W(M2) = 1.472 \times 10^{-8} A^{2/3} E_\gamma^5$$

$$\Gamma_W(E3) = 2.333 \times 10^{-14} A^2 E_\gamma^7$$

Play School

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typical lifetimes:(Pb 500 keV)

10^{-6} W.u. E1 23 ns

1 W.u. E2 0.4 ns

0.1 W.u. M1 0.002 ns

< 1 W.u. M2 > 32 ns

10 W.u. E3 8.6 μ s

a long lifetime does
not imply hindrance

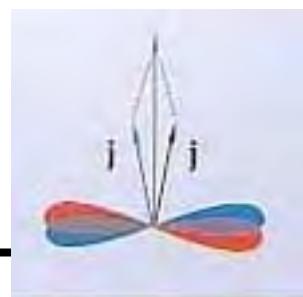
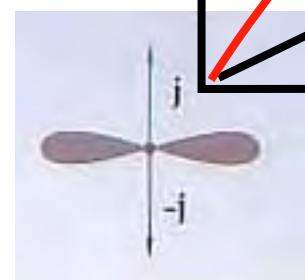
How

Yrast Line

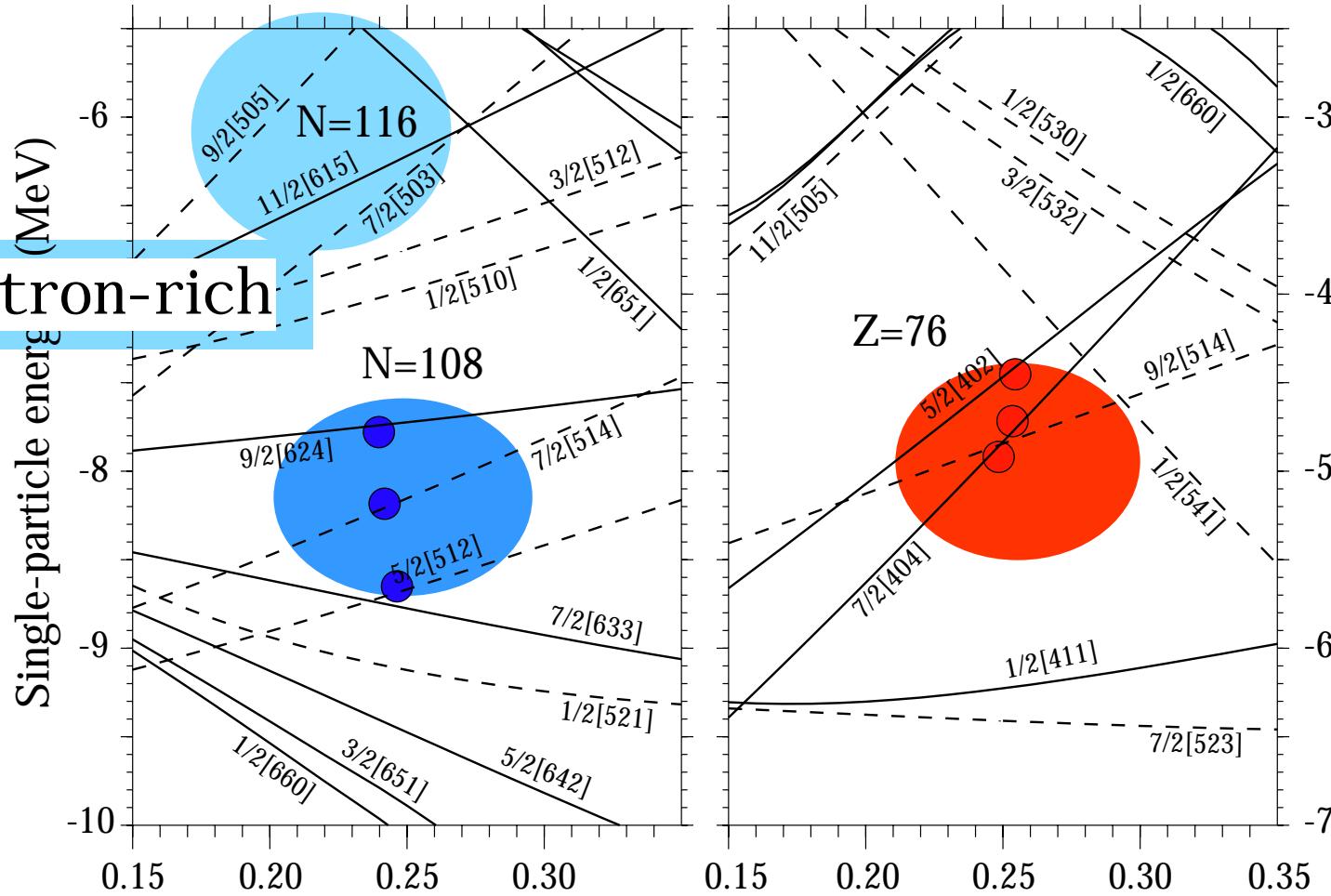
*Residual
spin-spin
 j or K*

*Non-uniform
levels -shell gaps
spherical or deformed*

*High- j or
High- K*



$Z \sim 76$; $N \sim 104$ favoured region for high-K



$$K = \nu |\Omega_1 + \Omega_2 + \Omega_3 \dots|^{S_2} + \pi |\Omega_1 + \Omega_2 + \Omega_3 \dots|$$

$$\sim 9/2 + 7/2 + 5/2 \dots + 7/2 + 9/2 + 5/2 \dots = 21$$

High-K Intrinsic states cost less energy than rotation

Z = 72
N ~ 104

common deformed building blocks

align high- Ω orbitals

$8^- \nu \quad 9/2^+[624] \times 7/2^- [514]$
 $8^- \pi \quad 9/2^- [514] \times 7/2^+ [404]$

16+

lower because of residual interaction →

^{178}Hf exp.

the 31-year isomer

$8^- 1147$
 $8^- 1479$

$16^+ 2626$ keV
cf: 2447 keV

Z = 72
N ~ 104

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Z ~ 82-
N ~ 124

common spherical building blocks

align high-j orbitals

^{212}Po exp.
the 65 sec. isomer

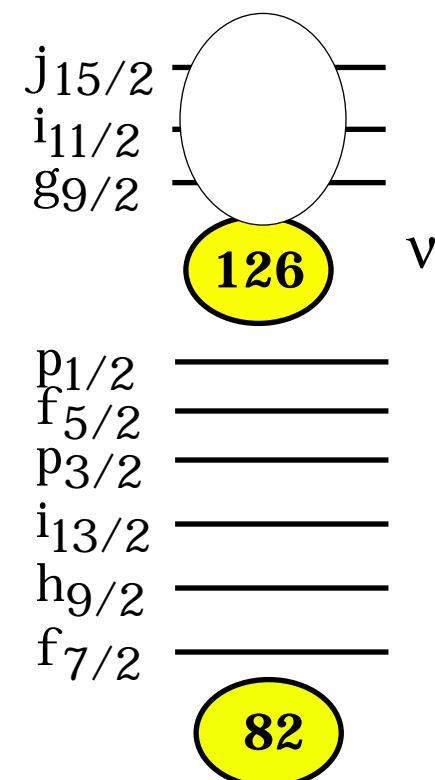
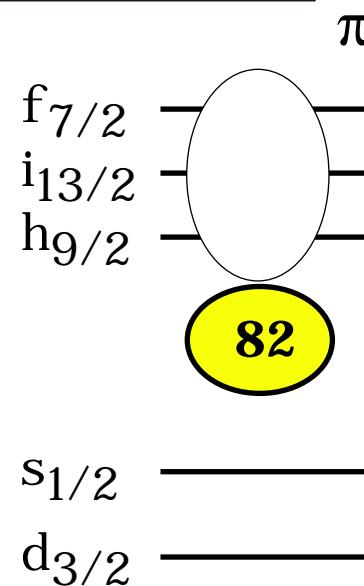
e.g.

$8^+ \pi h_{9/2}^2$
 $10^+ \nu g_{9/2} i_{11/2}$

18+

$8^+ 1476$
 $10^+ 1834$

 $18^+ 3310$ keV
cf: 2922 keV



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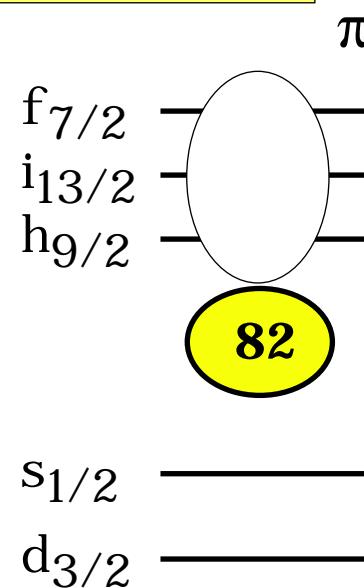
18+

$$8^+ 1476$$

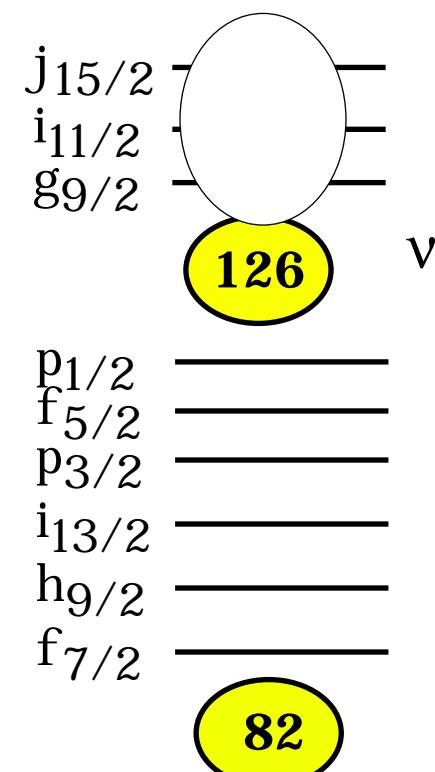
$$10^+ 1834$$

$$18^+ 3310 \text{ keV}$$

$$\text{cf: } 2922 \text{ keV}$$



11- $\pi h_{9/2} i_{13/2}$, 12+ $\nu i_{13/2}^{-2}$

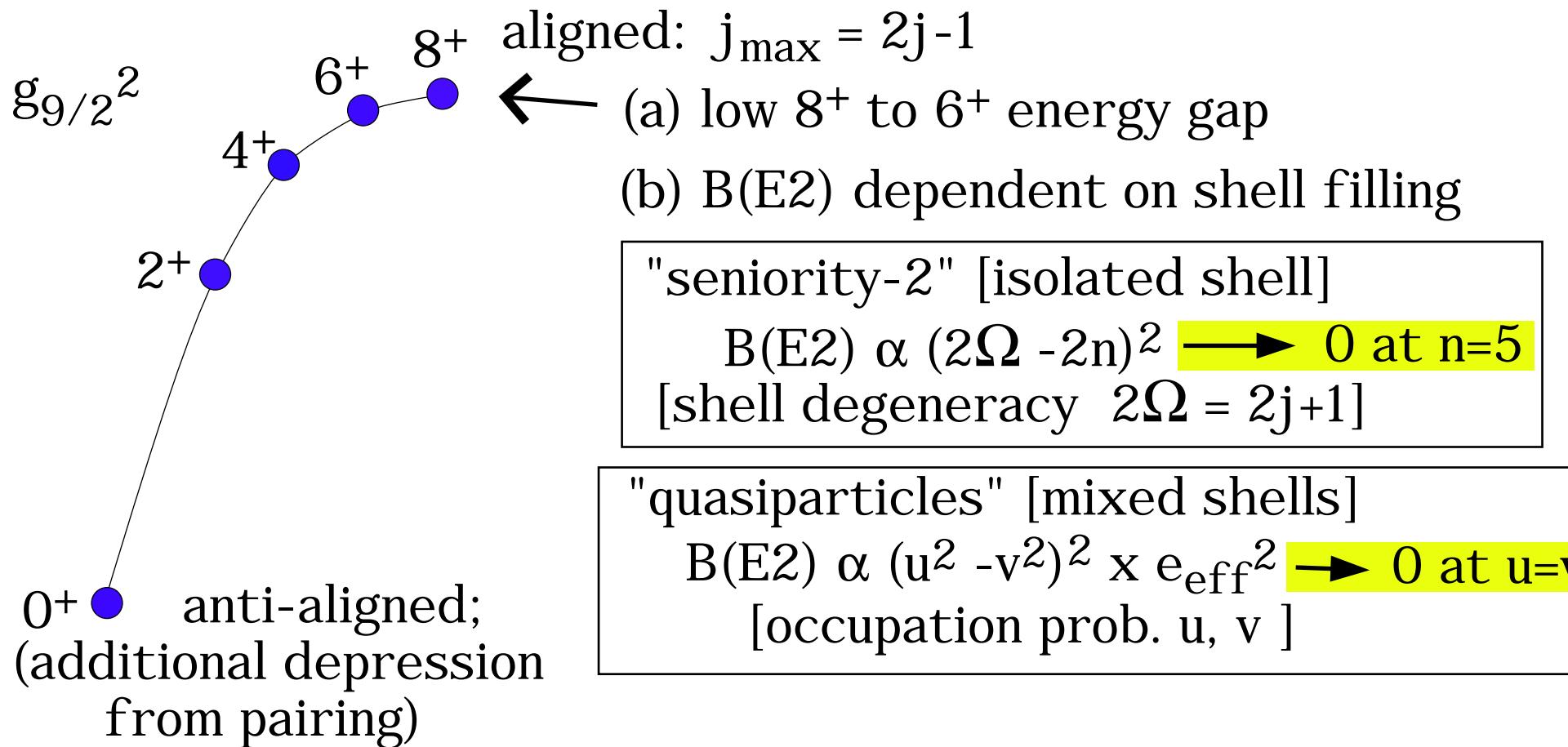


many factors...

Low transition energies - low orbitals + favourable interactions,
large configuration changes, K-hindrance, etc. etc.

e.g.

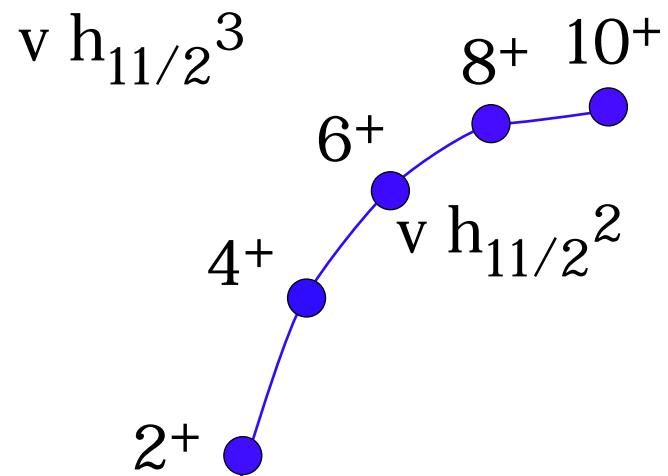
j^n multiplets - two (or more) identical particles in single j -shell



signature of shells

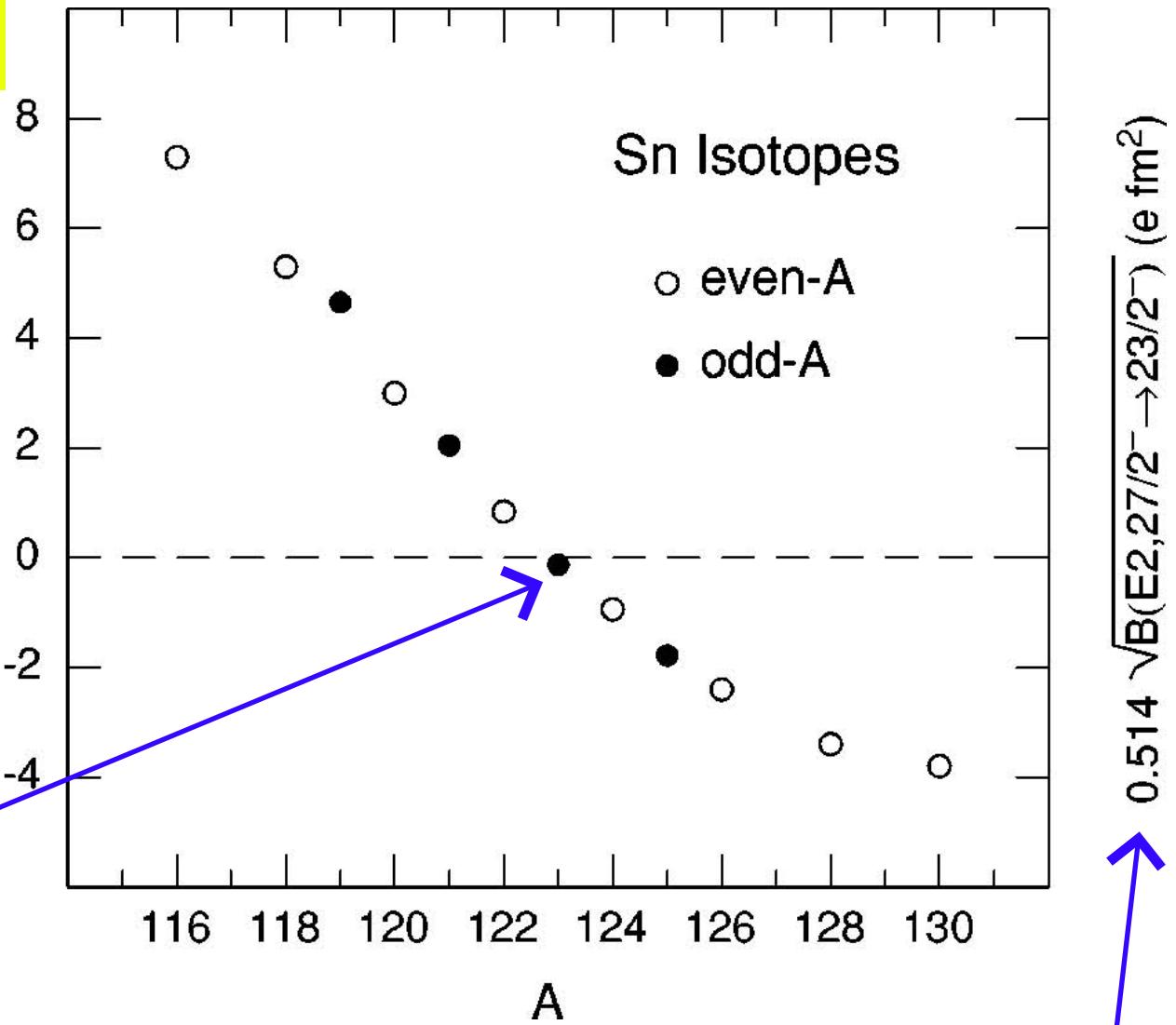
Zhang et al PRC 62 (2000) 057305: 10^+ and $27/2^-$ isomers

$$m = 11/2 + 9/2 + 7/2 = 27/2$$



negative sign assumed

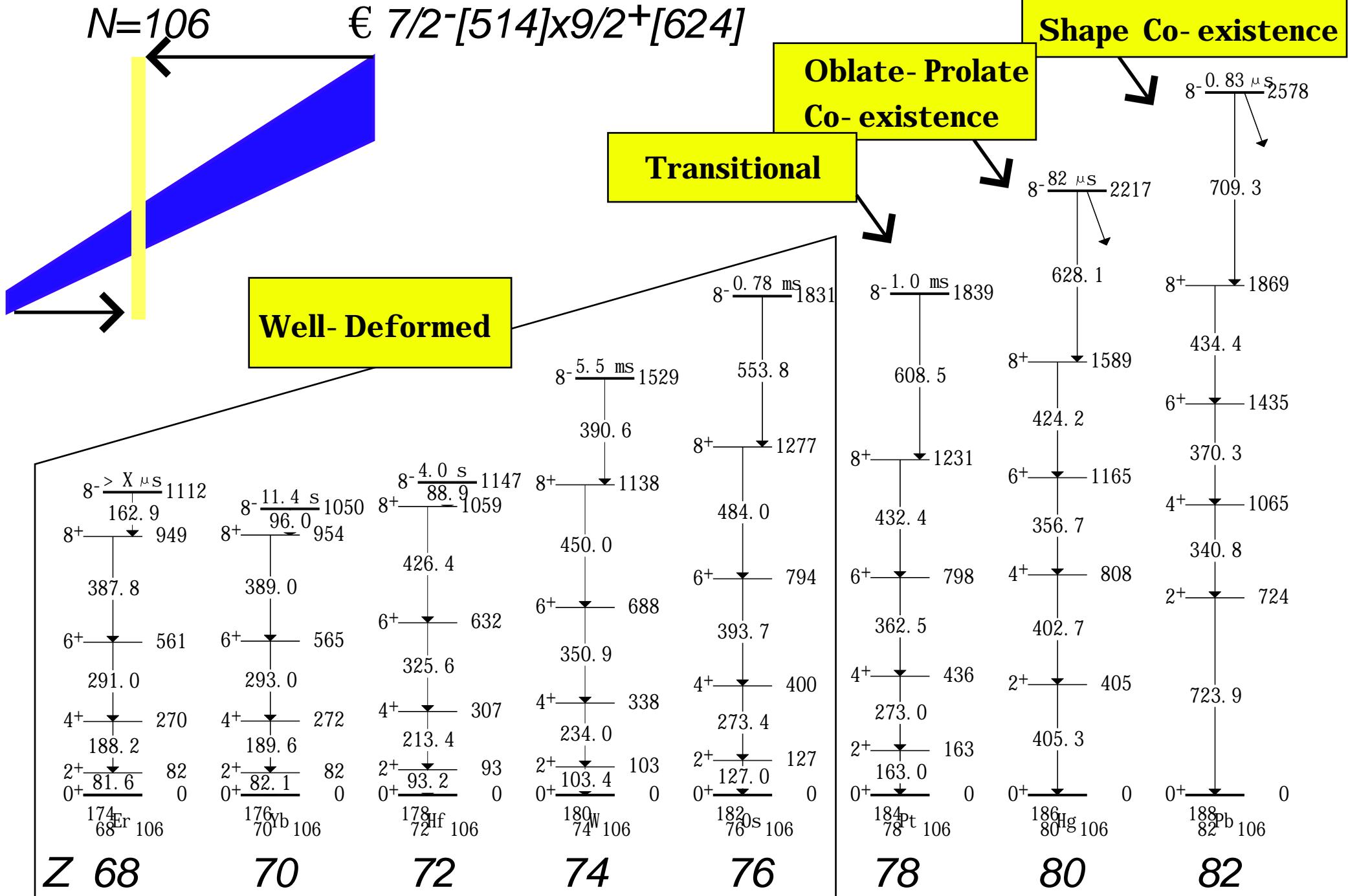
$$\sqrt{B(E2, 10^+ \rightarrow 8^+)} (\text{e fm}^2)$$



10^+ to 8^+ compared to $27/2^-$ to $23/2^-$
[adjusted for $v=3/v=2$ geometric factors]

