

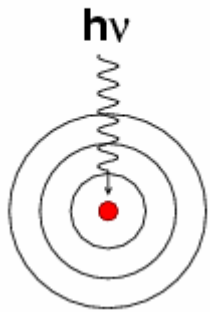


Nuclear Excitation in Plasmas

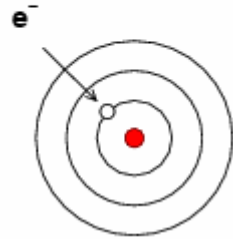
G. Gosselin, P. Morel, V. Méot
(CEA, France)



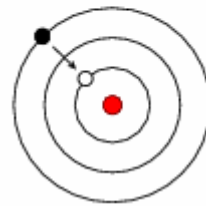
Nuclear Excitation Processes in Plasmas



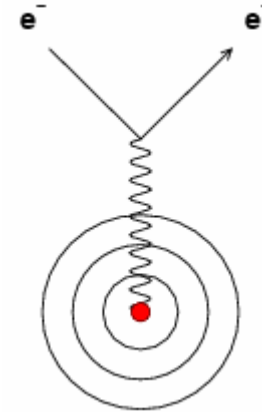
Photon
Absorption



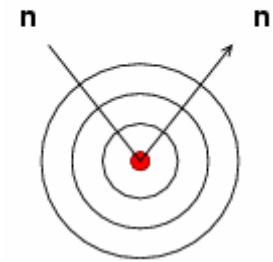
NEEC



NEET



Electron
Inelastic
Scattering



Neutron
Inelastic
Scattering



Thermodynamic Equilibrium

- Ions, electrons and photons are at thermodynamic equilibrium

$$T_i = T_e = T_r = T$$

- Nuclear populations are not $\propto$ Boltzmann law does not apply

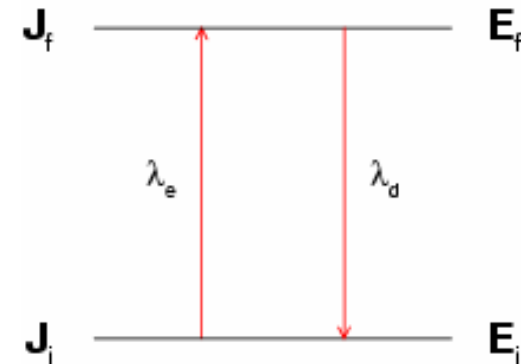
$$N_i \neq \frac{(2J_i + 1)e^{-\frac{E_i}{kT}}}{\sum_{j=1}^n (2J_j + 1)e^{-\frac{E_j}{kT}}}$$

- Equilibration time (under stationary conditions)

$$\tau = \frac{\ln 2}{\lambda_e + \lambda_d}$$

- Populations at equilibrium

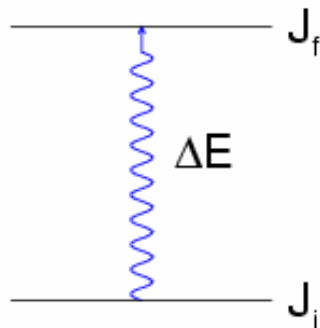
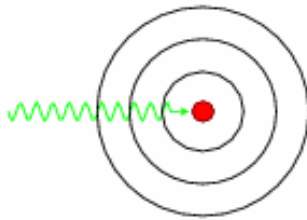
$$\frac{N_f(\infty)}{N_i(\infty)} = \frac{\lambda_e}{\lambda_d}$$



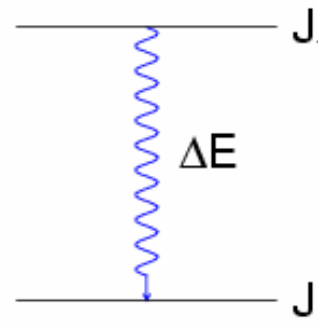
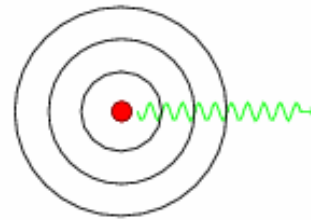


Radiative excitation and de-excitation

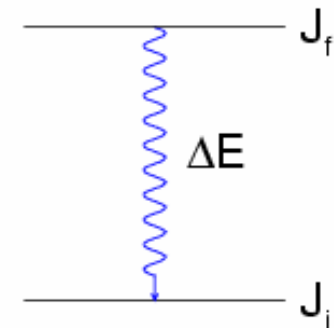
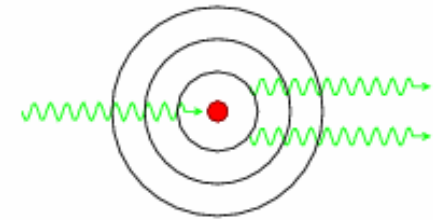
Absorption

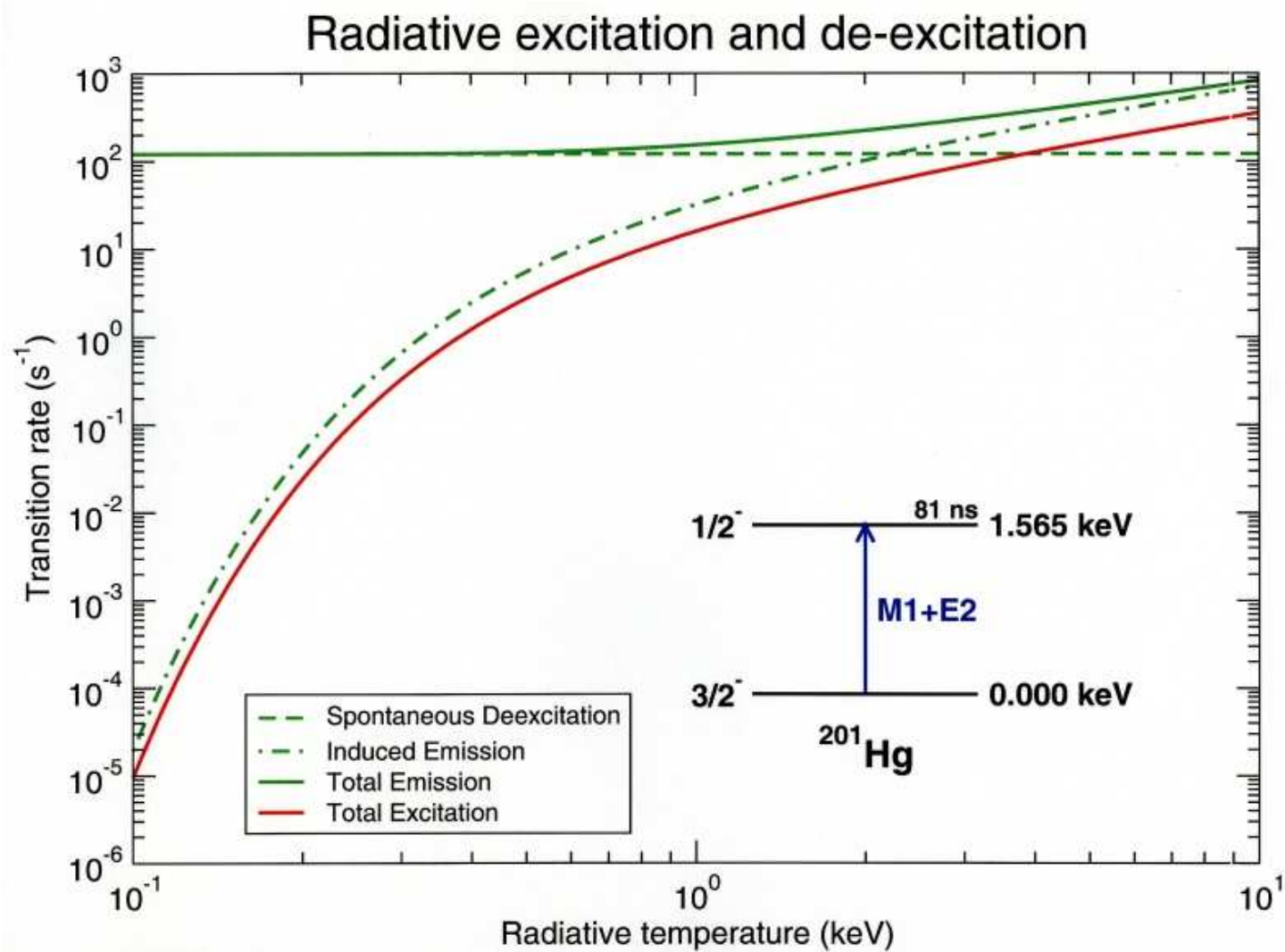


Spontaneous Emission



Induced Emission





G. Gosselin, P. Morel, Phys. Rev. **C70**, 064603 (2004)

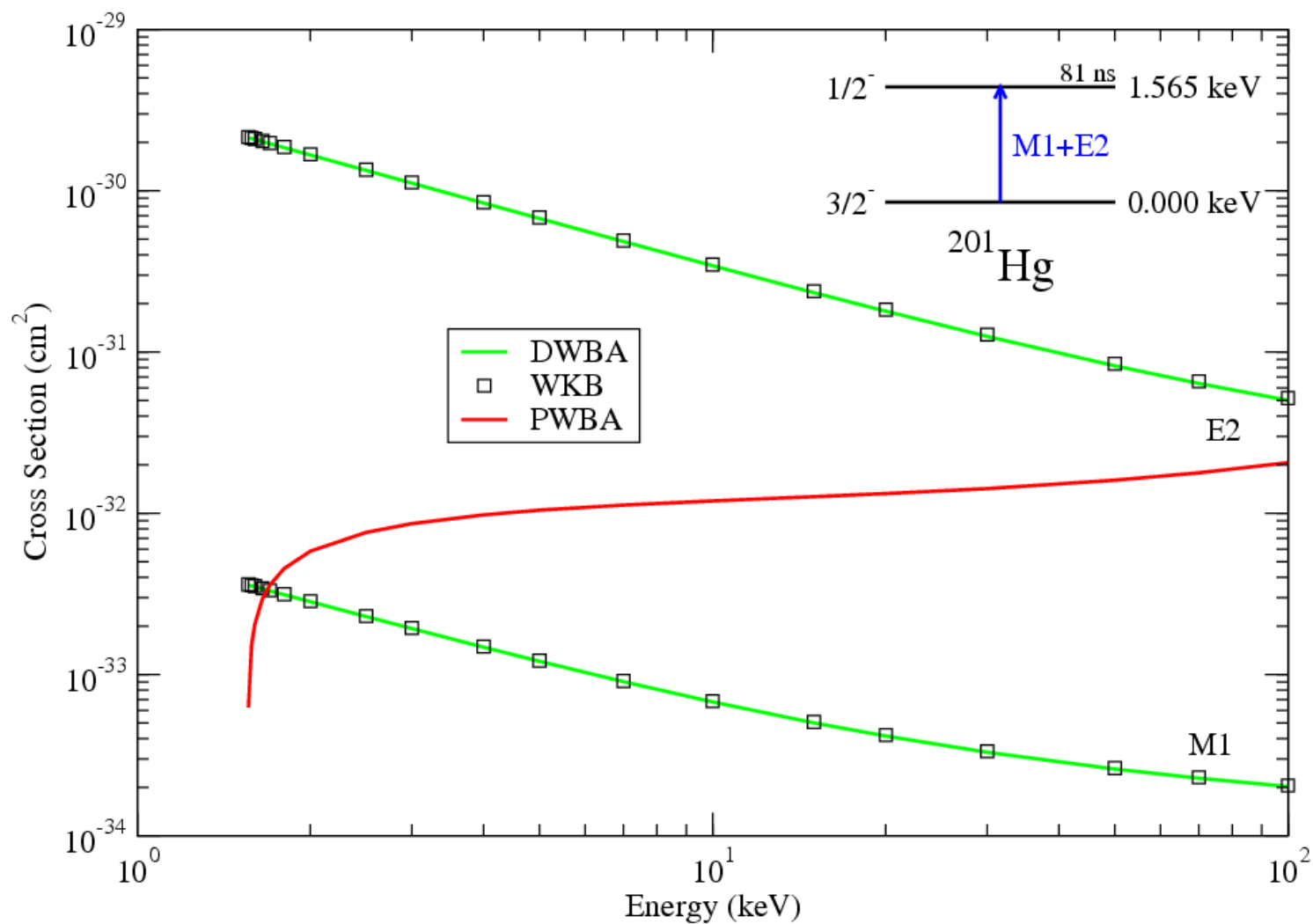


Quantum Coulomb Excitation

- PWBA Approximation
 - Plane wave : works well at high energy
- DWBA Approximation
 - Distorted wave function : works well over the whole energy domain
 - Radial equation can be solved by two methods