K^π = 8⁻ Isomer; Signature of prolate Deformation

Identification of ^{174}Er & 8- isomer in prolate Well



Triple shape co-existence - Isomers

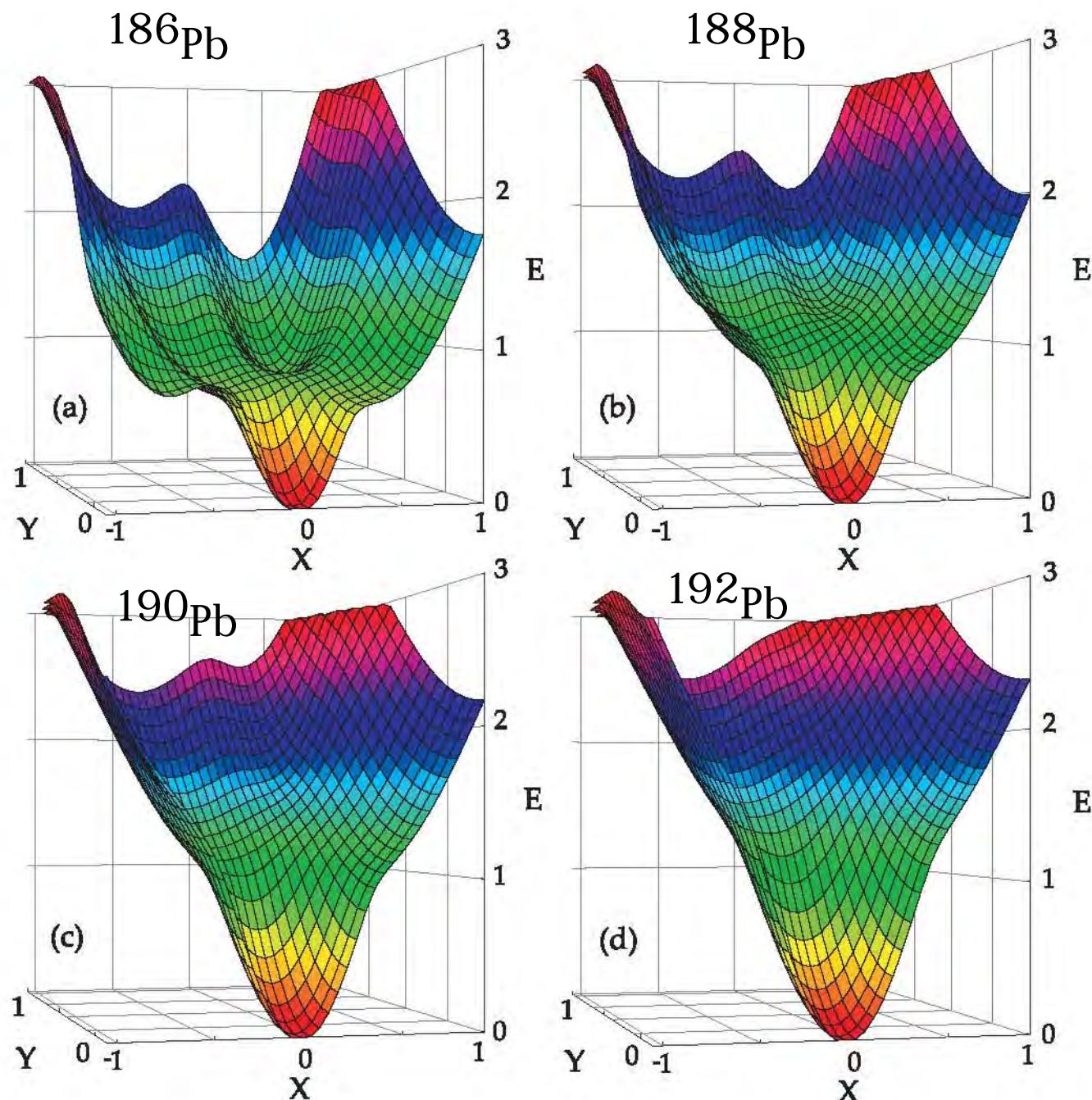
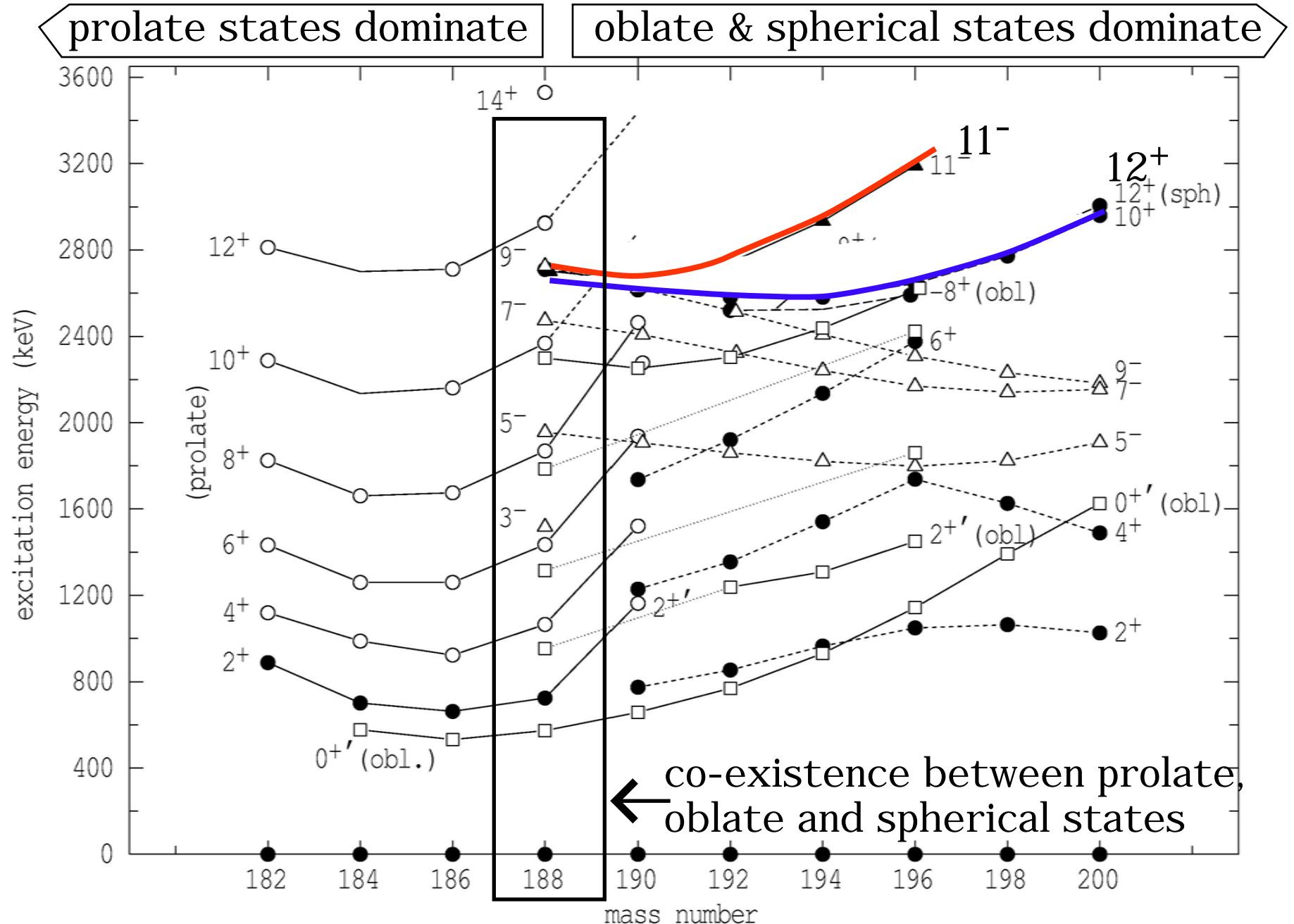


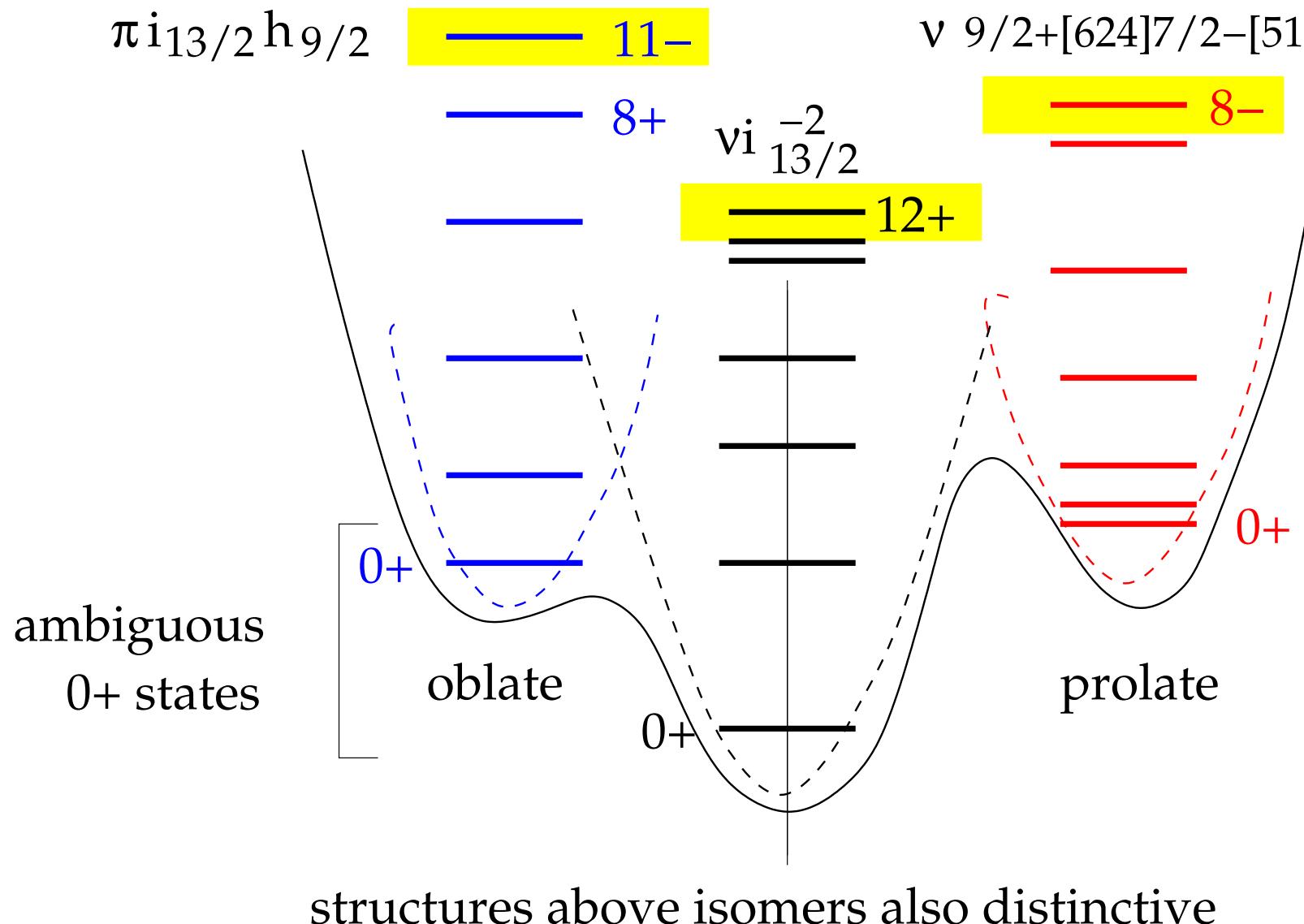
FIGURE 3.1: Energy surfaces for (a) ^{186}Pb , (b) ^{188}Pb , (c) ^{190}Pb , and (d) ^{192}Pb . On the horizontal axes, $X \equiv \beta \sin(\gamma + \pi/6)$, while $Y \equiv \beta \cos(\gamma + \pi/6)$. The vertical axis gives the energy in MeV. The figure is taken from [Fra04] and the parameters used to construct the surfaces were determined in [Fos03].

Pb Systematics



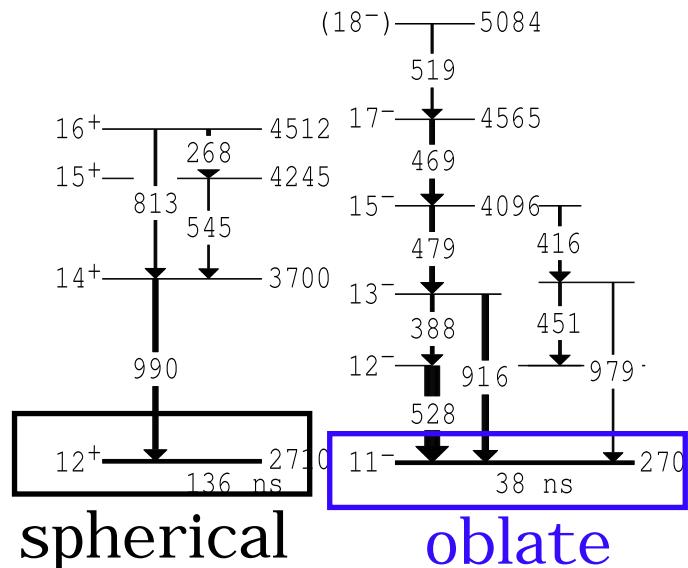
predicted Z=82 Triple Well

characteristic isomers in each: structure, spin and parity

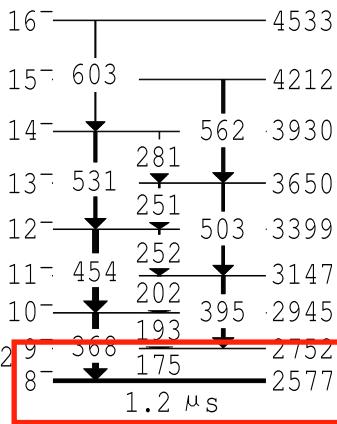


^{188}Pb spherical/oblate/prolate

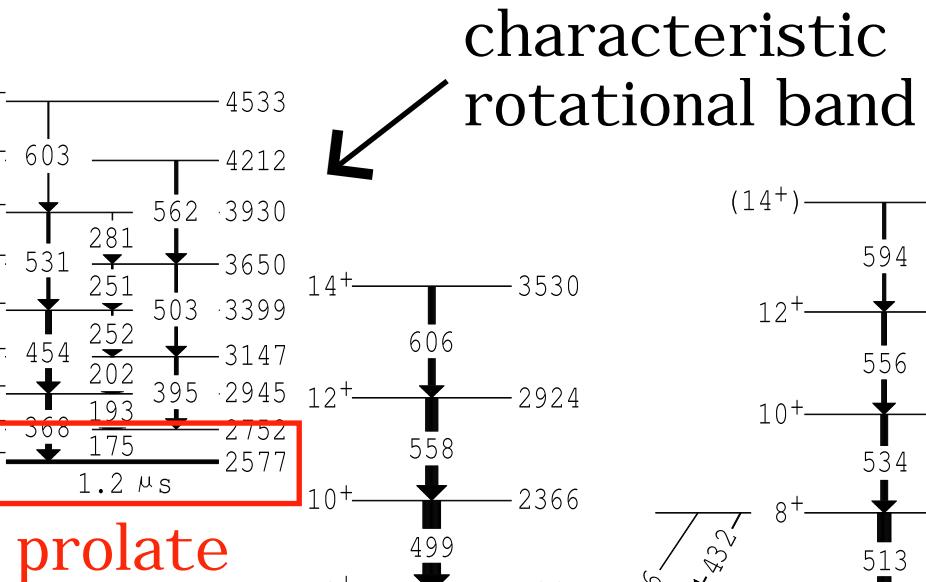
GDD et al. , PRC 67, 05130(R) (2003); PRC 69, 054318 (2004)



spherical

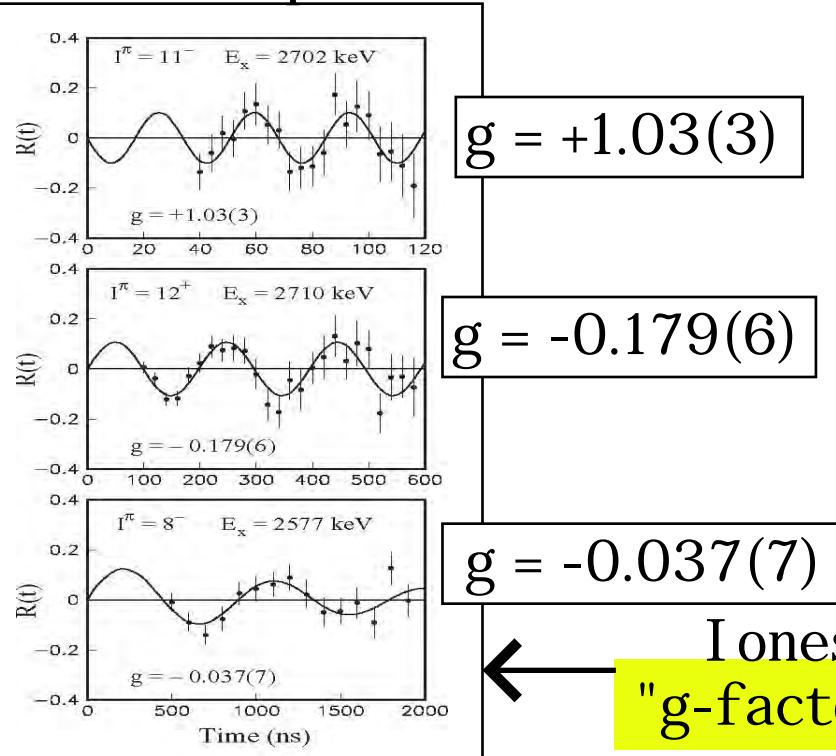


oblate



prolate

characteristic
rotational band



Ionescu-Bujor et al. PRC 81, 024323 (2010)
"g-factors of co-existing isomeric states in ^{188}Pb "

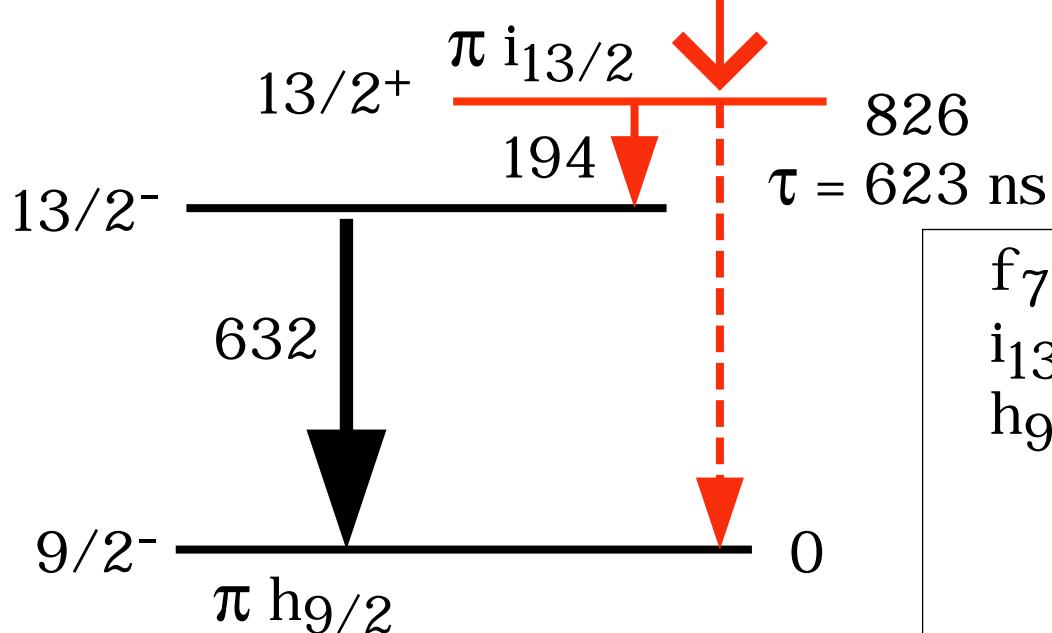
CASE - I: ^{209}Fr
[5 valence protons; 4 neutron holes]

Case I: ^{209}Fr Z=87, N=122 , $i_{13/2}$ proton state?

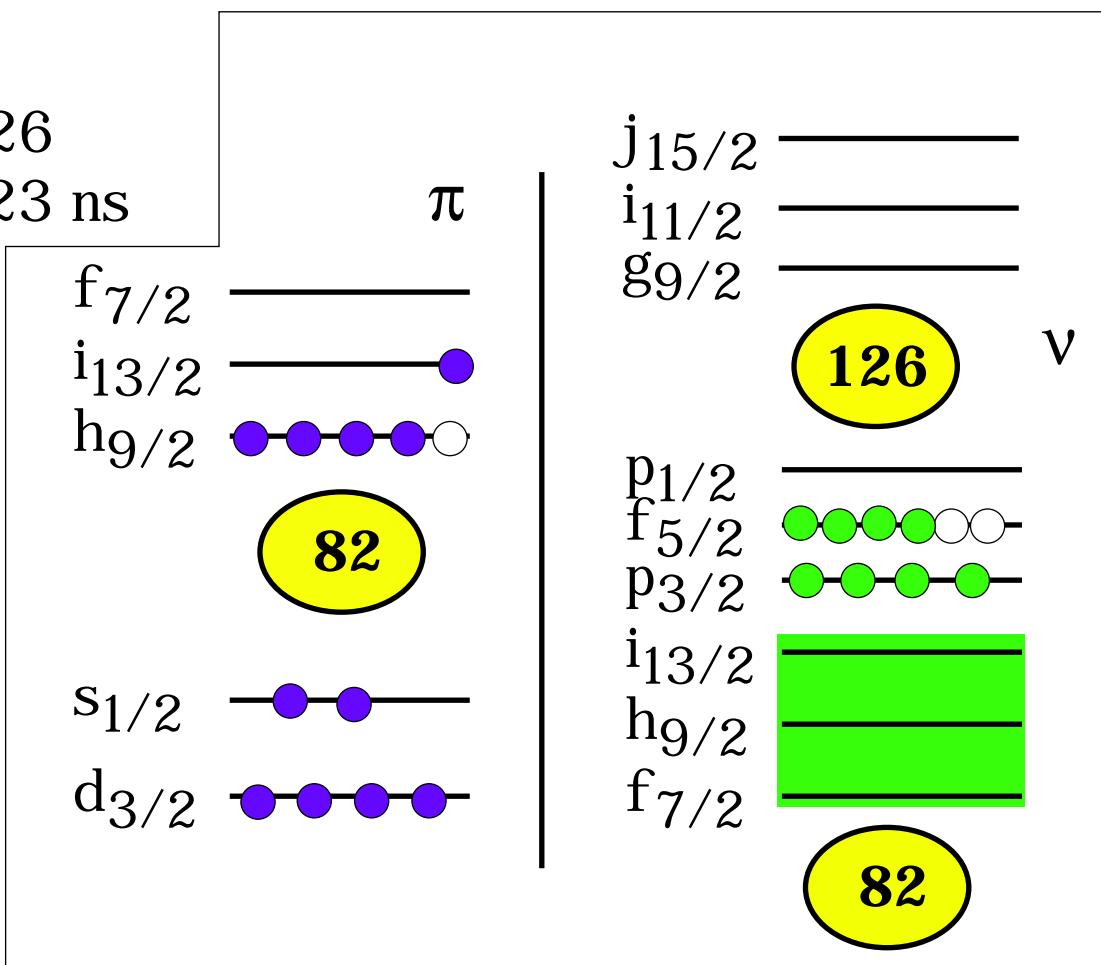
$$\Gamma(E1,194) = 1 \times 10^{-9} \text{ eV}$$

$\pi i_{13/2}$ to $h_{9/2}$ M2 ~ 0.2 W.u. ←
 $\Gamma(M2,826) \sim 40 \times 10^{-9} \text{ eV}; \tau_\gamma = 17 \text{ ns}$

M2 BRANCH SHOULD DOMINATE



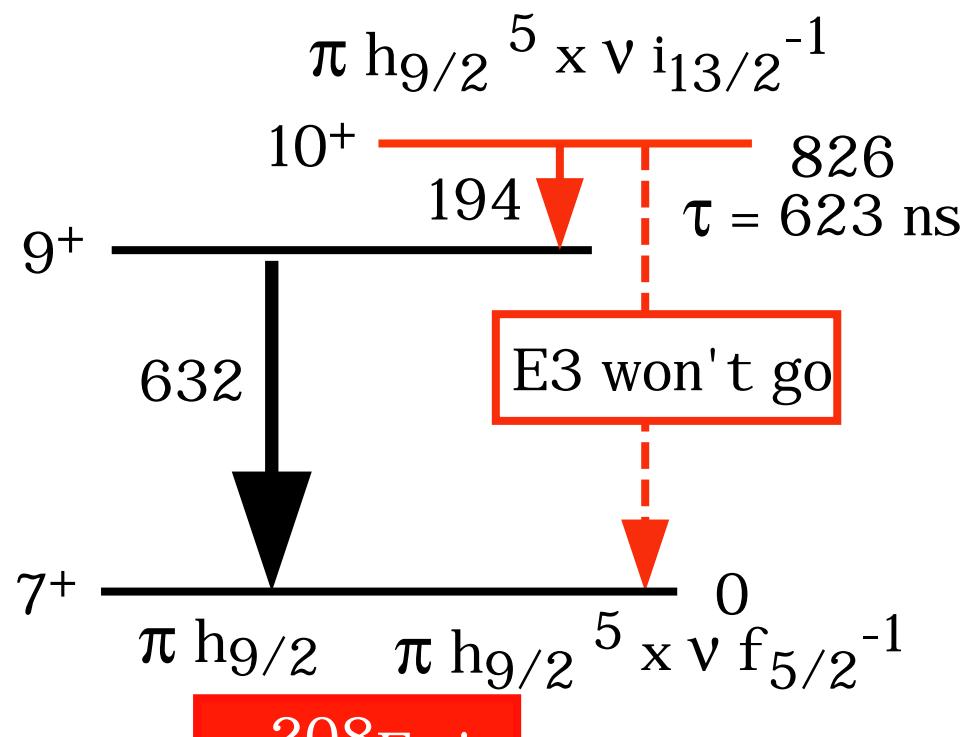
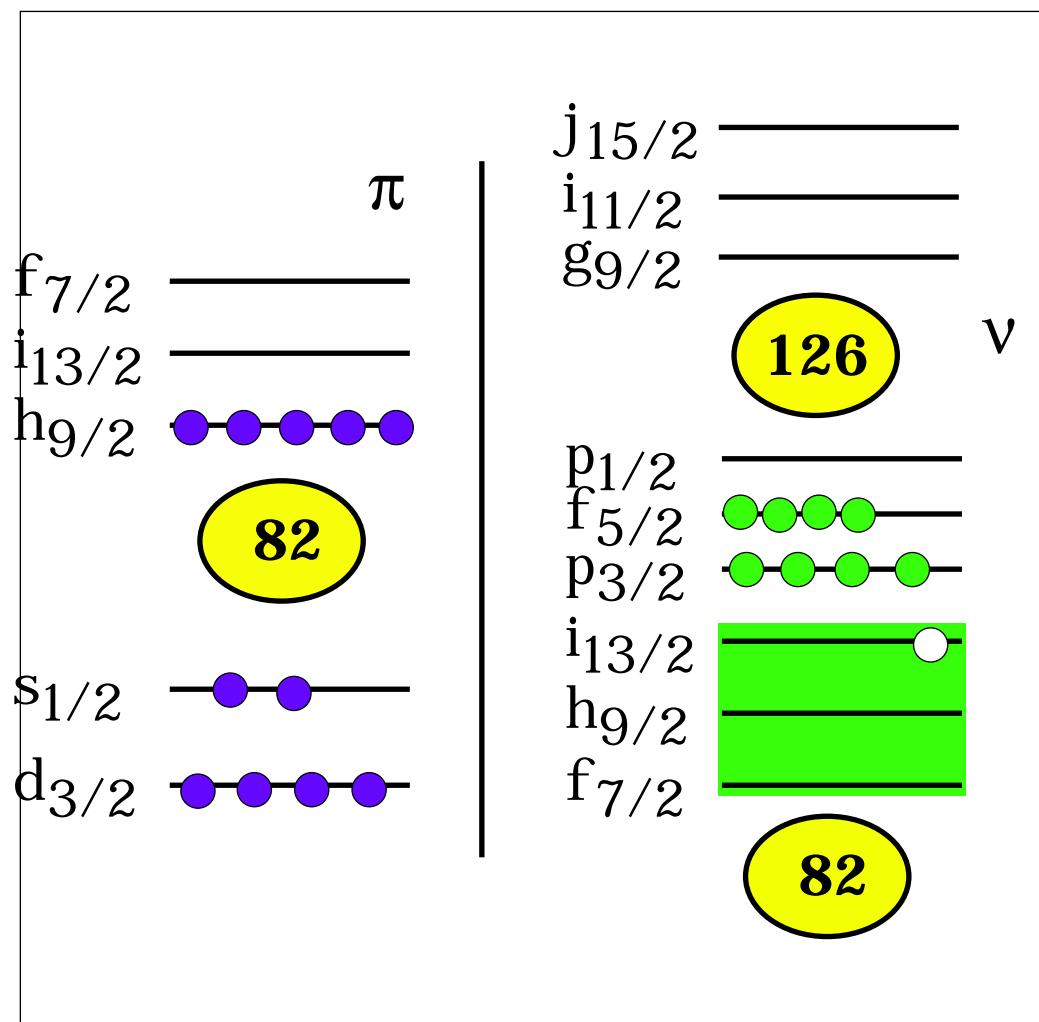
209 Fr
Meyer et al
Phys Rev C 73 (2006) 024307



Case I: the drastic solution...

$$\Gamma(E1, 194) = 1 \times 10^{-9} \text{ eV}$$

STRUCTURE IS IN ^{208}Fr !!



$^{208}\text{Fr} !$
GDD et al
EPJA 40(2009)127

CASE - II: ^{212}Rn
[4 valence protons; 0 neutron holes
plus neutron core excitations]
Signature E3 Transitions

See Phys Rev C-80-2009-05430
Phys Lett B 662-2008-19

Spherical levels and octupole coupling - E3 signatures

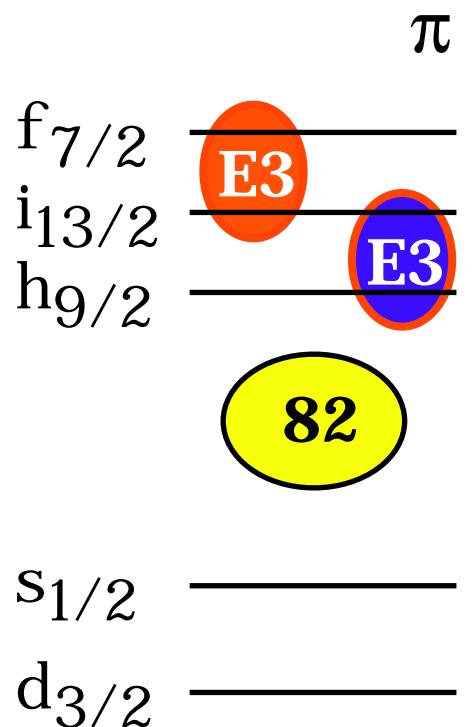
$$|i_{13/2}\rangle = \alpha |i_{13/2}\rangle + \beta |f_{7/2}\rangle \otimes |3^-\rangle + \gamma \dots$$

"type-A" E3 ~ 22 W.u.
(stretched)

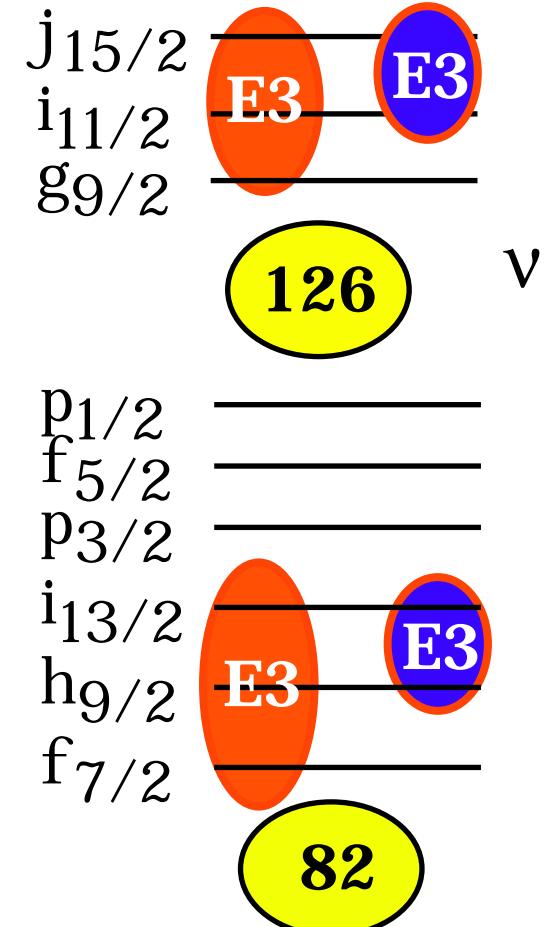
$\pi i_{13/2} \rightarrow f_{7/2}$
or
 $\nu j_{15/2} \rightarrow g_{9/2}$

"type-B" E3 ~ 3 W.u.
(spin-flip)

$\pi i_{13/2} \rightarrow h_{9/2}$
or
 $\nu j_{15/2} \rightarrow i_{11/2}$



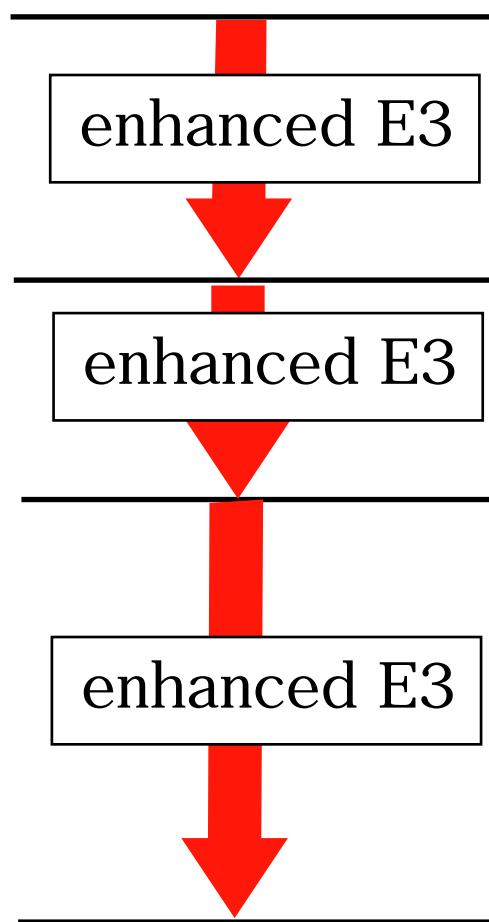
← collective E3 admixtures



see Bergstrom and Fant, Phys. Scr. 31 (1985) 26 and many others

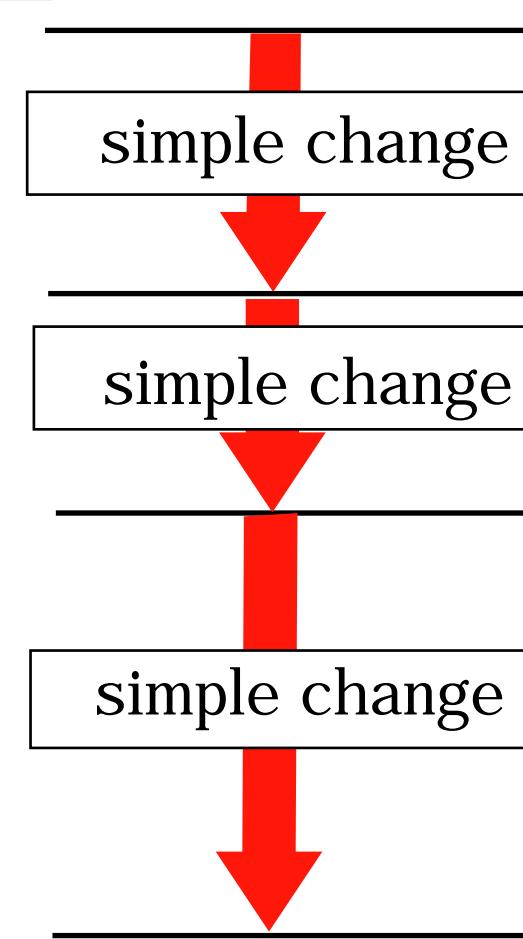
A simple proposition...

EXPERIMENT



must have
specifically
related
Configurations

"THEORY"



Low and High- Spin States (Isomers) in ^{212}Rn
Shell Model View
[Extremely Brief]

Semi-Empirical Shell-Model (ESM):

- Single-Particle Energies + weighted sum of two-body interaction matrix elements

The shell-model Hamiltonian for a many-body system may be written as

$$H = \sum_i H_i + \sum_{i>j} H_{i,j}, \quad (1)$$

where H_i and $H_{i,j}$ are single-particle and two-body terms, respectively.

In terms of the excitation energy of a state Ψ , we can write:

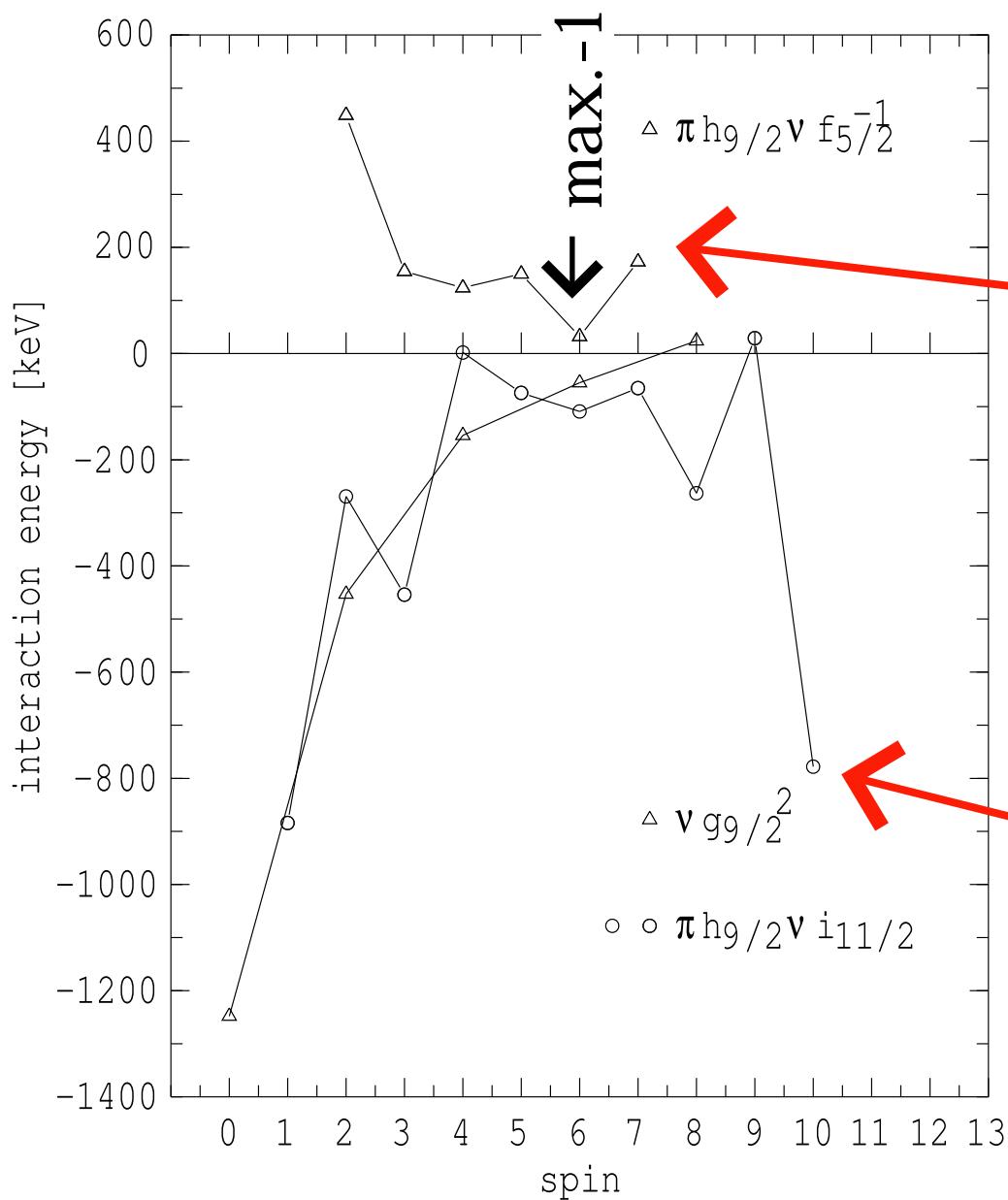
$$E_{\text{ex}}(\Psi) = \varepsilon_0 + \sum \langle j_a, j_b \rangle_J, \quad (2)$$

where ε_0 denotes the difference between the single-particle energy and the ground-state energy and $\langle j_a, j_b \rangle_J$ are the two-body residual interaction energies. In the present calculations interaction energies are taken from experiment where possible, and from theoretical models (Kuo and Herling [20]) where it is not. A full list of the interactions used in this work is tabulated in [21].

Blomqvist :
Proc Argonne Symposium, 1979 ANL/PHY-79-4, p155

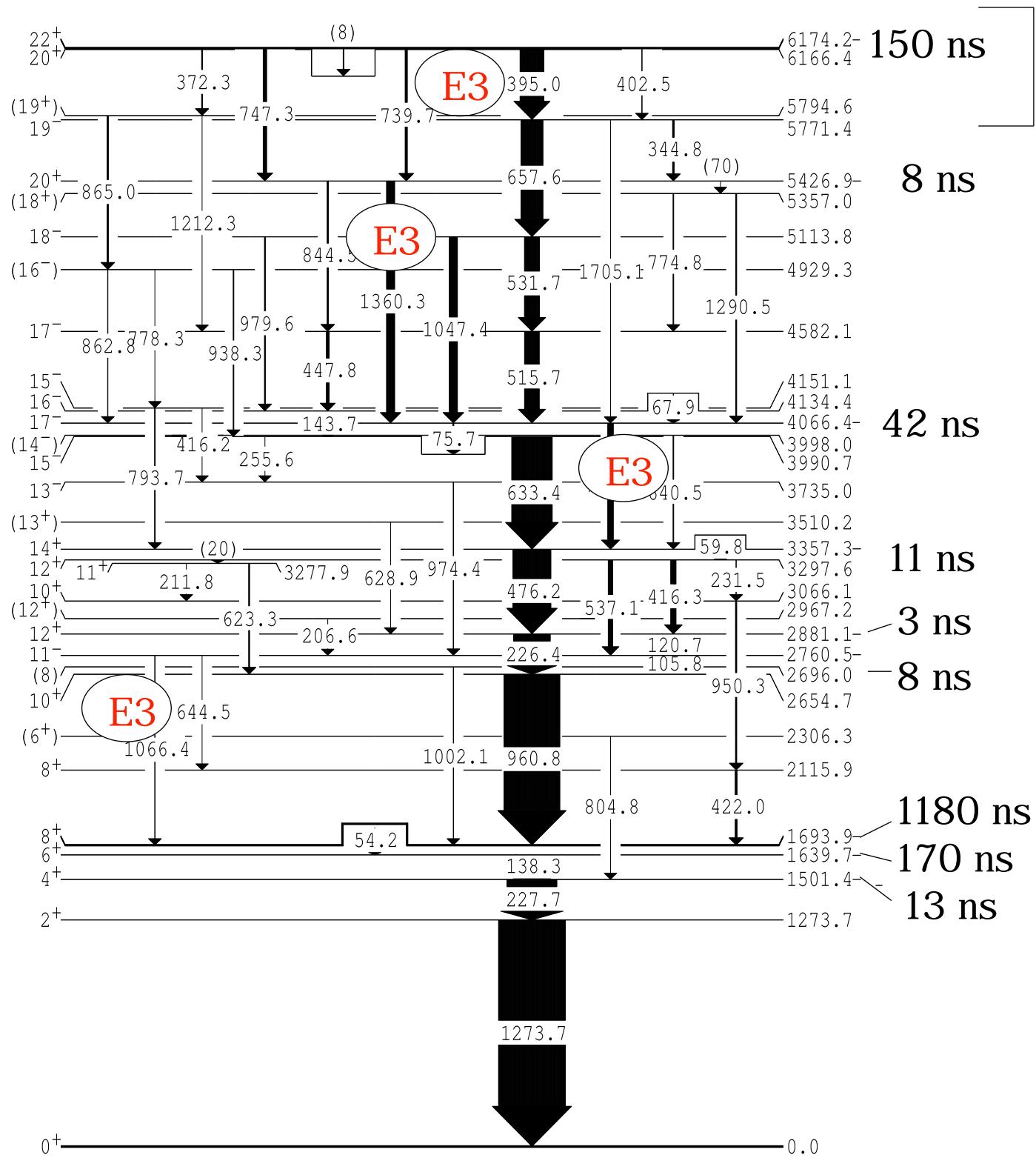
See: Bayer et al, NPA 694 (2001) 3; NPA 650 (1999) 3; etc.

Typical residuals



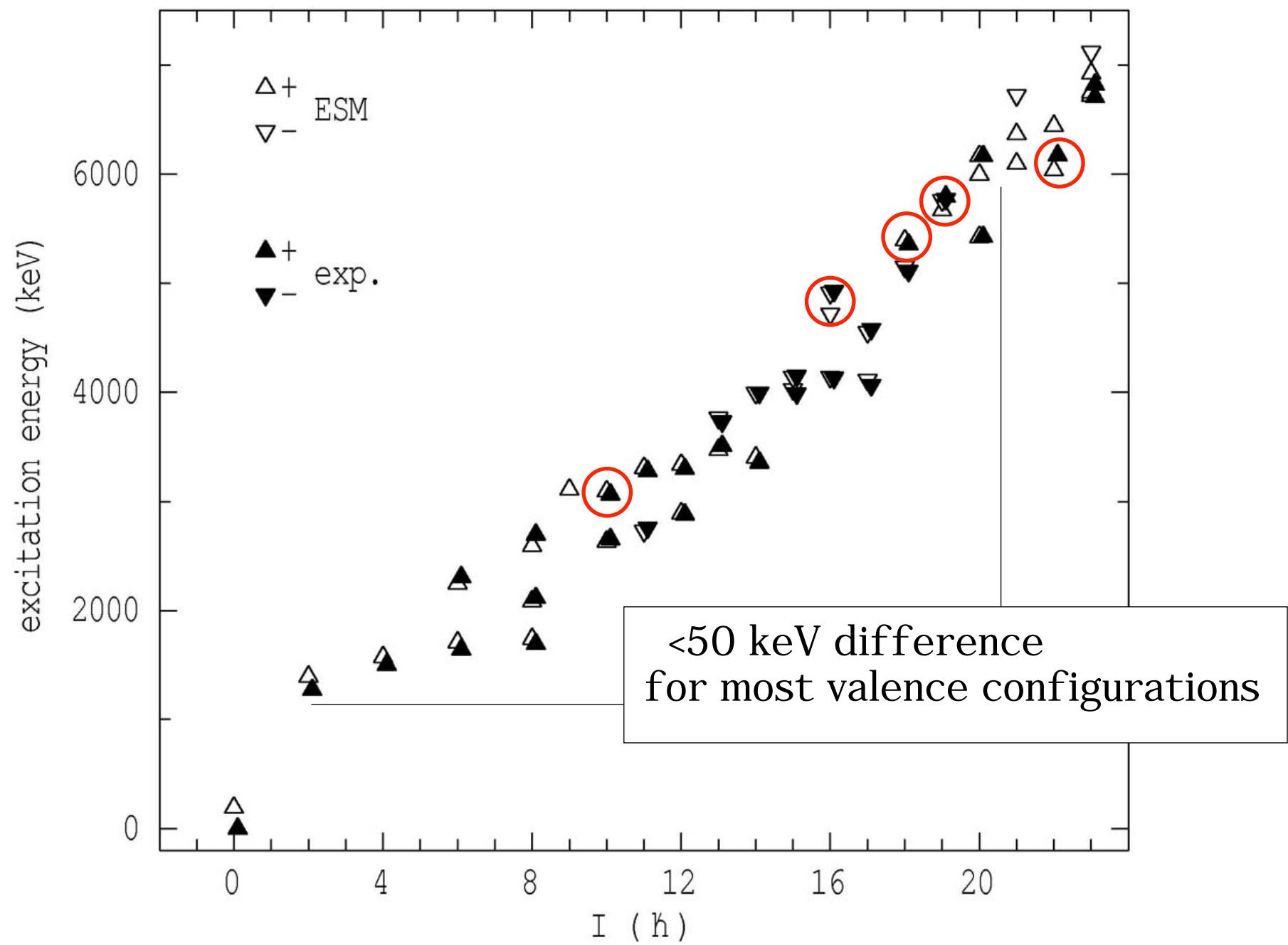
proton, neutron-hole
repulsive
for mutual alignment,
even for low-j neutron-hole

proton, neutron
strongly-attractive
for mutual alignment



22^+ , 20^+ and 19^-
core-excited states

^{212}Rn Z=86, N= 126 - Experiment and ESM - low



Strengths (sampler)

J^π		E(keV)	$\sigma\lambda$	W.u.	within multiplet E2 type
22^+	20^+_2	7.6	E2	4(1)	$\pi[h^3i] \nu p^{-1}g$ multiplet
20^+_1	747		E2	$3.6(7) \times 10^{-5}$	core-excited to valence
19^-_1	403		E3	44(11)	core-excited to core-excited $\pi i_{13/2} \tilde{\rightarrow} f_{7/2}$ type-A
20^+_1	18^+	70	E2	2.7(5)	valence multiplet
17^-_2	845		E3	37(8)	valence to valence $\pi i_{13/2} \tilde{\rightarrow} f_{7/2}$
17^-_1	1360		E3	7(1)	valence to valence $\pi i_{13/2} \tilde{\rightarrow} h_{9/2}$ type-B