 - Exact solution (only with an unscreened potential)
 - WKB method
- Radial exact solution
 - Hypergeometric functions
 - Numerical evaluation very difficult for low energies
- WKB Method
 - Approximate radial equation resolution
 - Solution by Langer
 - Asymptotic solutions Bessel functions
 - No divergence at turning point

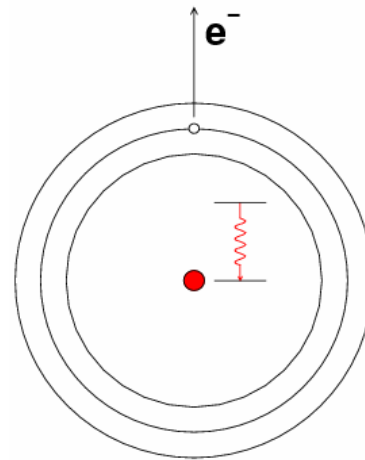
$$\frac{d^2\varphi_\ell}{dr^2} + \left[\frac{2m}{\hbar^2} E - \frac{2m}{\hbar^2} V(r) - \frac{\ell(\ell+1)}{r^2} \right] \varphi_\ell = 0$$

Inelastic Electron Scattering

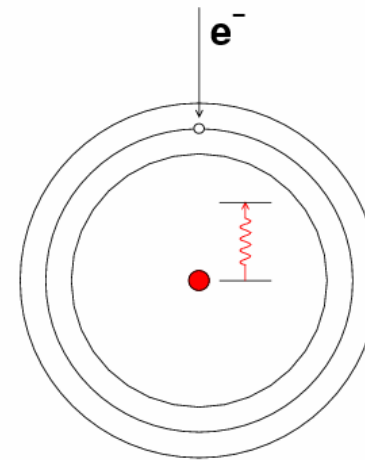


G. Gosselin, N. Pillet, V. Méot, P. Morel, A. Dzyublik, Phys. Rev. C79, 014604 (2009)

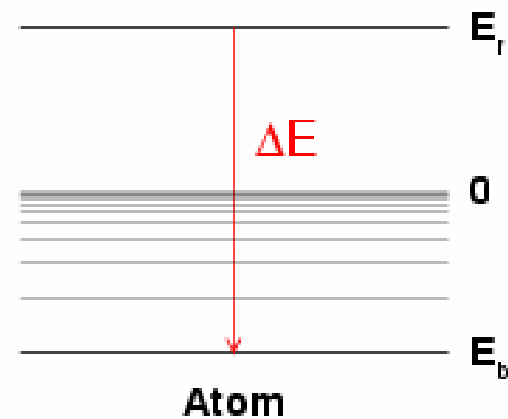
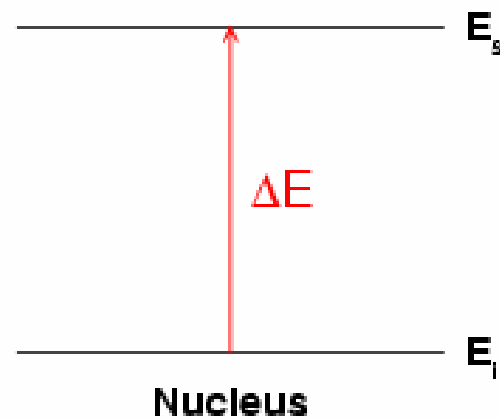
Internal Conversion / Electronic Capture



Internal Conversion



NEEC





Internal conversion / NEEC

- Internal conversion coefficient

$$\alpha = \frac{\lambda_{J_f \rightarrow J_i}^{\text{IC}}}{\lambda_{J_f \rightarrow J_i}^{\gamma}} = \frac{T_{J_f \rightarrow J_i}^{\gamma}}{T_{J_f \rightarrow J_i}^{\text{IC}}}$$