Maximum- J ; Single (I), Double (II), Triple(III) - Neutron Core Excitations

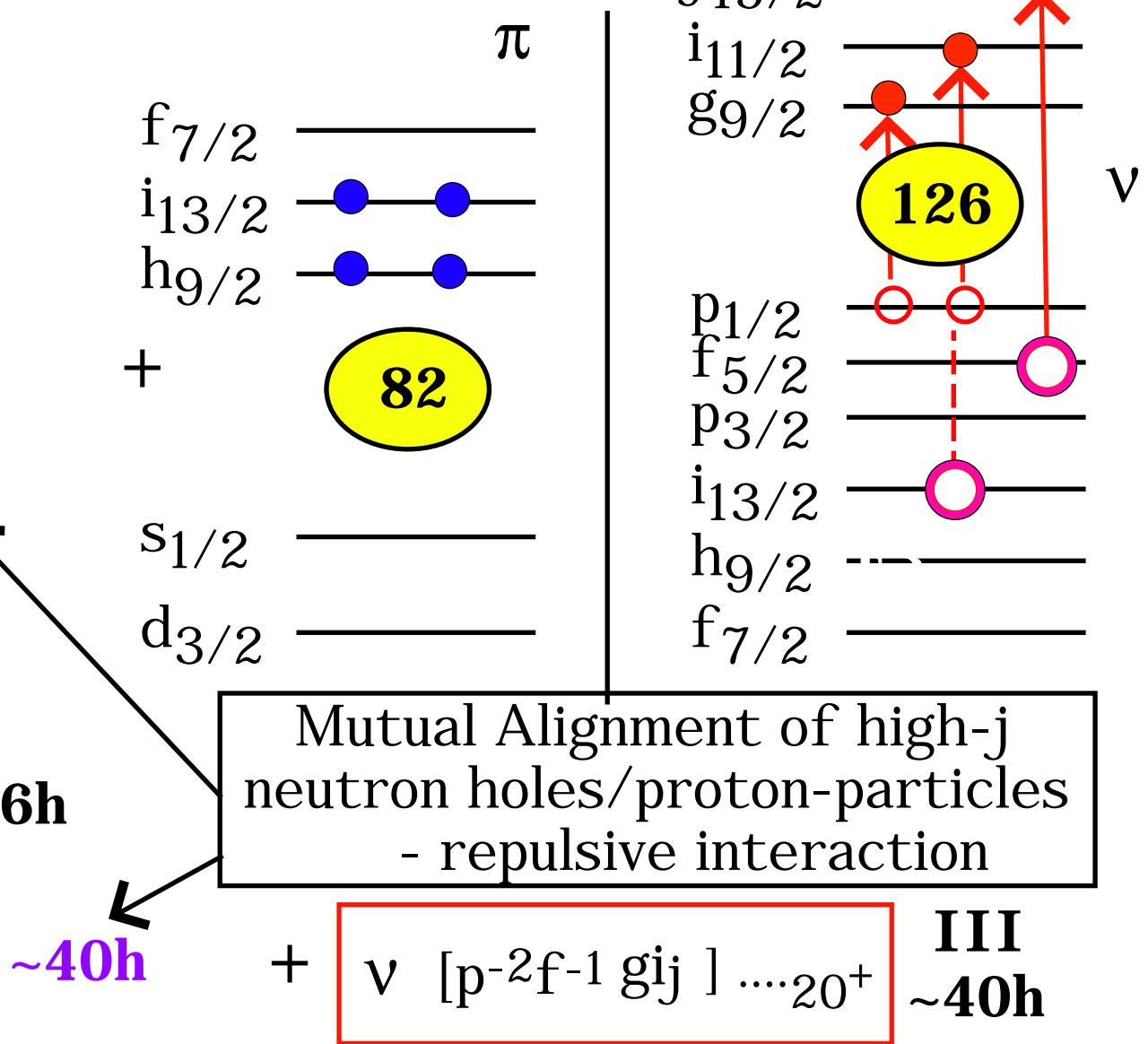
valence protons

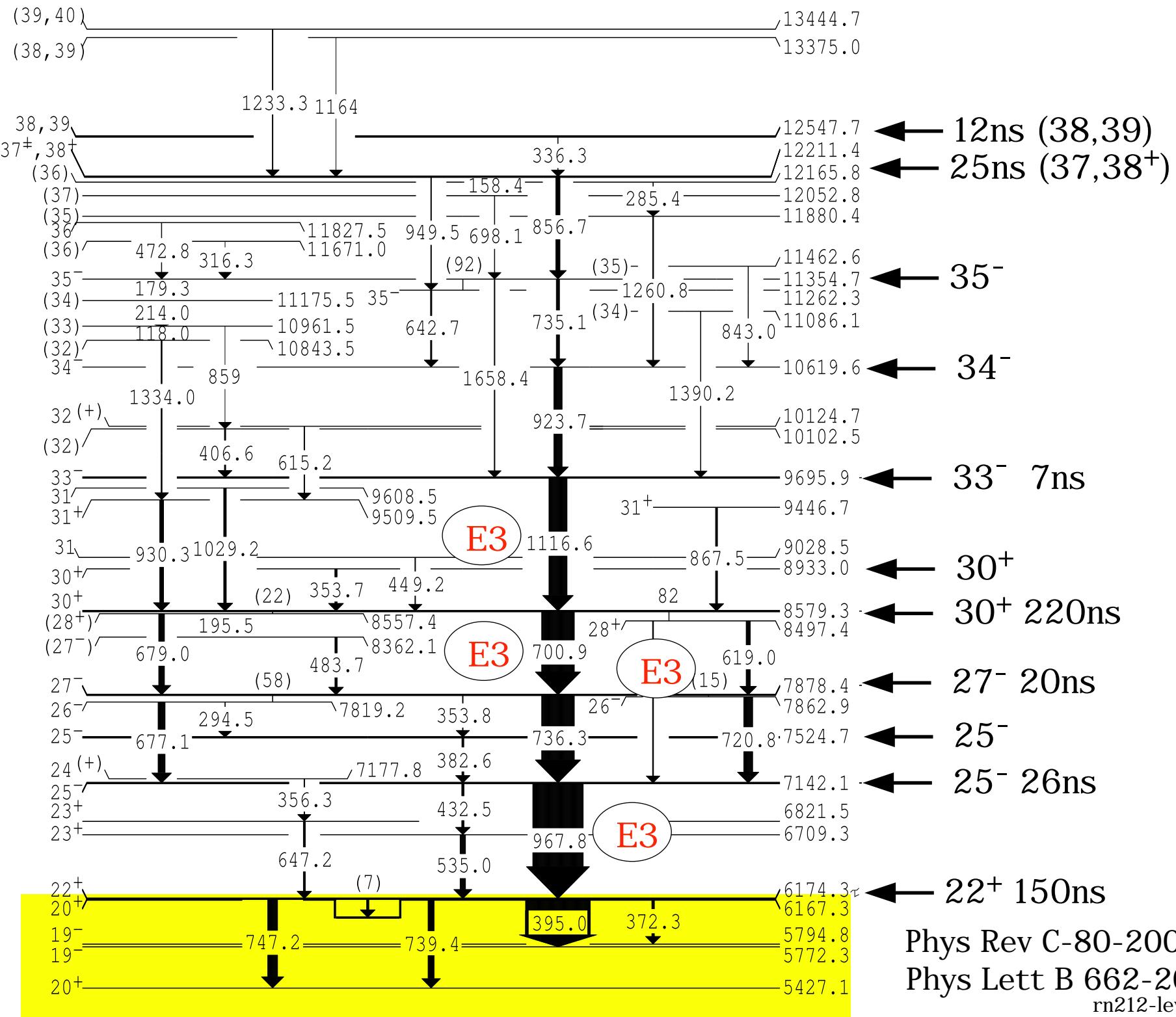
$\pi [h^3 f]_{17^-}$

$\pi [h^2 i^2]_{20^+}$

+ $\nu [p^{-1} g]_{5^-}$
 $\nu [p^{-1} i]_{6^-}$
 $\nu [p^{-1} j]_{8^+}$
 $\nu [f^{-1} g]_{7^-}$ **I**
~28h

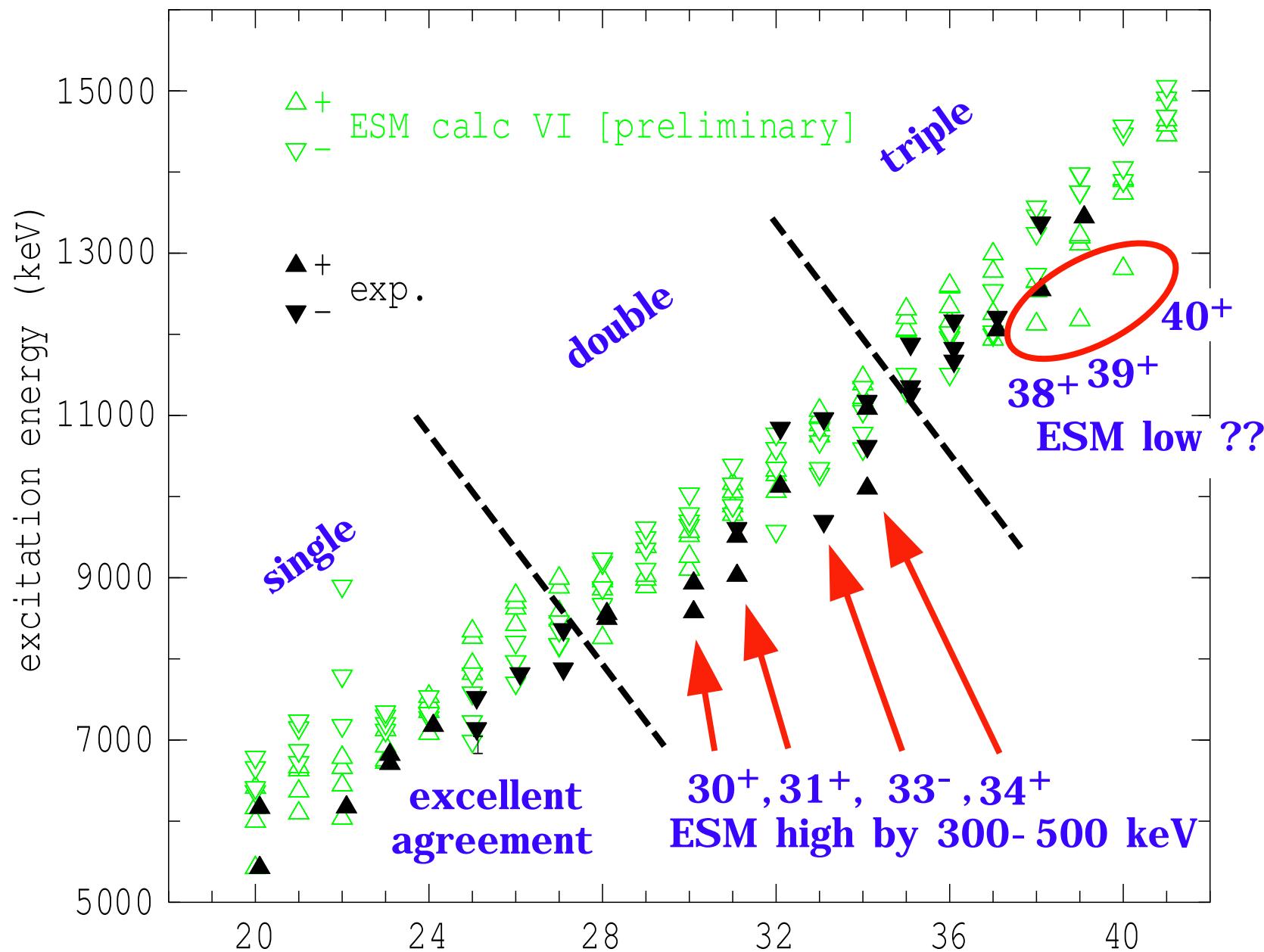
+ $\nu [p^{-2} g i]_{10^+}$
 $\nu [p^{-2} j^2]_{14^+}$
 $\nu [p^{-1} f^{-1} g j]_{16^-}$
 $\nu [p^{-1} i^{-1} i j] \dots 20^+$ **II**
~36h





Phys Rev C-80-2009-05430
 Phys Lett B 662-2008-19
 rn212-levels-surrey-I.md

^{212}Rn Z=86, N= 126 - Experiment and ESM - preliminary



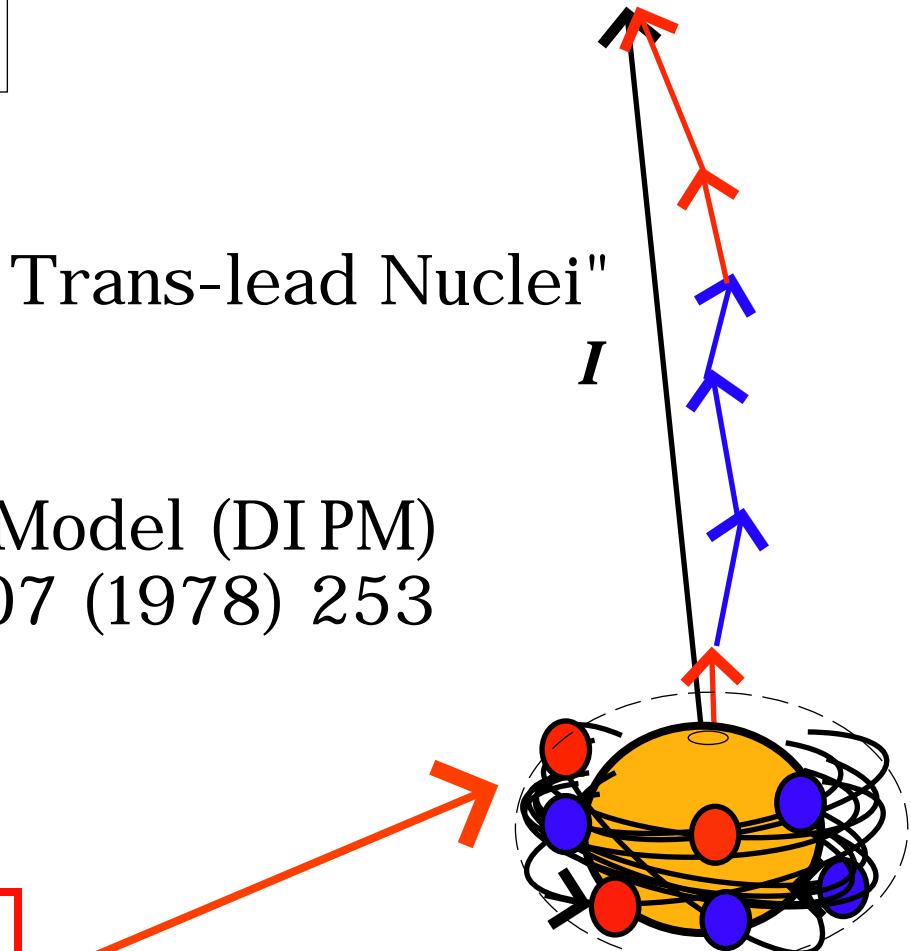
Deformed View (DIPM)

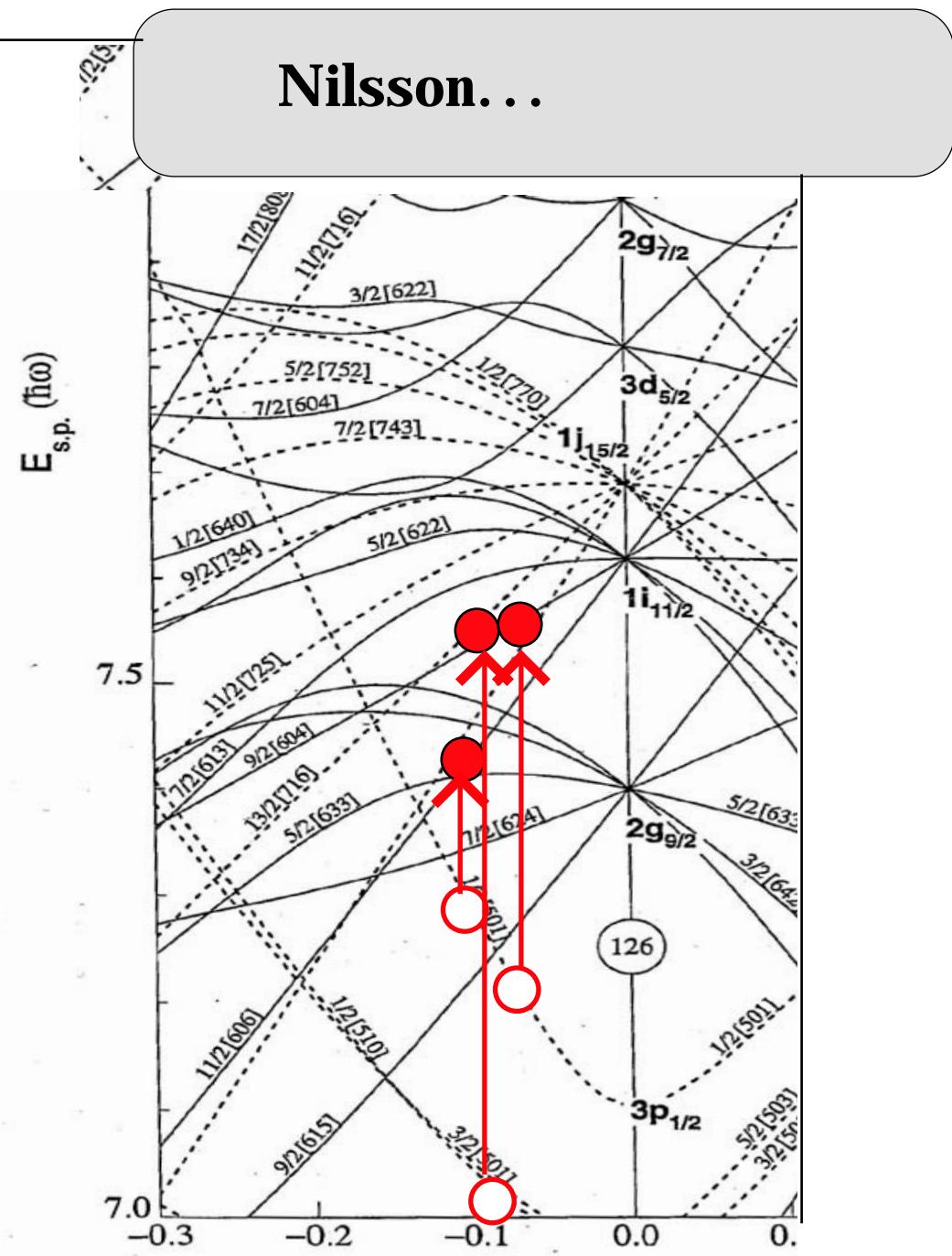
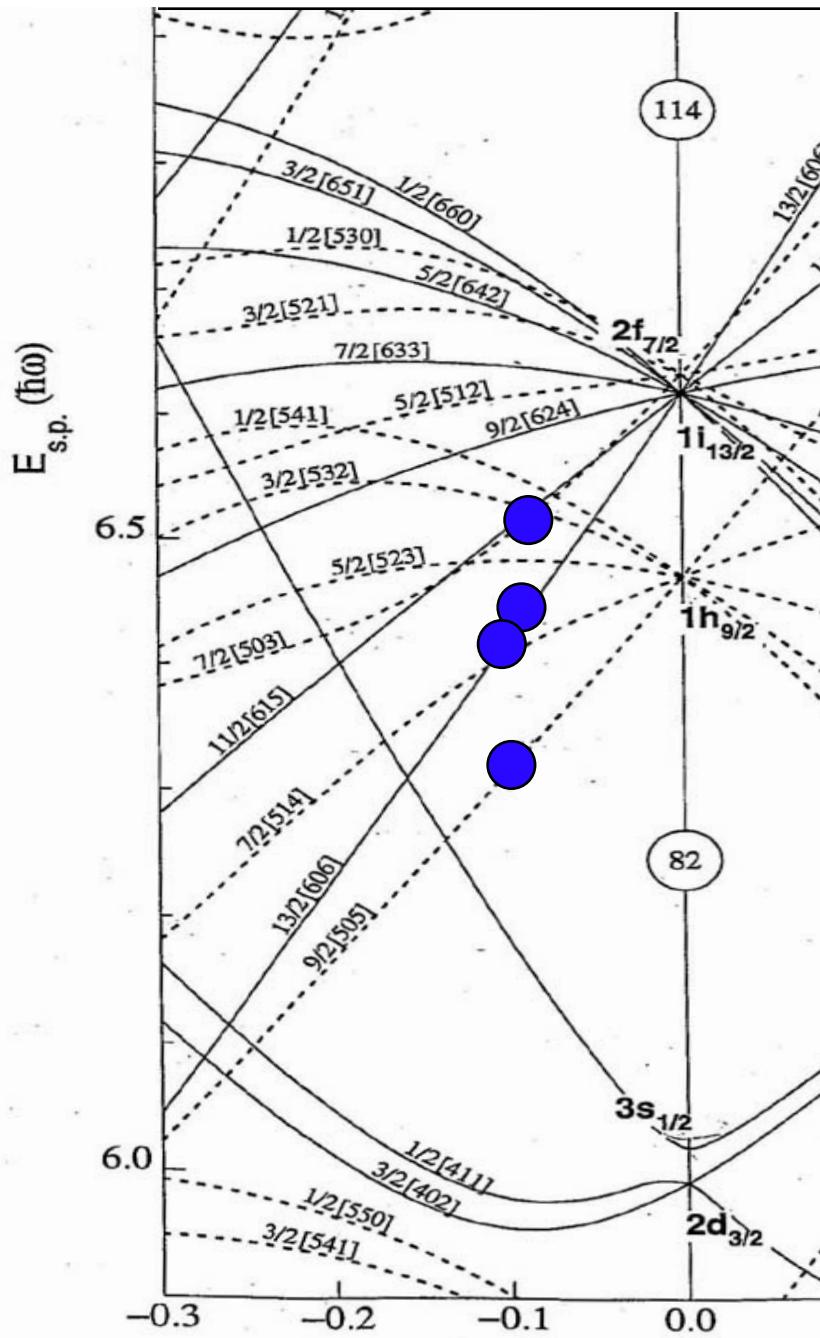
Induced Deformation

"New Era of High-Spin Studies in Trans-lead Nuclei"

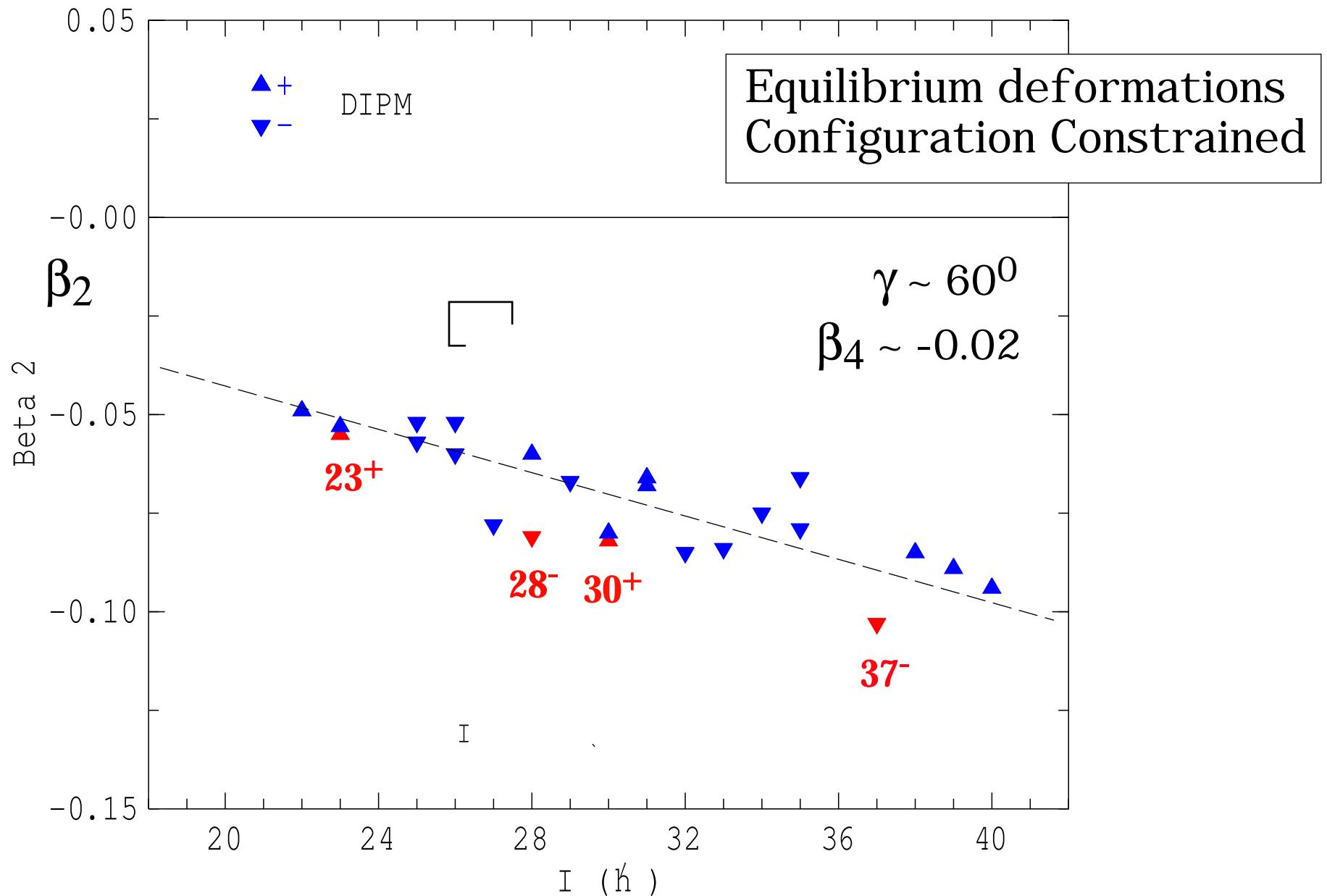
Deformed Independent Particle Model (DIPM)
Matsuyanagi et al Nucl Phys A 307 (1978) 253

aligned high-spin states
equatorial nucleons congregate
- oblate deformation ?

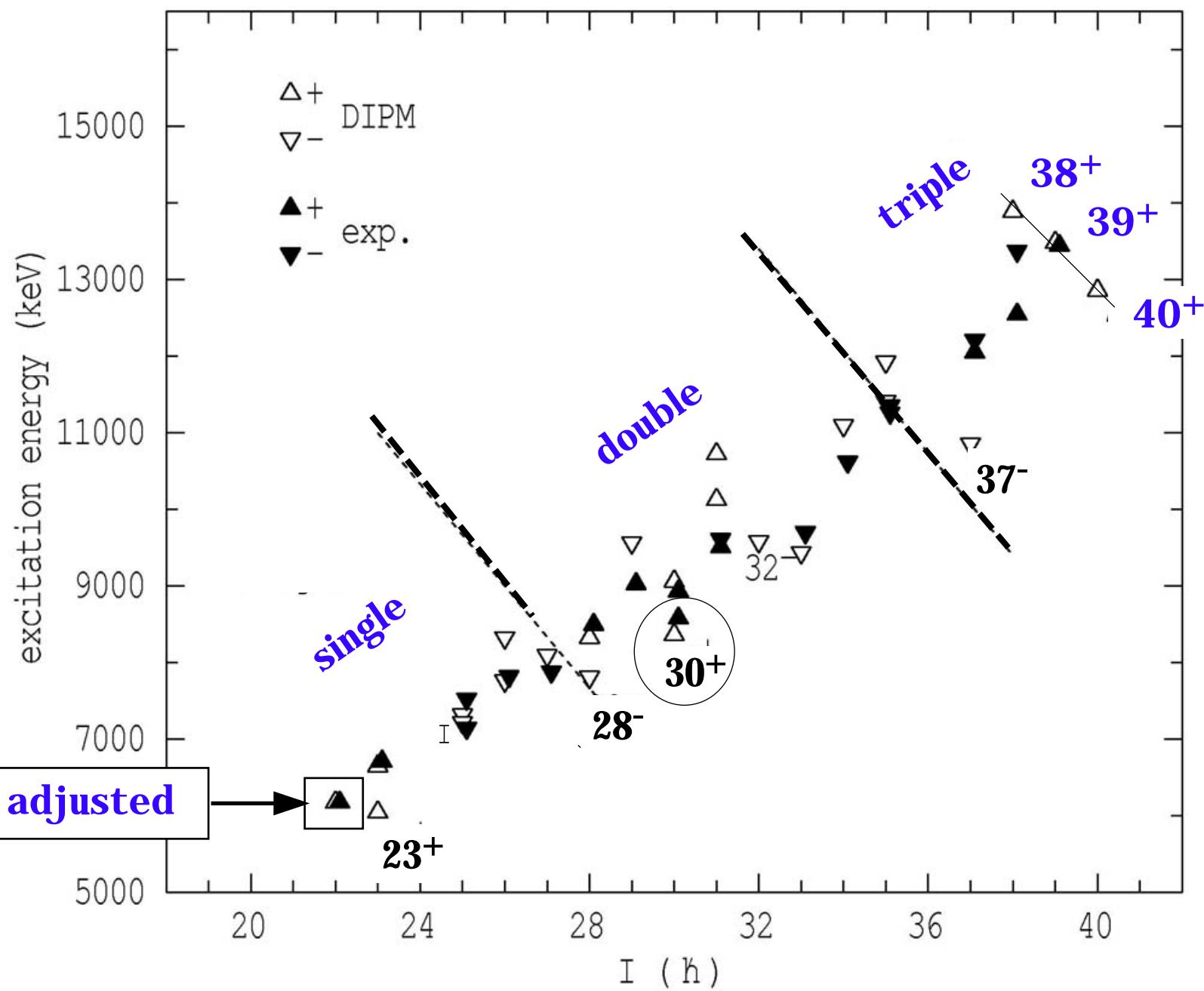




DIPM - predicted deformations (Xu, Liu)

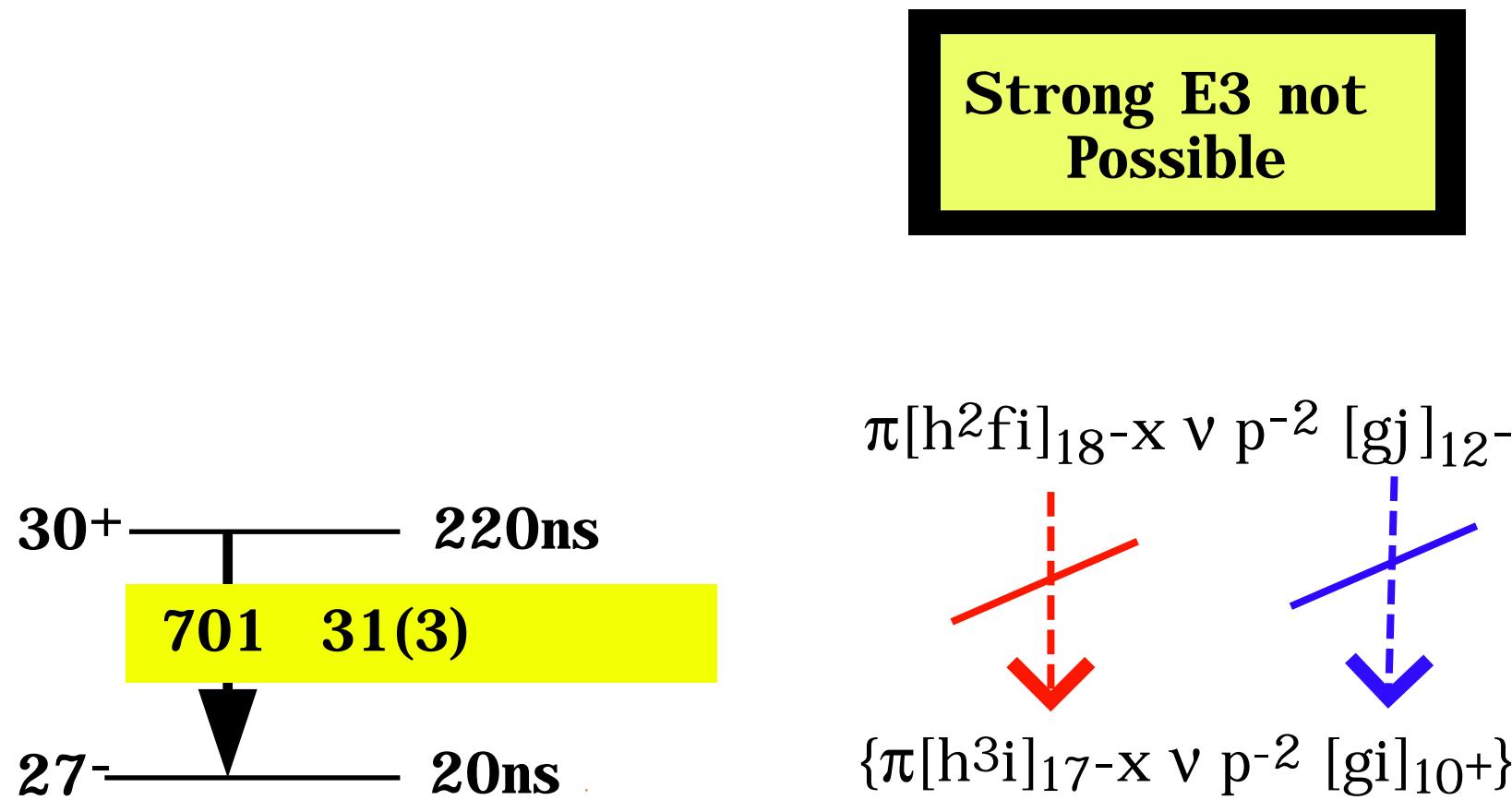


^{212}Rn Z=86, N= 126 - Experiment and DIPM

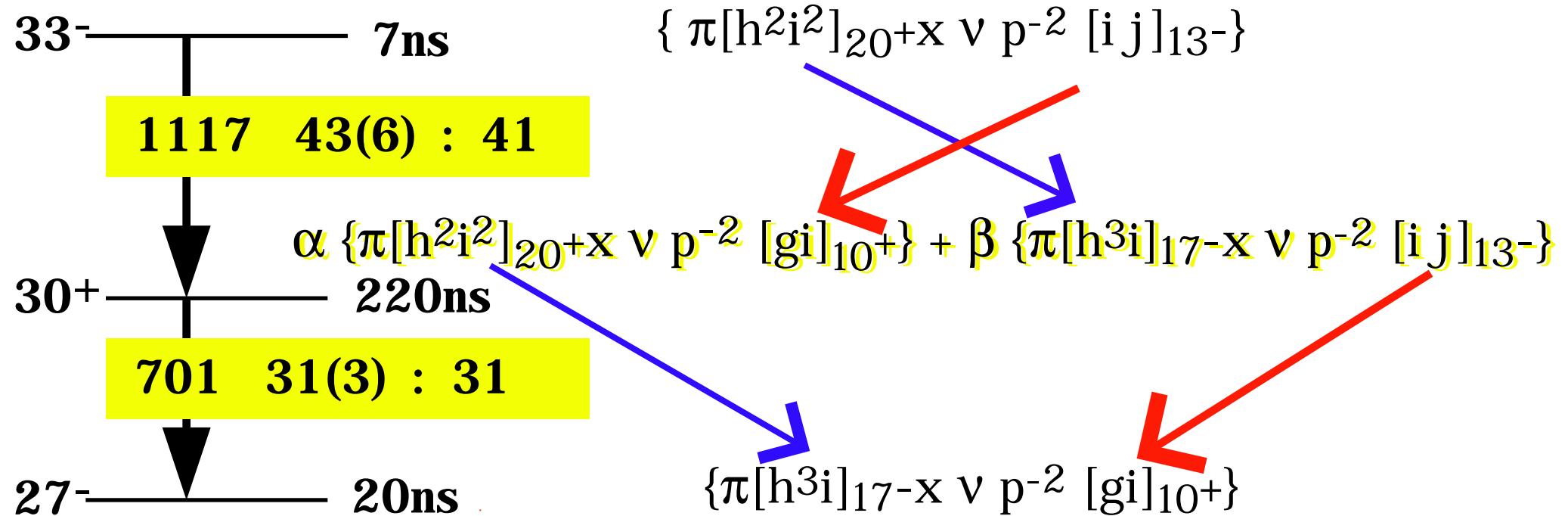


**DIPM Configuration assignments
do not match E3 transitions**

**DIPM Assignment - Matsuyanagi et al.,
Andersson et al; deVoight review...**



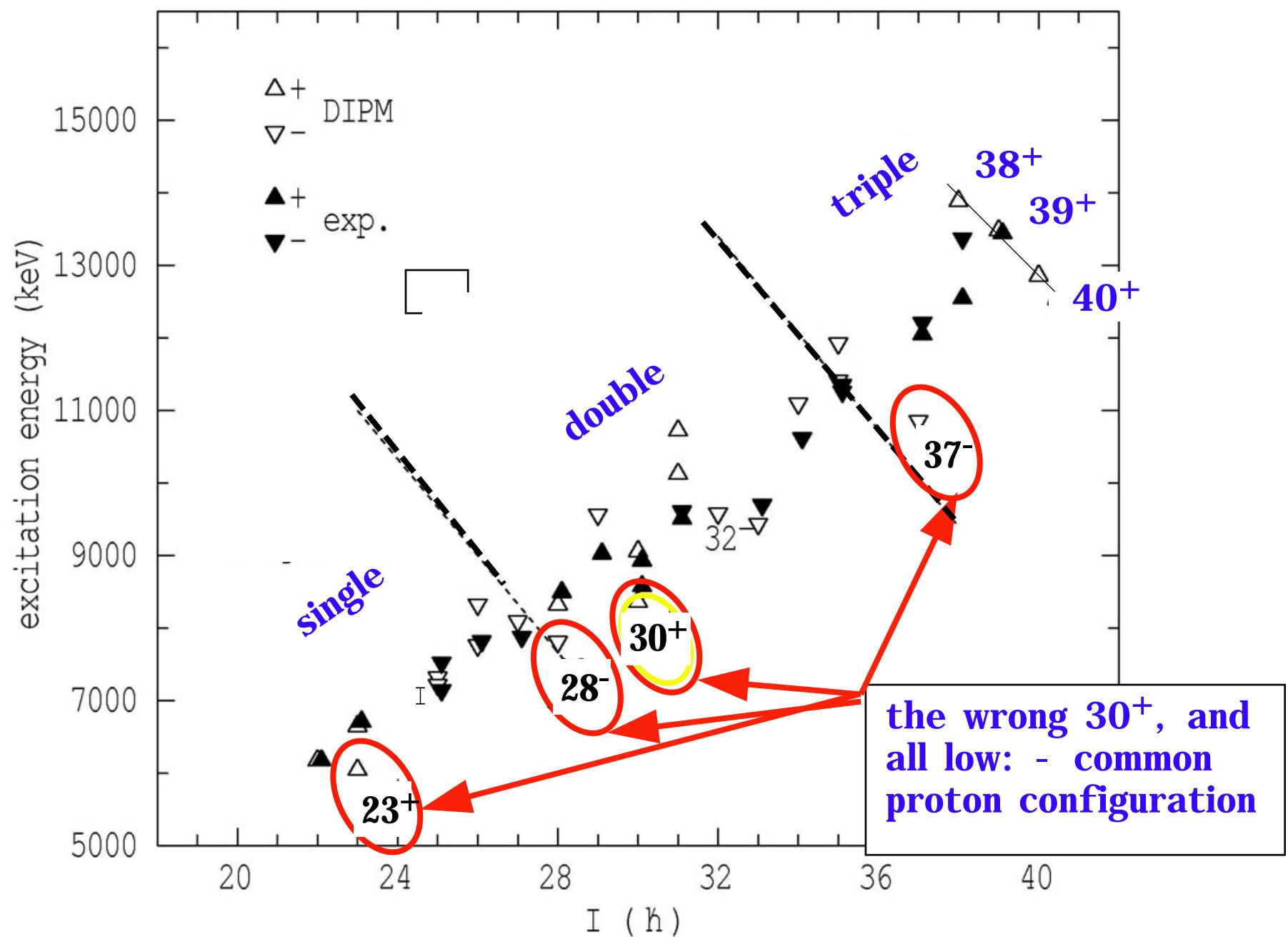
Multiple- Particle- Octupole- Coupling: MPOC



Poletti et al NPA 448 (1986) 189
 Dracoulis et al PLB 246 (1990) 31

"Automatic" mixing
 Constructive Interference

^{212}Rn Z=86, N= 126 - Experiment and DIPM - preliminary



Why is 18⁻ too low in DI PM?

$$23^+ \pi (h^2 f i)_{18^-} \times v(p^{-1} g)_{5^-}$$

$$28^- \pi (h^2 f i)_{18^-} \times v(p^{-2} g i)_{10^+}$$

$$30^+ \pi (h^2 f i)_{18^-} \times v(p^{-2} g i)_{10^+} \leftarrow \text{the original problem}$$

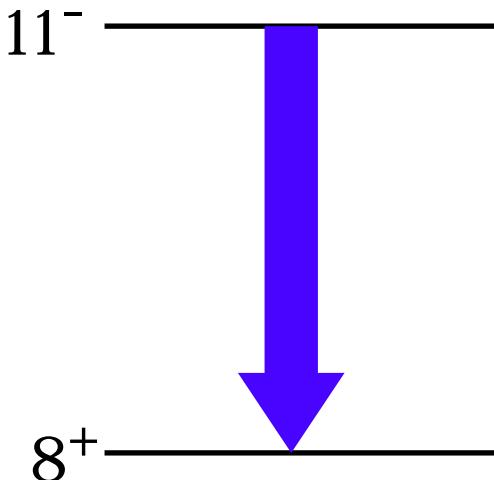
$$37^- \pi (h^2 f i)_{18^-} \times v(p^{-2} f^{-1} g i j)_{19^+}$$

CASE III; Pb- Po, 11⁻

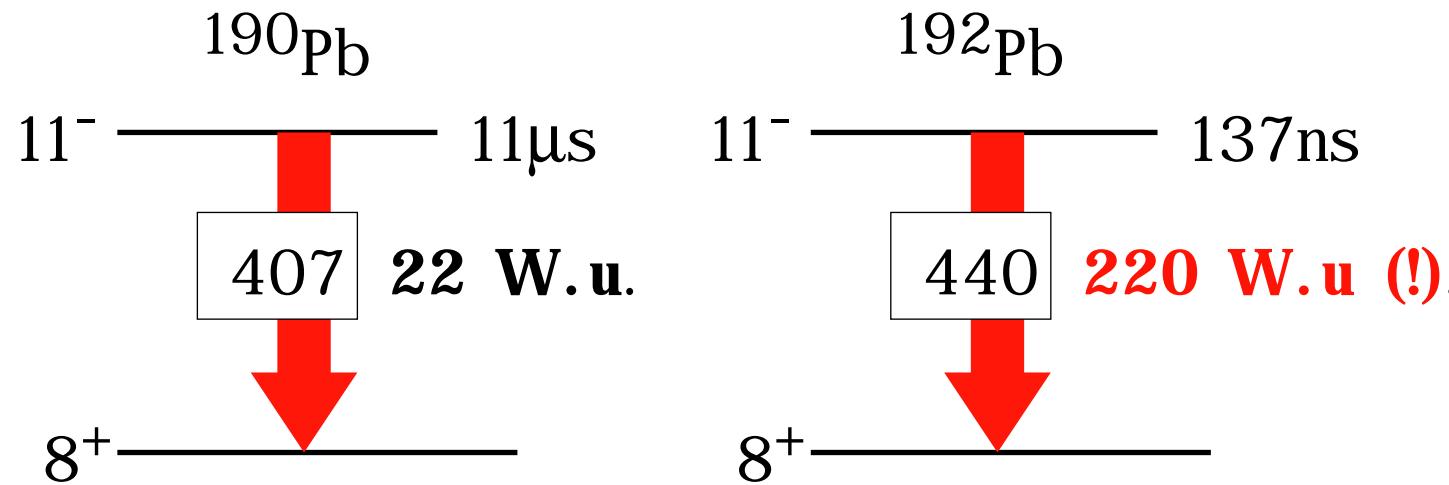
11⁻ to 8⁺ E3 ; naive expectation

$\pi [i_{13/2} h_{9/2}] 11^- \text{ to } [h_{9/2}^2] 8^+$

"type-B" E3 ~ 3 W.u.
(spin-flip)



11^- to 8^+ E3 ; emerging reality

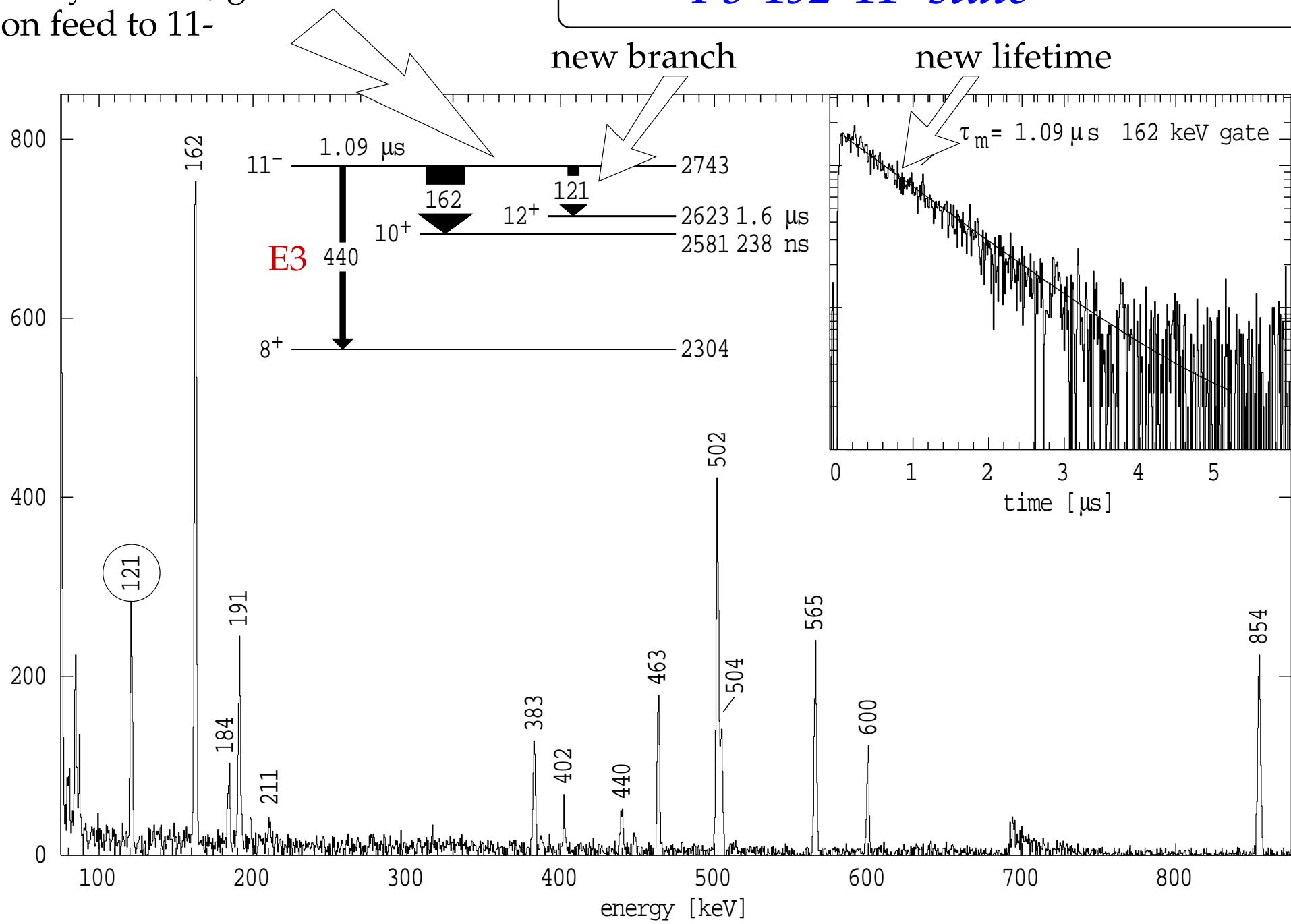


Dracoulis et al
PLB 432 (1998)37

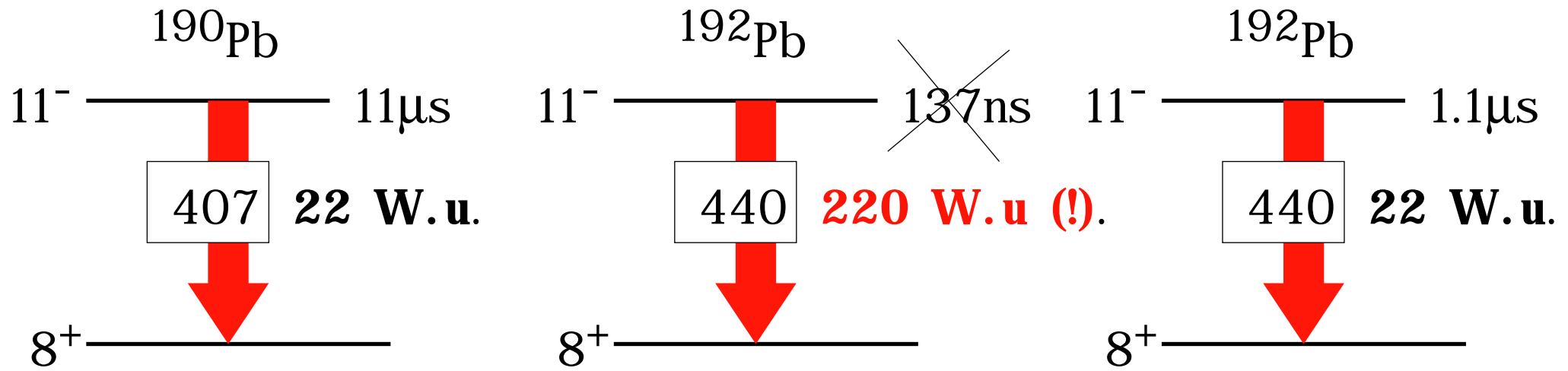
Literature
values of Lifetime and
Branching