- Internal conversion rate in laboratory

$$\lambda_{\text{IC}}^{\text{nuc}} = \frac{\text{Log } 2}{T_{J_f \rightarrow J_i}^{\text{IC}}} = \frac{2\pi}{\hbar} \left| \langle \psi_i \phi_r | H | \psi_f \phi_b \rangle \right|^2 \rho_r(E_r)$$

- NEEC cross section in laboratory

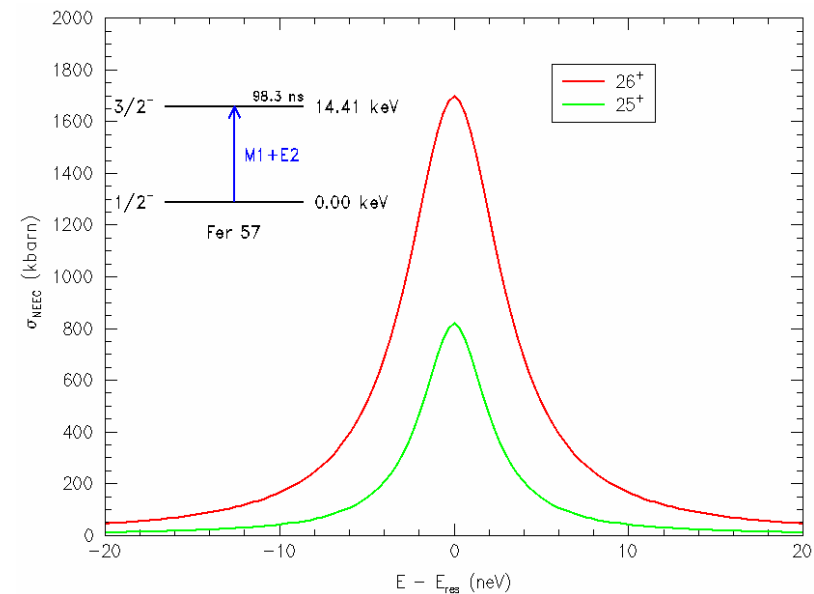
$$\sigma_{\text{NEEC}}(E) = \frac{2\pi}{\hbar v_e} \left| \langle \psi_f \phi_b | H | \psi_i \phi_r \rangle \right|^2 \rho_b(E)$$

NEEC cross section

$$\sigma_{\text{NEEC}}(E) = \frac{\pi \hbar^2}{2m_e E} \frac{2J_f + 1}{2J_i + 1} \frac{\hbar \alpha \text{Log} 2}{T_{J_f \rightarrow J_i}^\gamma} \frac{\Gamma}{(E - E_r)^2 + \left(\frac{\Gamma}{2}\right)^2}$$

$$\sigma_{\text{NEEC}} = \frac{\pi^2 \hbar^2}{m_e E} \frac{2J_f + 1}{2J_i + 1} \frac{\hbar \alpha \text{Log} 2}{T_{J_f \rightarrow J_i}^\gamma} \delta(E - E_r)$$

- Laboratory conditions
 - Non excited atom
 - Natural level widths
 - Internal conversion coefficient α





Relativistic Average Atom

- Thomas-Fermi-Dirac model
 - The atom lies in a spherical box whose radius is dictated by density
 - Average atom : mean value of binding energies and occupation numbers
 - Relativistic : necessary for heavy nuclei internal conversion calculations

B. F. Rozsnyai, Phys. Rev. A5 (1972) 1137

Macroscopic transition rates

- Excitation rate : NEEC

$$\lambda_e = \int_0^{\Delta E} \sigma_{\text{NEEC}}(E) v_e f(E) P(E - \Delta E) dE$$

$$\lambda_e = \frac{2J_f + 1}{2J_i + 1} \frac{\alpha(T_e) \text{Log } 2}{T_{J_f \rightarrow J_i}^\gamma} g(E_b) f_{\text{FD}}(E_r) [1 - f_{\text{FD}}(E_b)] \left[\text{erf}\left(\frac{E_r}{\sigma\sqrt{2}}\right) - \text{erf}\left(\frac{E_b}{\sigma\sqrt{2}}\right) \right]$$