Pb-192 11- state

delayed lines; gated
on feed to 11-



11⁻ to 8⁺ E3 ; emerging reality



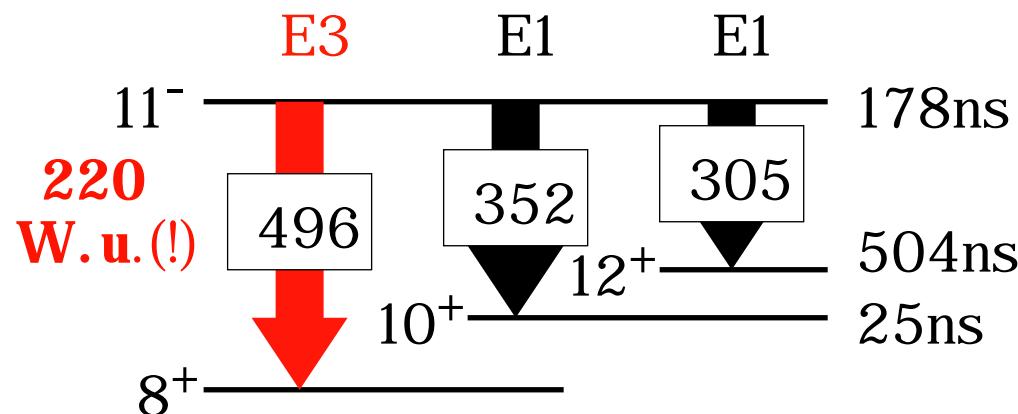
Dracoulis et al
PLB 432 (1998)37

Literature
values of Lifetime and
Branching

Dracoulis et al
PRC 63 (2001)
061302

"type-A" E3 ~ 20 W.u

11^- to 8^+ E3 ; ^{194}Pb

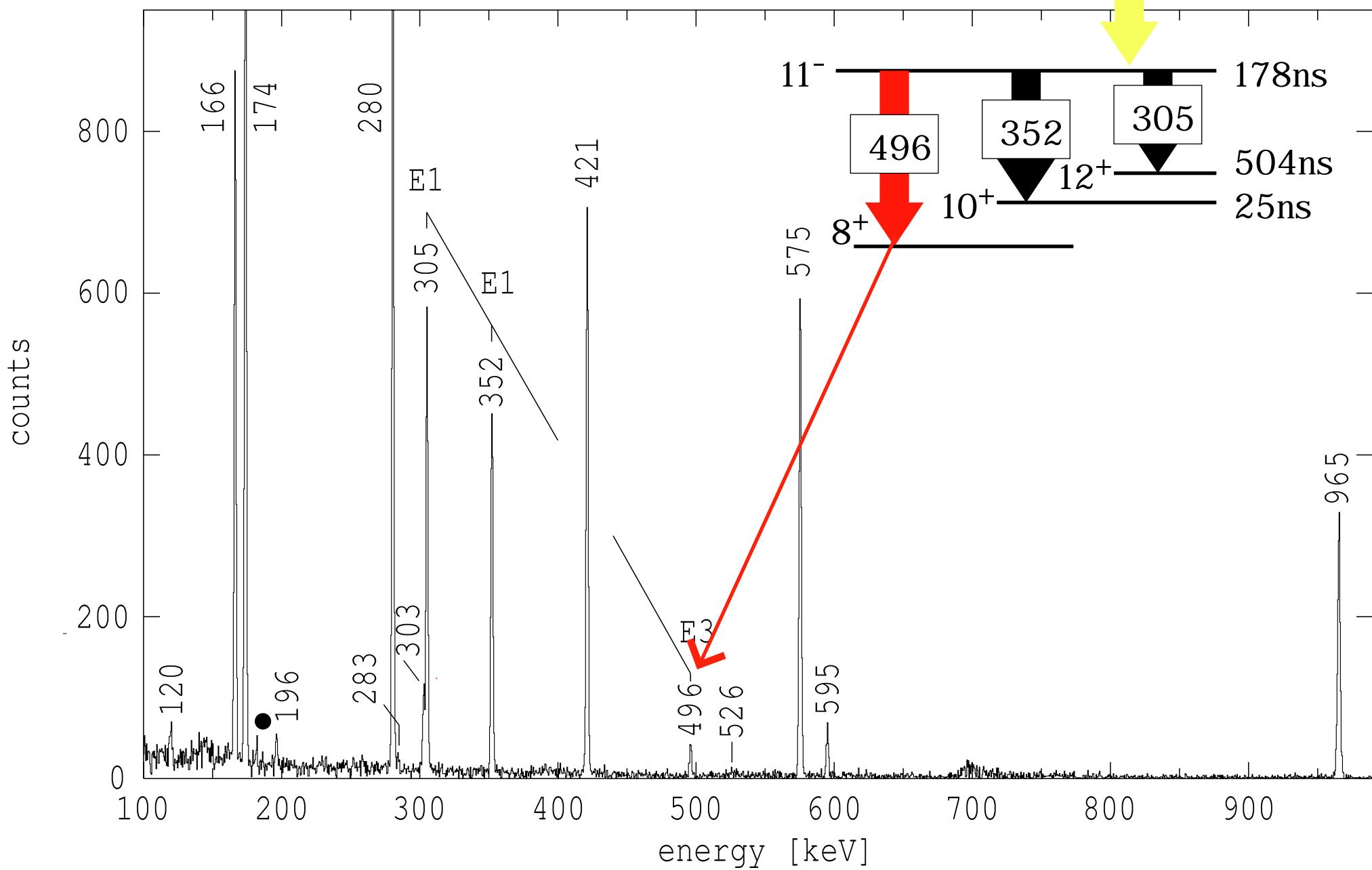


$$I\gamma(496) = I\gamma(352) = I\gamma(305) ?$$

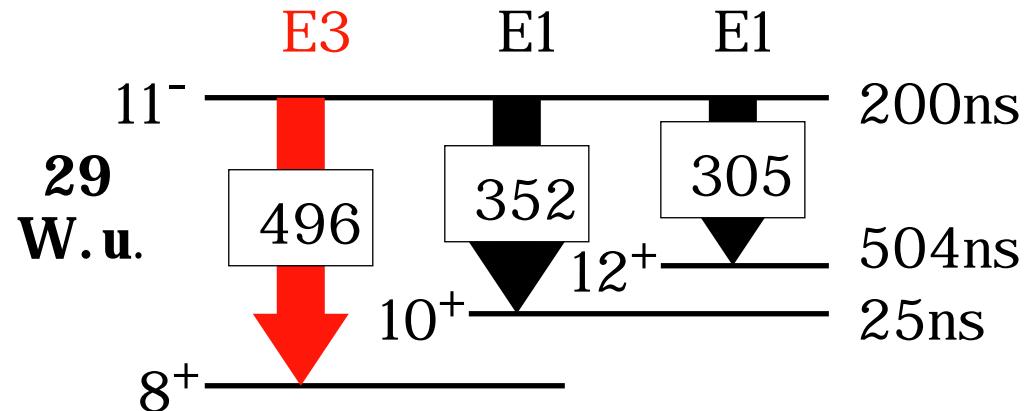
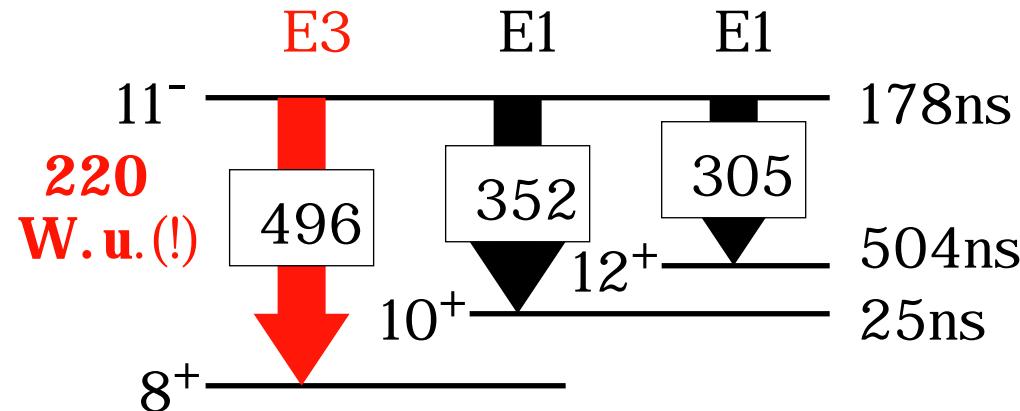
Van Ruyven et al (1986)
Fant et al (1991)
Kaci et al (2002)

11^- Isomer decay ^{194}Pb

γ - ray gates above plus time gate



11⁻ to 8⁺ E3 ; ¹⁹⁴Pb



I γ (496) = I γ (352) = I γ (305) ?

I γ (496) << I γ (352) = I γ (305) ?

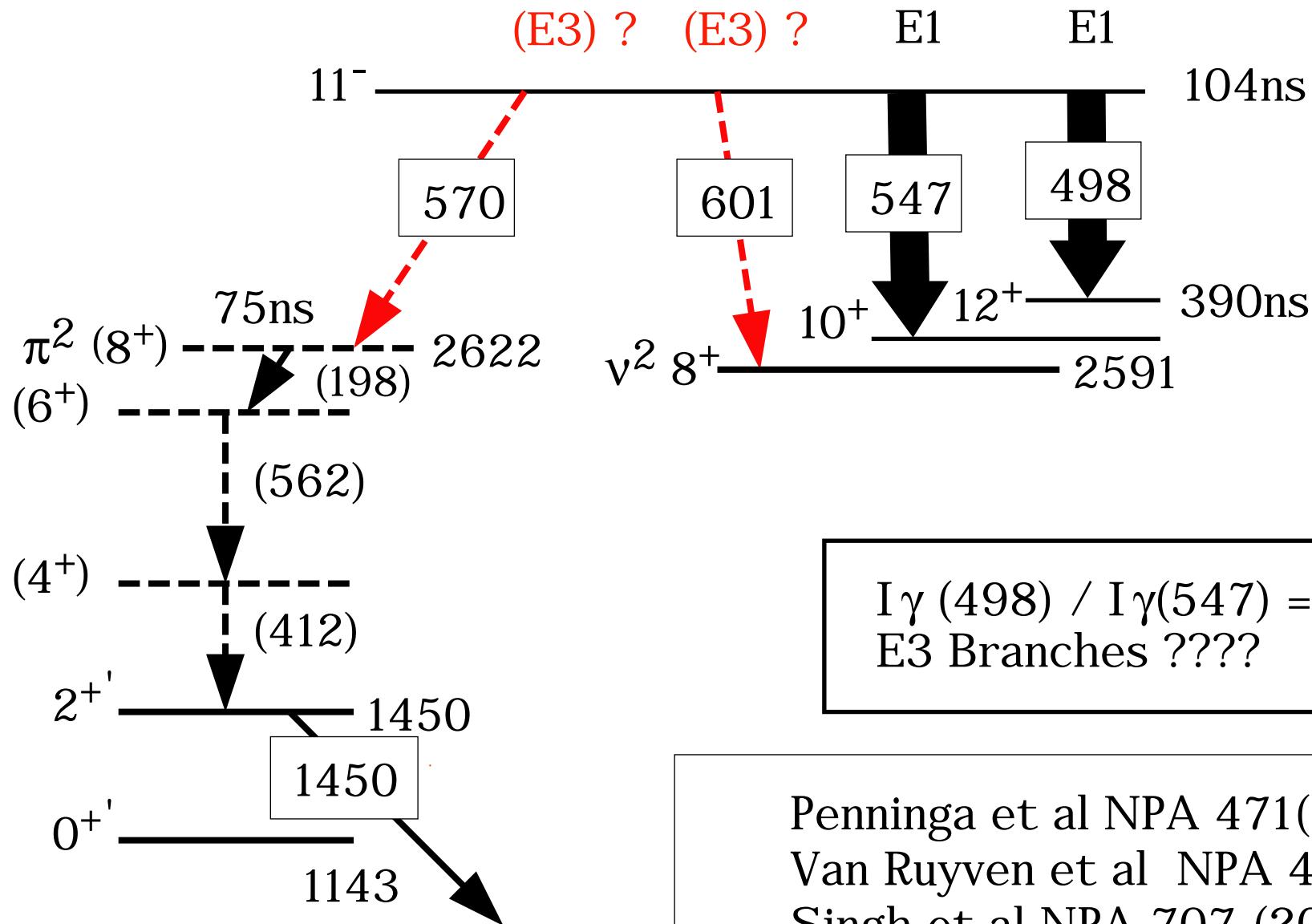
Van Ruyven et al (1986)
 Fant et al (1991)
 Kaci et al (2002)

"type-A" E3 ~ 20 W.u

... to be published

11⁻ to 8⁺ E3 ; ¹⁹⁶Pb challenge

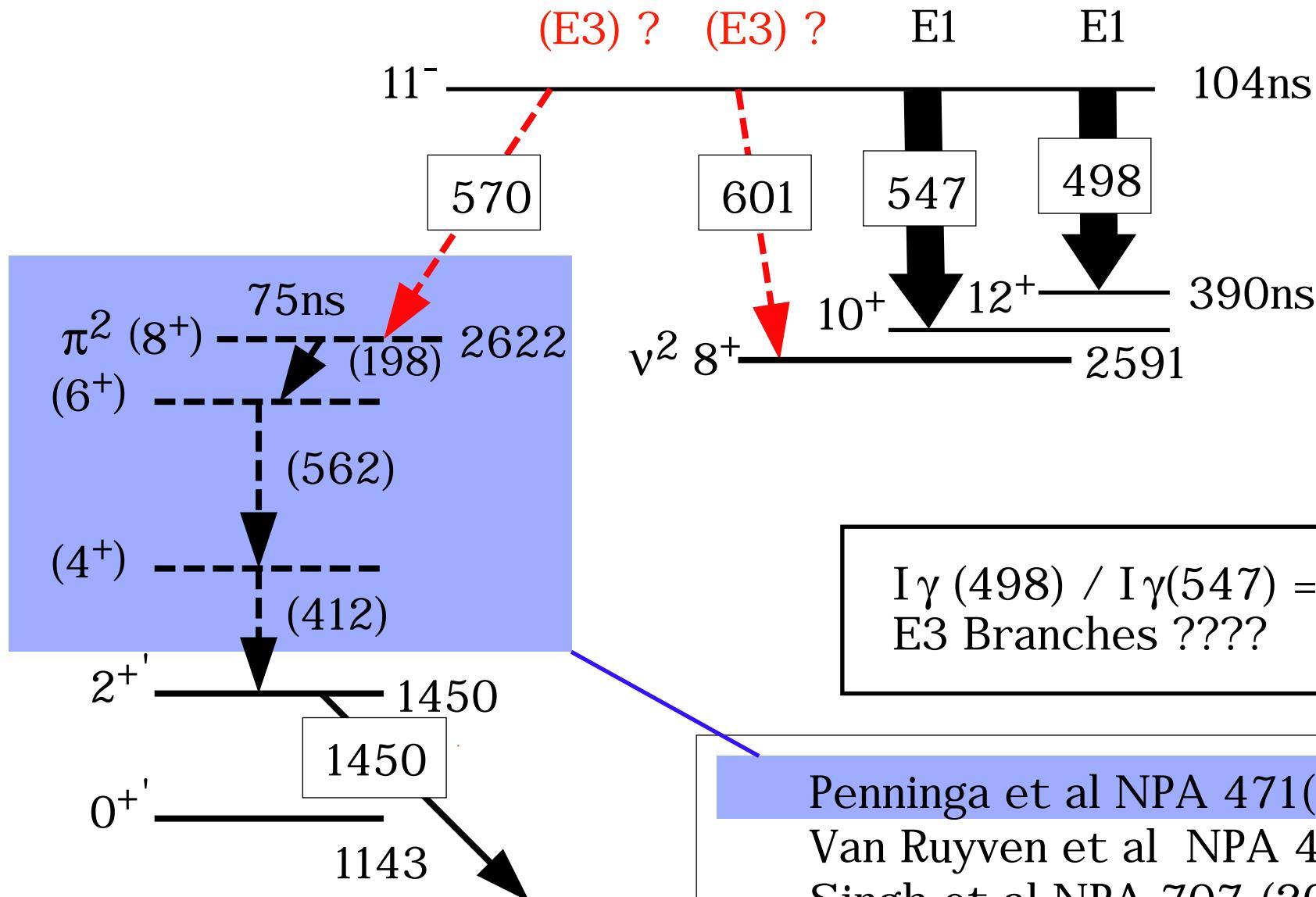
short- τ = weak branch



$I\gamma(498) / I\gamma(547) = ???$
E3 Branches ???

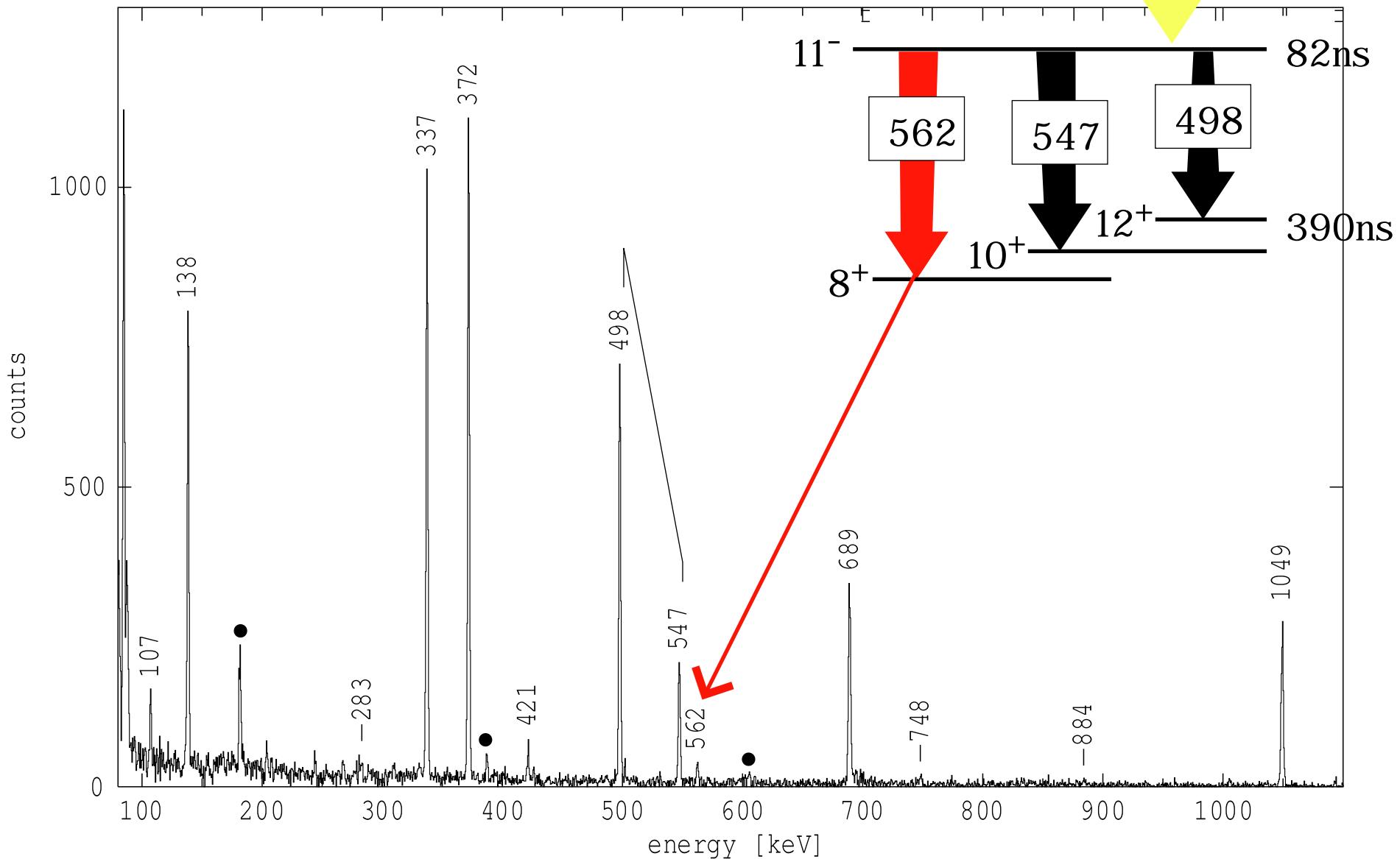
Penninga et al NPA 471(1987)535
Van Ruyven et al NPA 449(1986)
Singh et al NPA 707 (2002) 3

11⁻ to 8⁺ E3 ; ^{196}Pb challenge

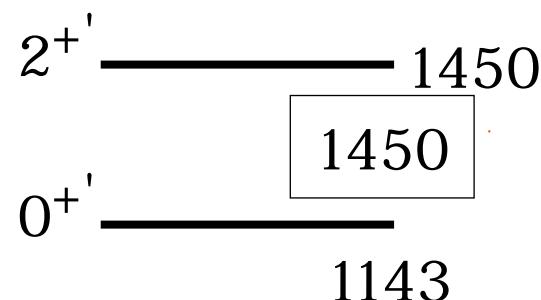
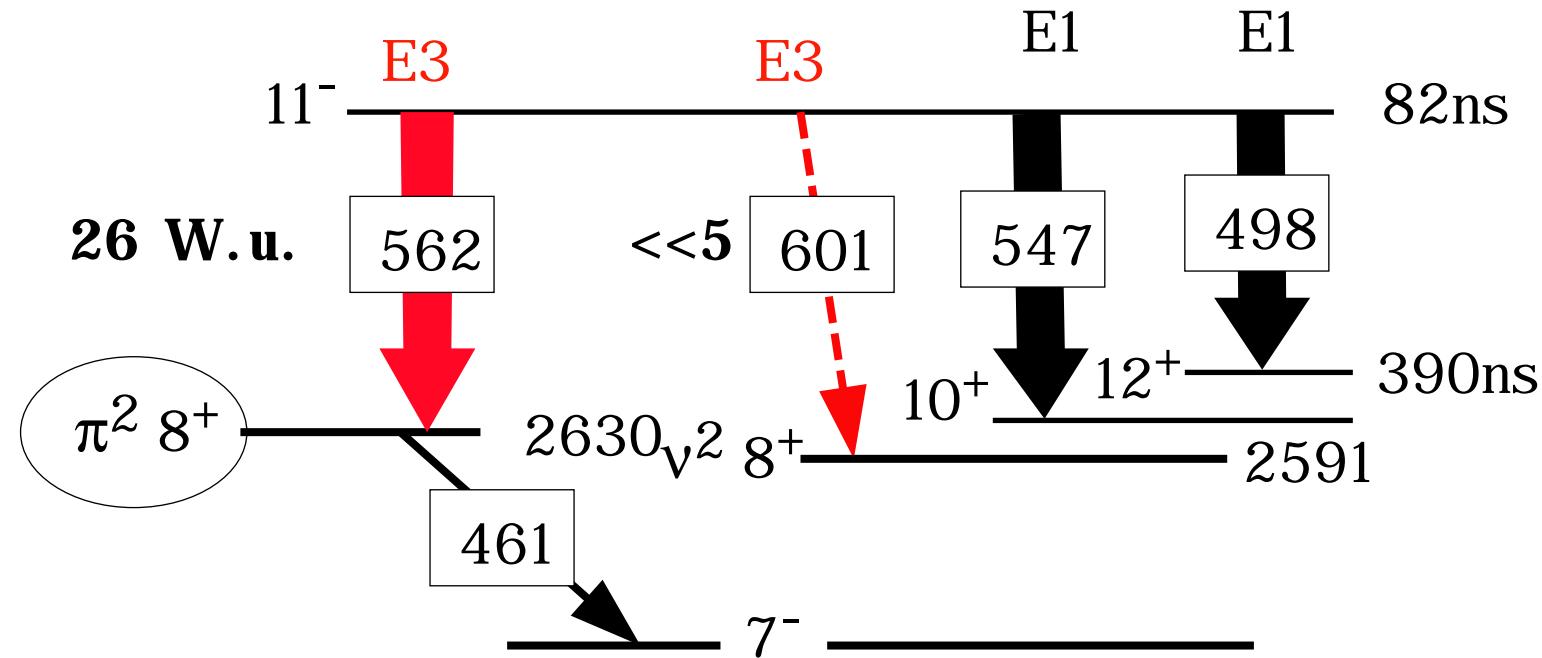


11^- Isomer decay ^{196}Pb

γ -ray gates above plus time gate



11⁻ to 8⁺ E3 ; ¹⁹⁶Pb solution ?

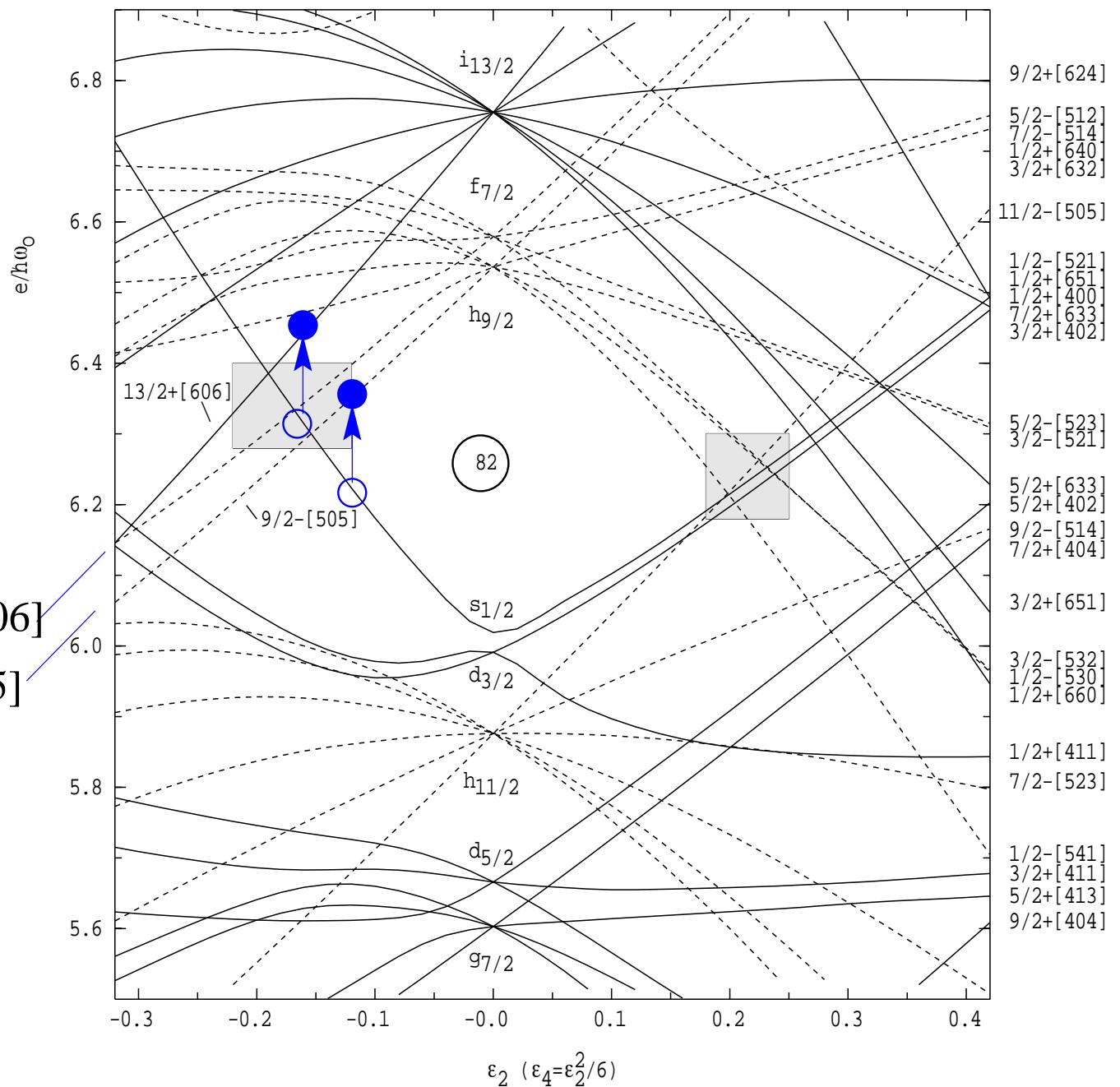


$I\gamma(498) / I\gamma(547) = 3$
E3 Branch 562 keV

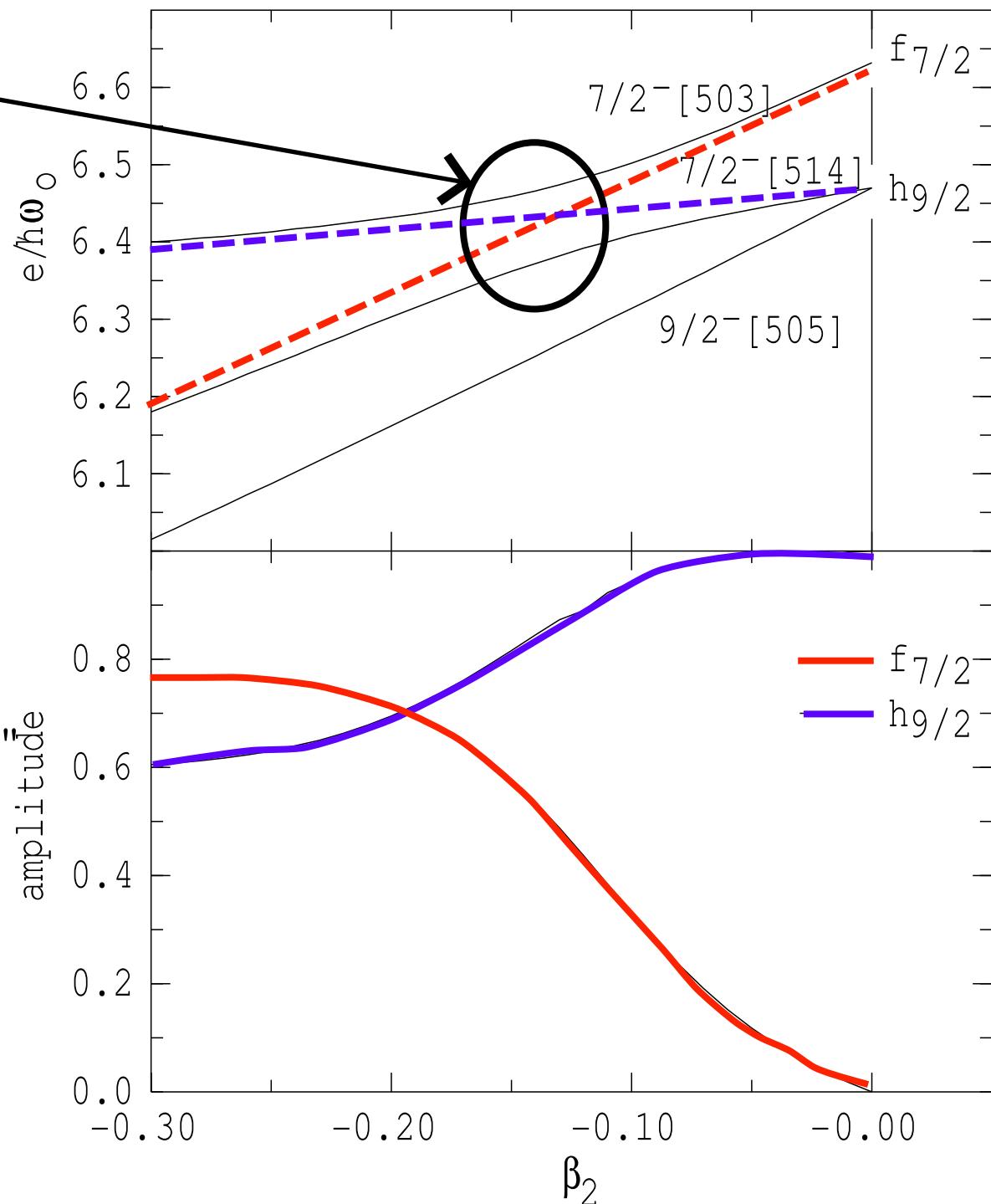
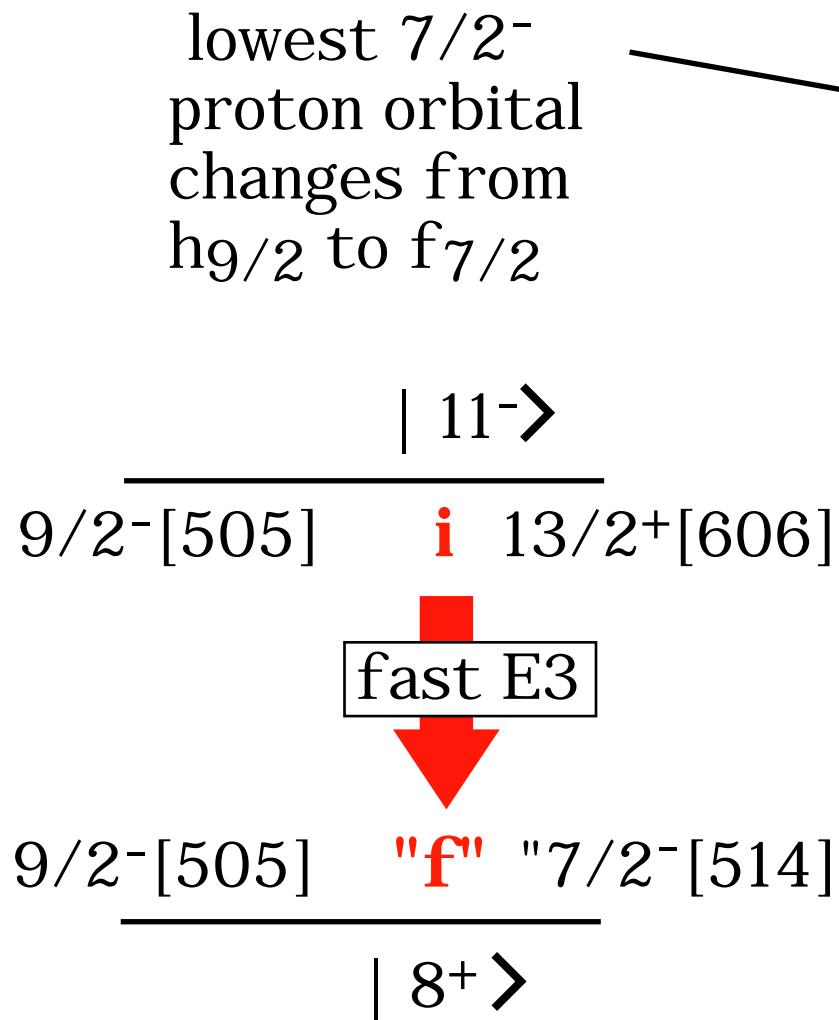
Type-A E3 ~ 20 W.u.

$K = 11^- \{$

$13/2+[606]$
 $9/2-[505]$

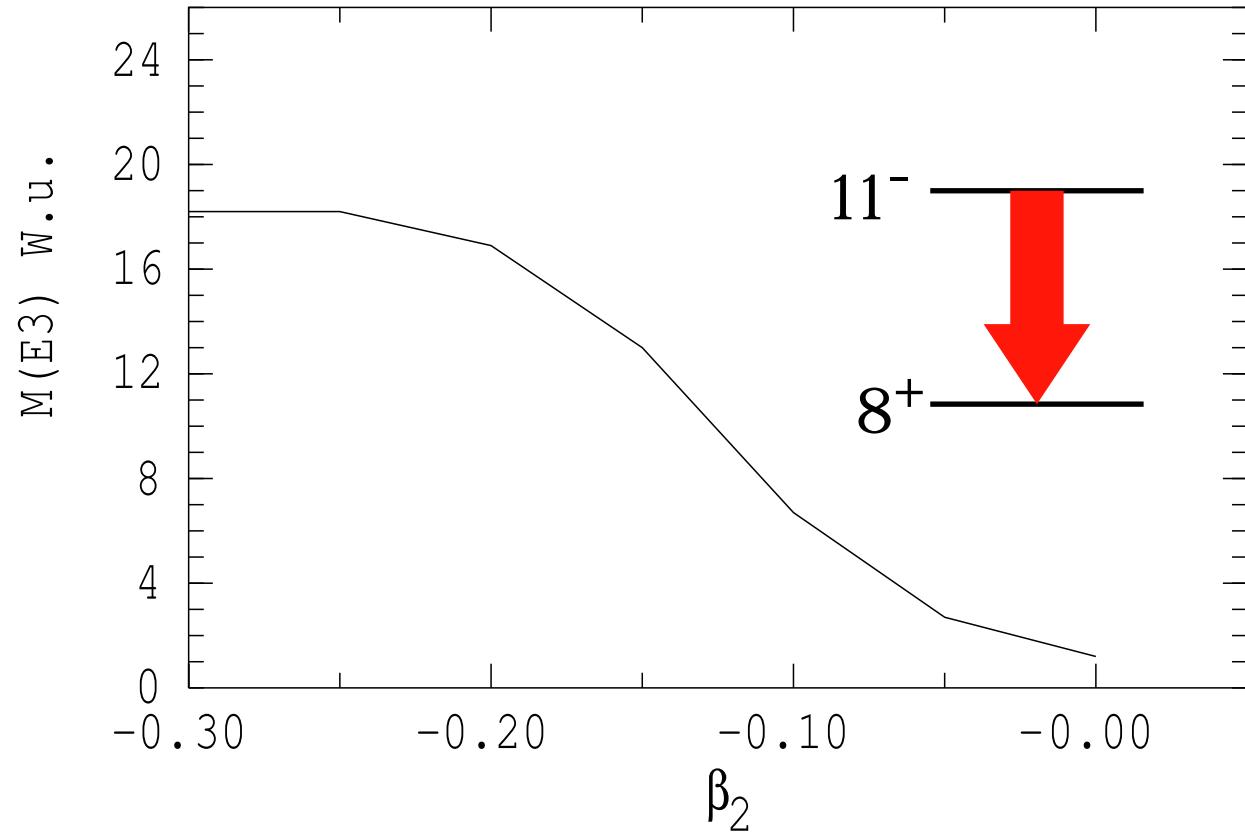


potential- mixing with oblate deformation



approximate dependence on oblate deformation

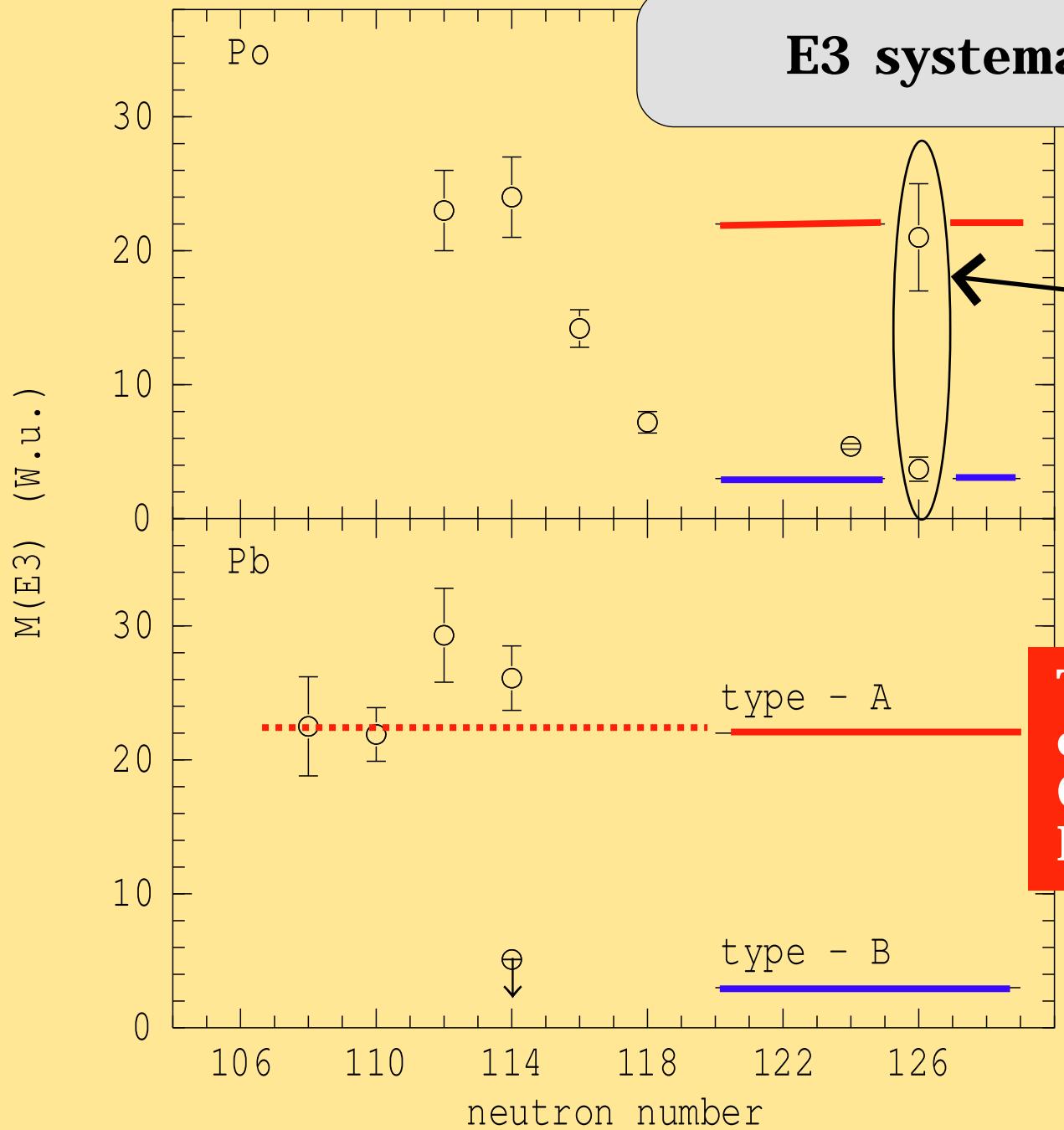
Taking calculated amplitudes α and β of $h9/2$ and $f7/2$ configuration in "7/2-[514]" orbital in typical Nilsson Potential



$$B(E3, 11^- \text{ to } 8^+) \sim \{\alpha U(\dots) \sqrt{B(E3; i_{13/2} \text{ to } h_{9/2})} + \beta U(\dots) \sqrt{B(E3; i_{13/2} \text{ to } f_{7/2})}\}^2$$

$$M(E3, 11^- \text{ to } 8^+) = \{\alpha 0.632 \sqrt{3} + \beta 1.0 \sqrt{22}\}^2 \text{ W.u.}$$

E3 systematics ; Pb- Po

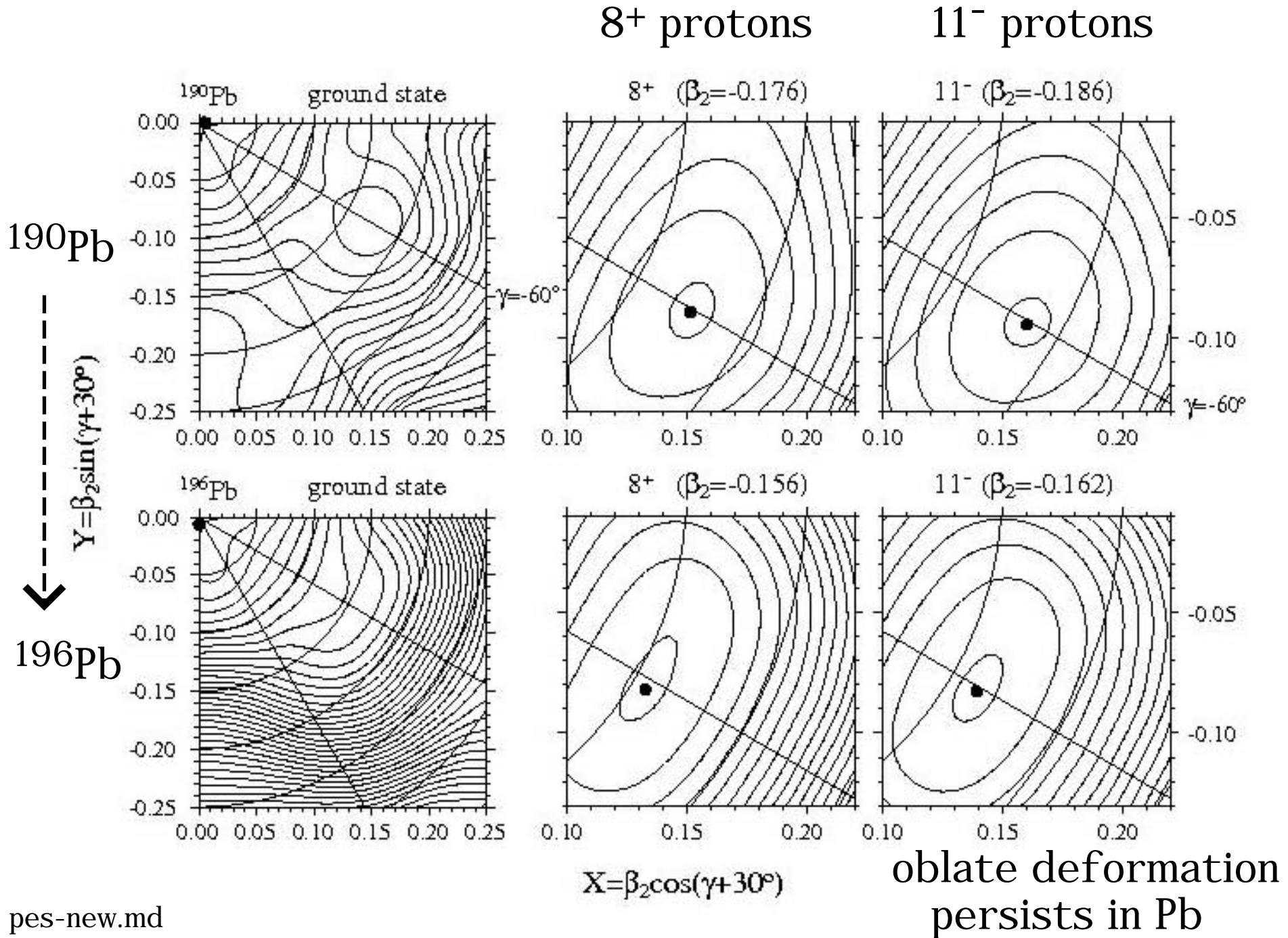


N=126 reproduced
by shell model with
octupole coupling
Stuchbery et al.
NPA 555 (1993)355

Type- A $E3 \sim 22$ W.u.
consistent with ~constant
Oblate deformation
N< 116 ?

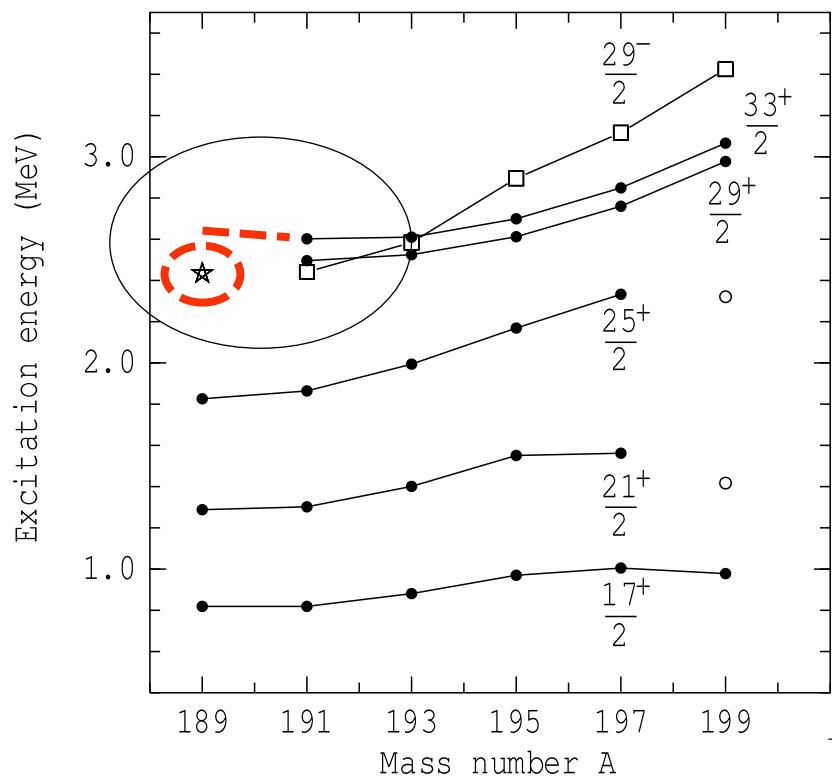
to be published...

PES - configuration constrained



CASE IV; ^{189}Pb isomer

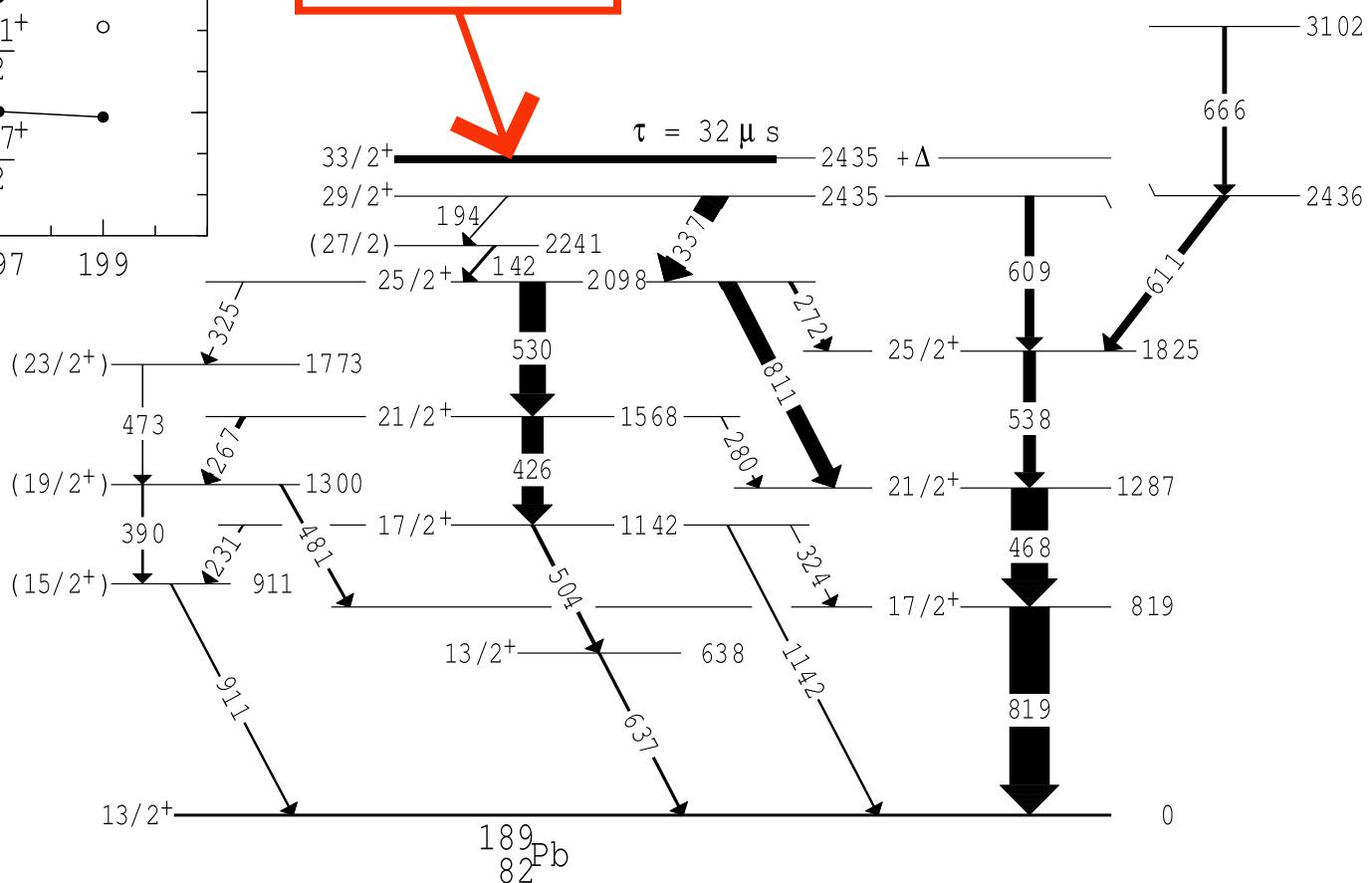
^{189}Pb : neutron $i_{13/2}^{-3}$



Baxter et al,
Phys Rev. C 71, 054302 (2005)

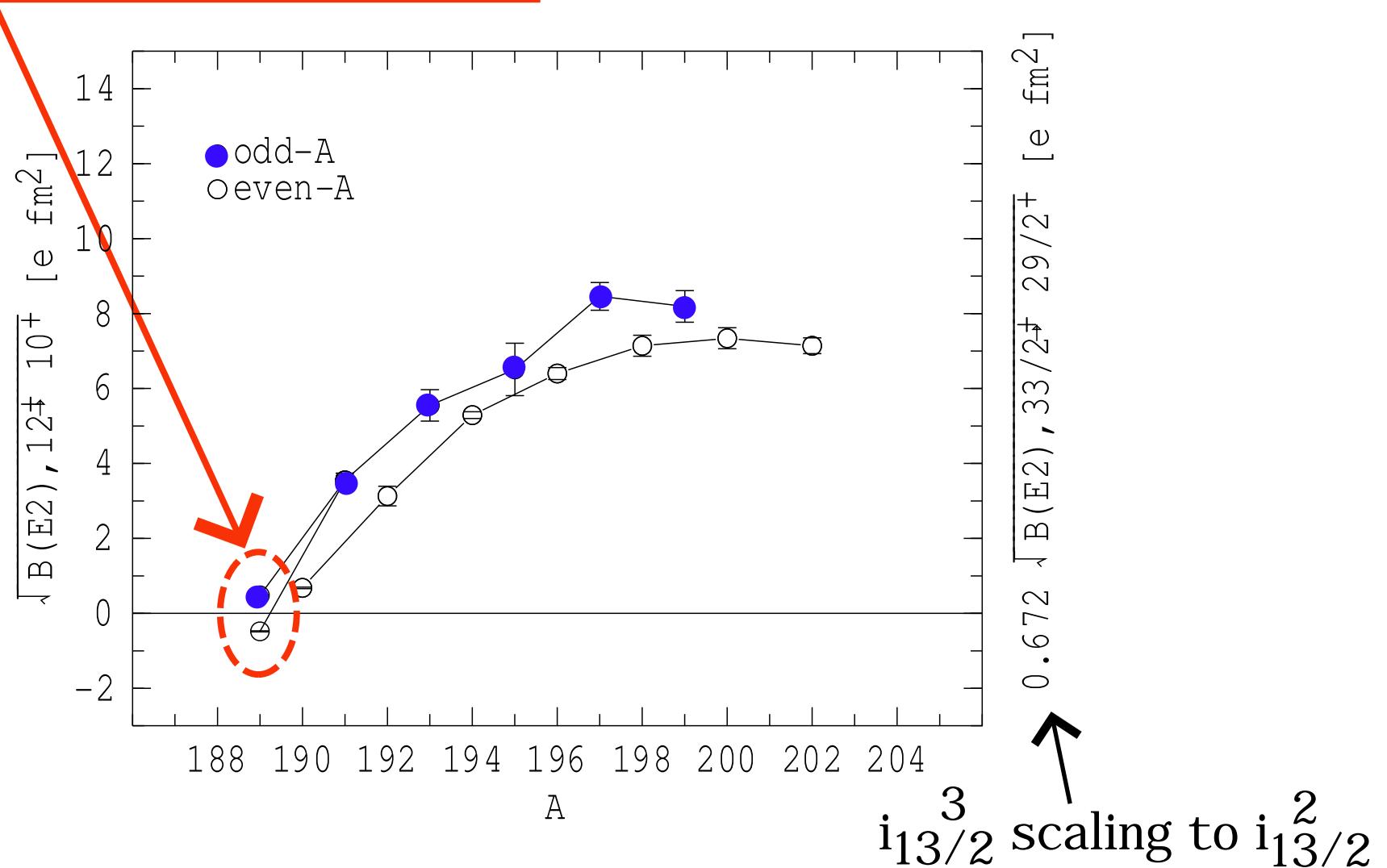
Tentative $J^\pi = 33/2^+$ assignment with unobserved E2 transition

$$\tau = 32 \mu\text{s}$$



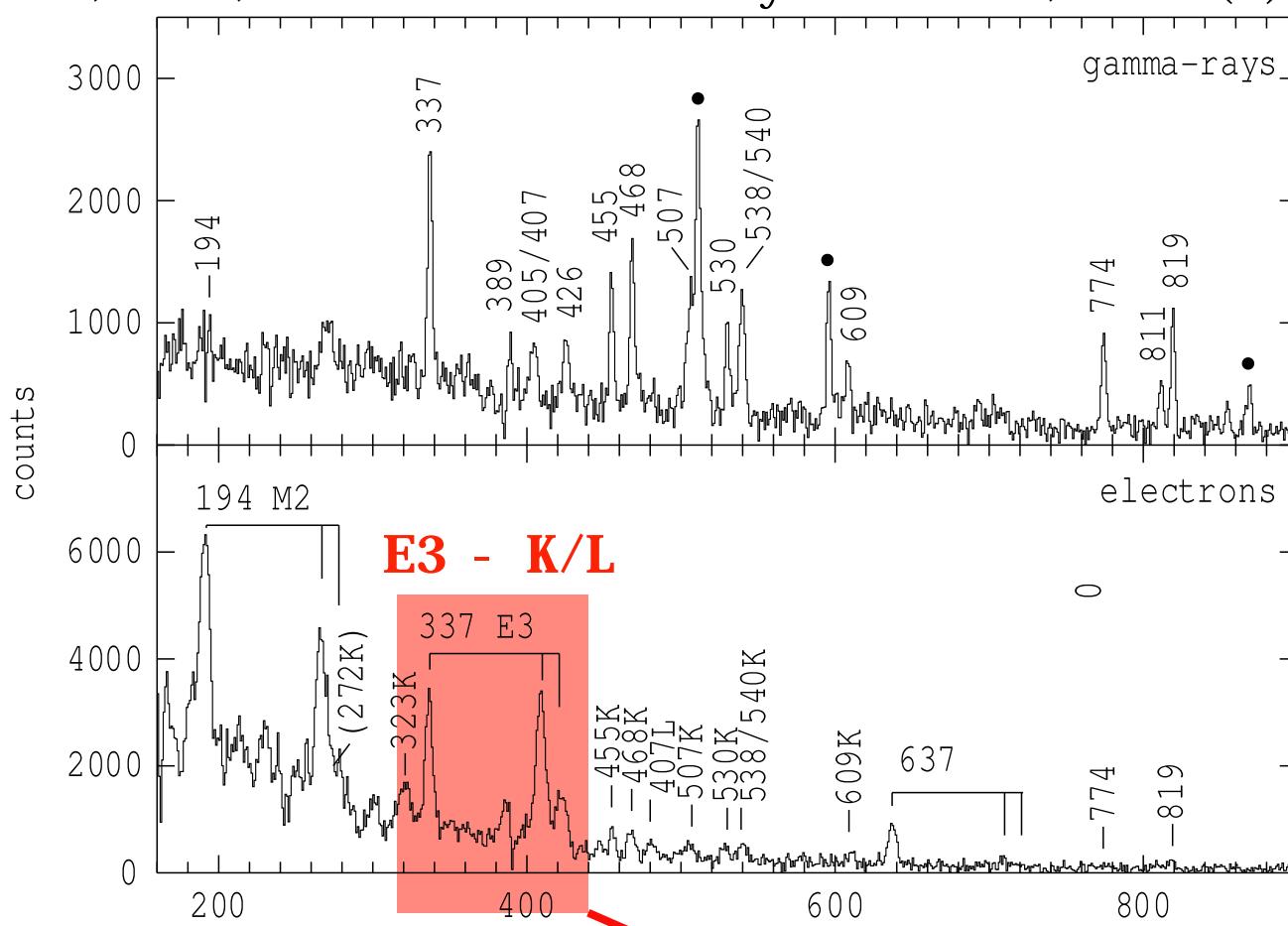
B(E2) mid- shell cancellation...

189Pb proposed 33/2⁺ isomer

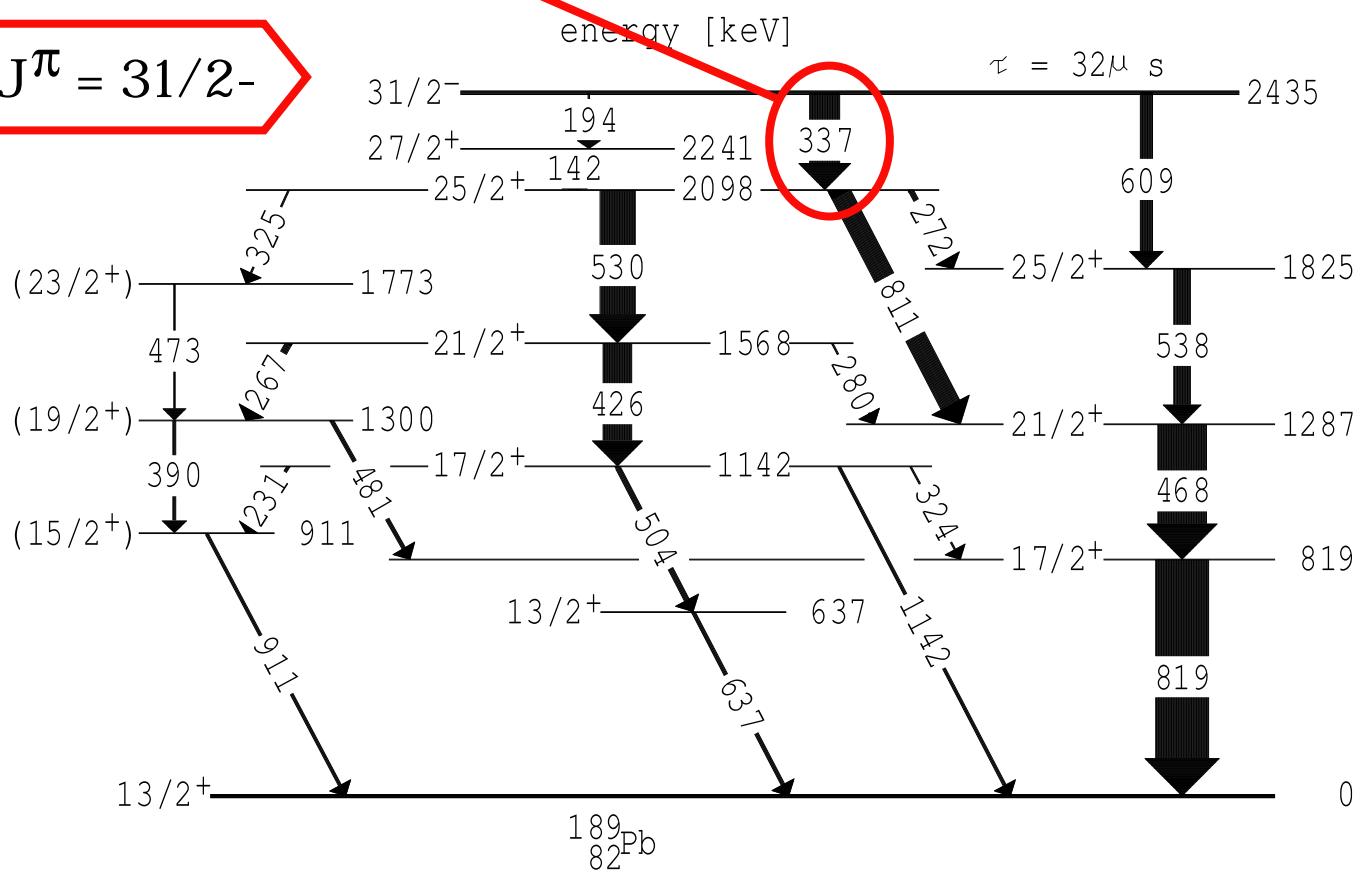


Solenogam test; ^{189}Pb 32 μs Isomer

Characterisation as a Shears-mode bandhead
 GDD, Lane, Kibedi, Nieminen. Phys Rev C 79, 01302(R) (2009).

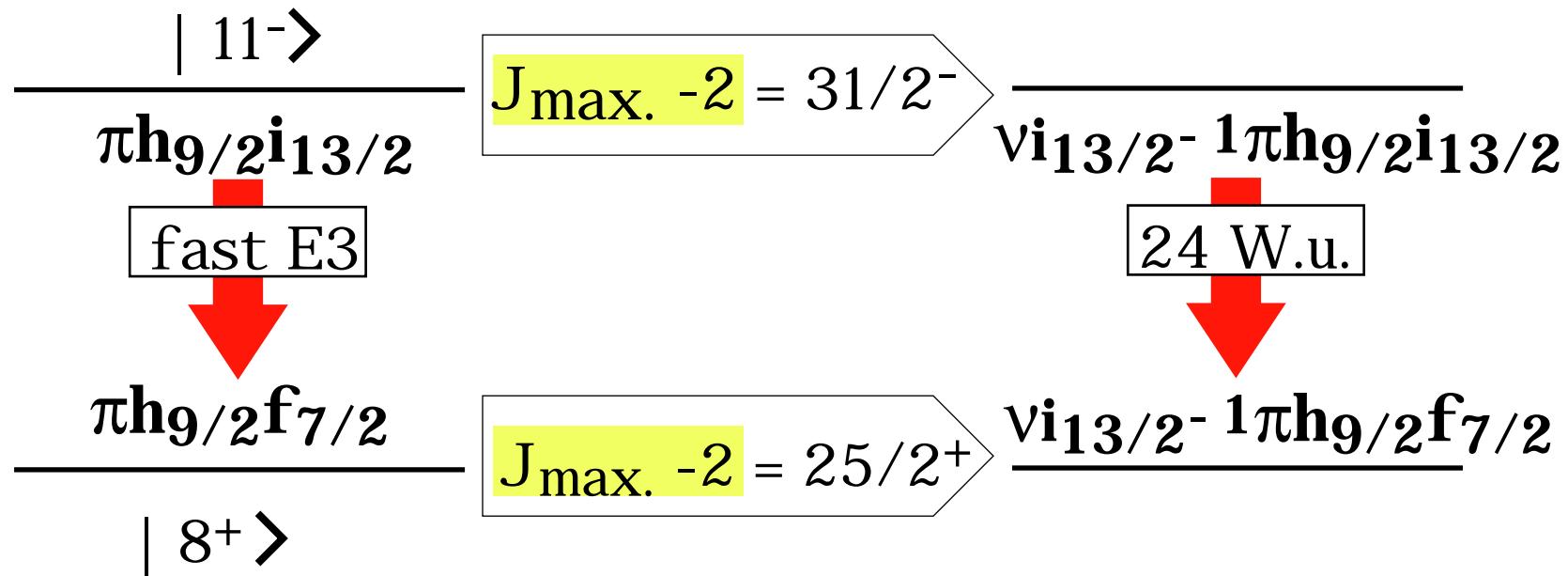


$J^\pi = 31/2^-$



Configuration

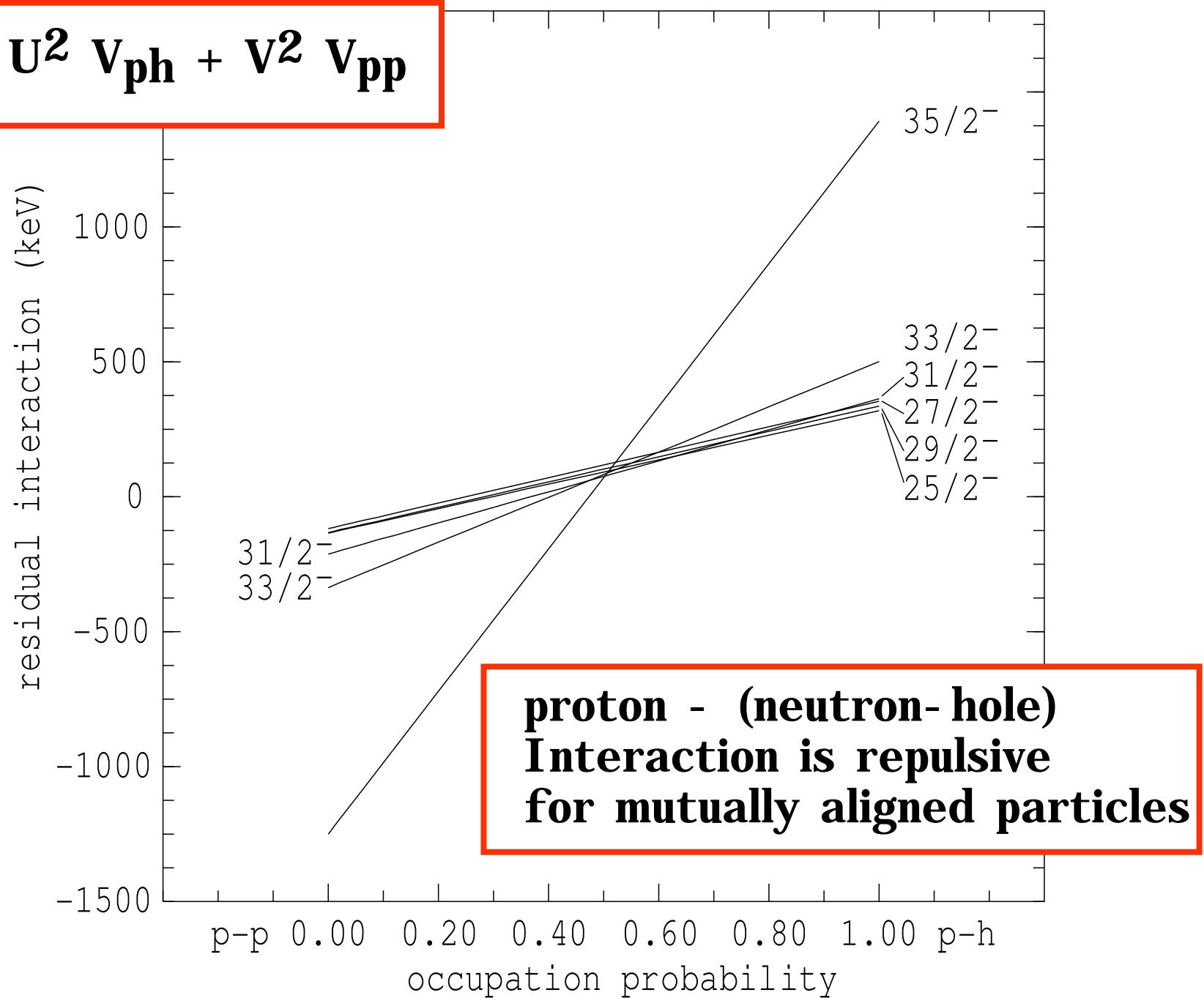
naive expectation:
Maximum coupling of
 11^- core and $i_{13/2}$ neutron hole $J^\pi = 35/2^-$



Non-Maximal Coupling -
balance of residual interactions

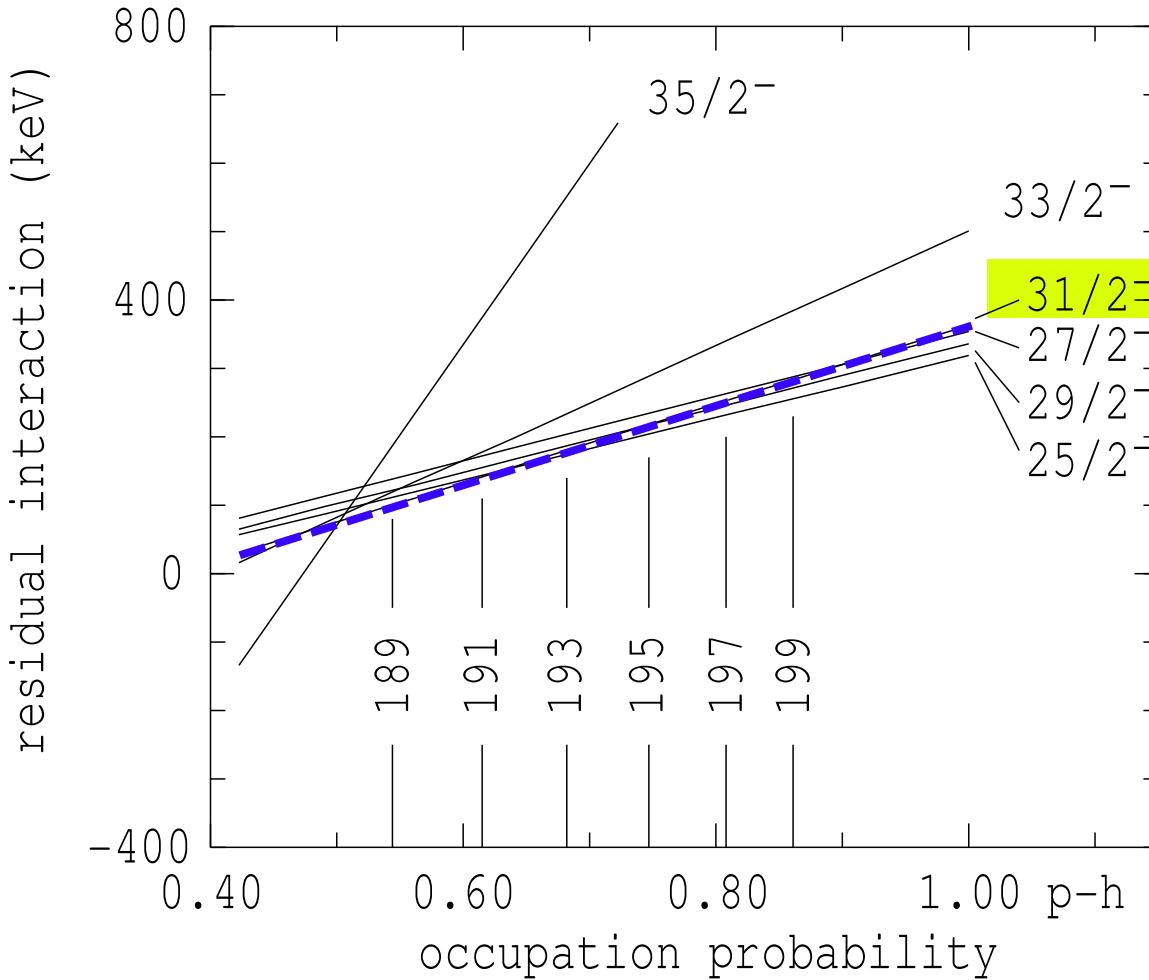
Residual Interactions: Particles/Holes

$$V_{\text{eff.}} = U^2 V_{\text{ph}} + V^2 V_{\text{pp}}$$

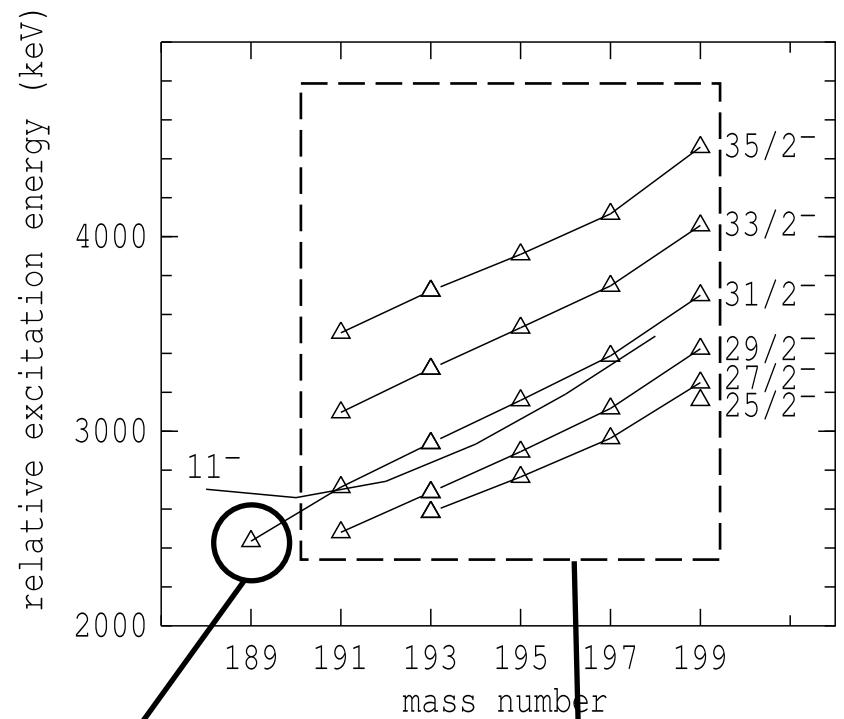
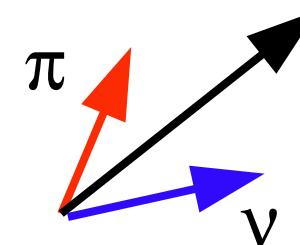


Residual Interactions: Detail/shears

$\pi [h_{9/2}i_{13/2}]_{11^-} \otimes \nu i_{13/2}$



31/2⁻ Isomer



shears bands...

Final Comments...

**Isomers are a useful and important tool
in Nuclear Structure**

The detail matters

Expectation is not the same as Measurement

A Plea:

Don't ignore what has gone before

a thought for the day

Isomers have a life of their own

Dafni et al. PRL 55 (1985) 1271
 Quadrupole Moment.. 63/2⁻ isomer in ²¹¹Rn
 "Absence of Core Deformations at Very High Spins"

$$\Sigma_l Q_{sp}$$

$^{211}\text{Rn} (63/2^-) = ^{212}\text{Rn} (30^+) + \nu (3/2^-)$
Spherical (eff. charges)
 $Q(\text{shell}) = -186$

Single Particle
Moments (e.fm²)

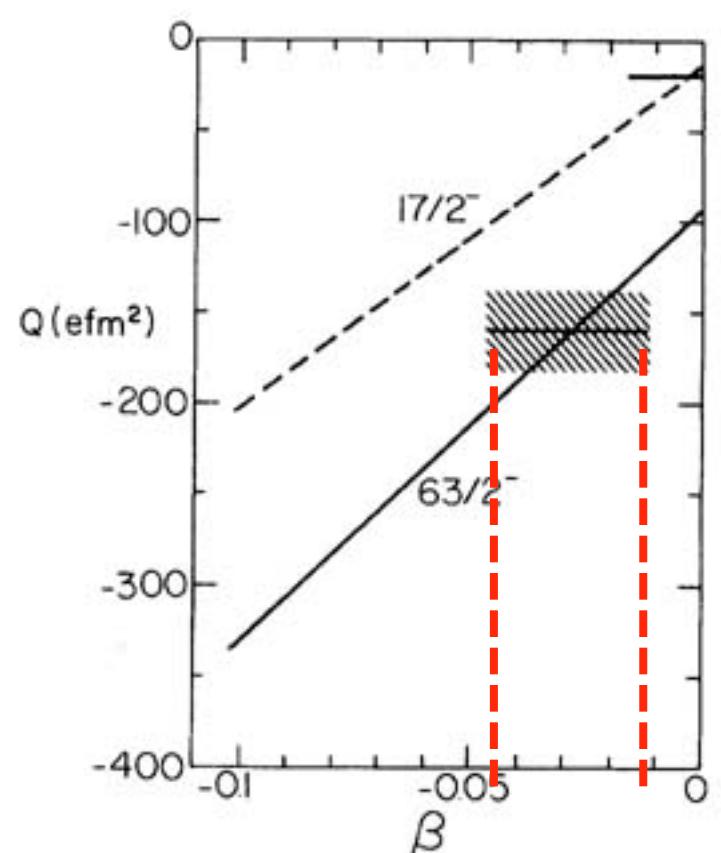
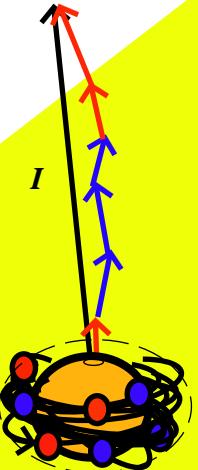
$\pi h_{9/2}$	-43
$\pi l_{13/2}$	-62
$v f_{5/2}$	+20
$v g_{9/2}$	-29
$v i_{11/2}$	-28
$v j_{15/2}$	-42

$$|Q_{\text{exp}}| = 160(22)$$

X

Deformed (no eff. charges)
 $Q = f(I)Q_0 = f(I)(Q_I + Q_{\text{core}})$
 $Q_{\text{core}} \propto Z_{\text{core}} R^2 (\beta + 0.36\beta^2)$

Also, see Sagawa&Arima PLB 202 (1988) 15



$$\beta \sim 0.027(10)$$