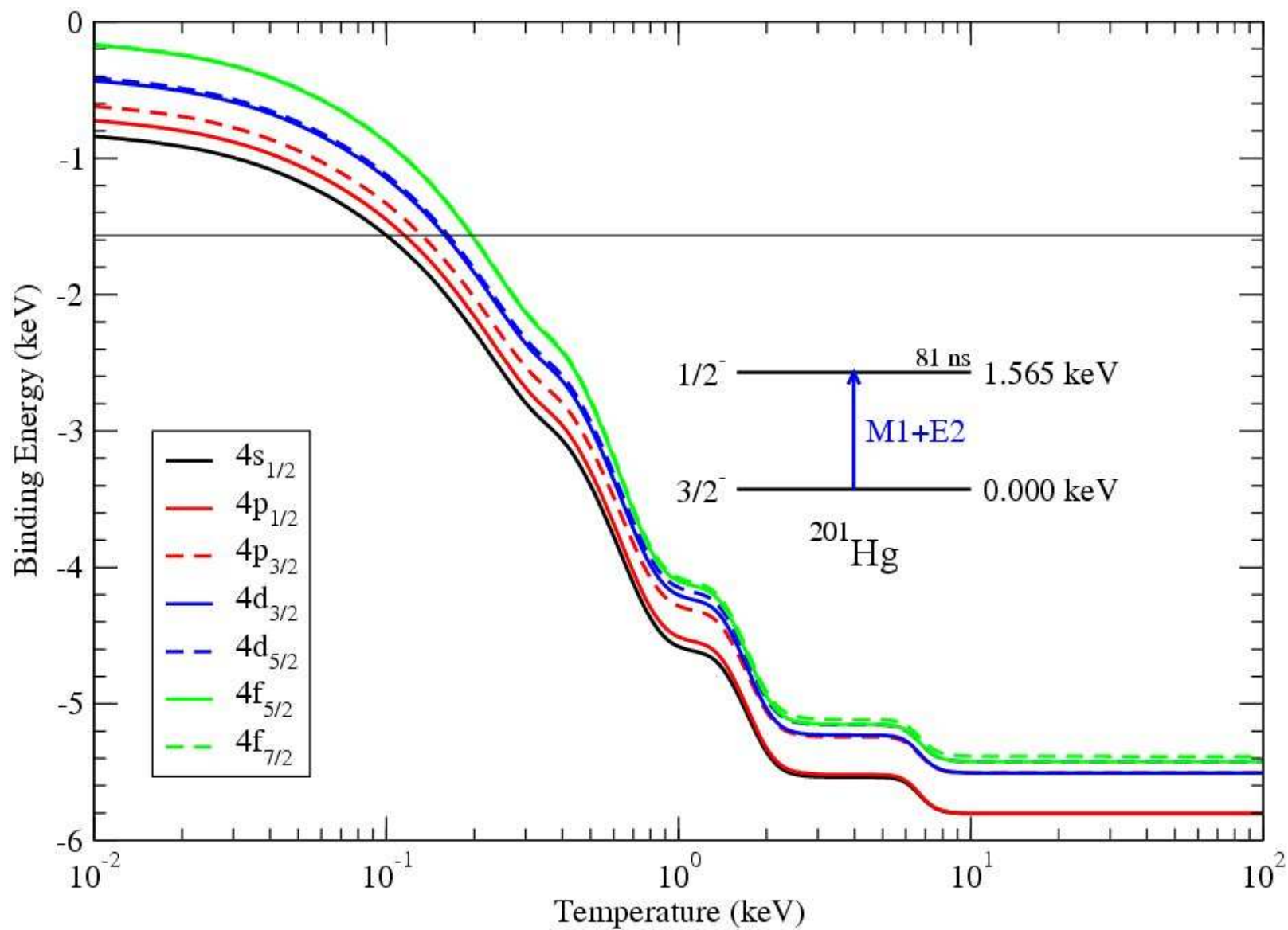
- De-excitation rate : **Calculated with Average Atom wavefunctions**

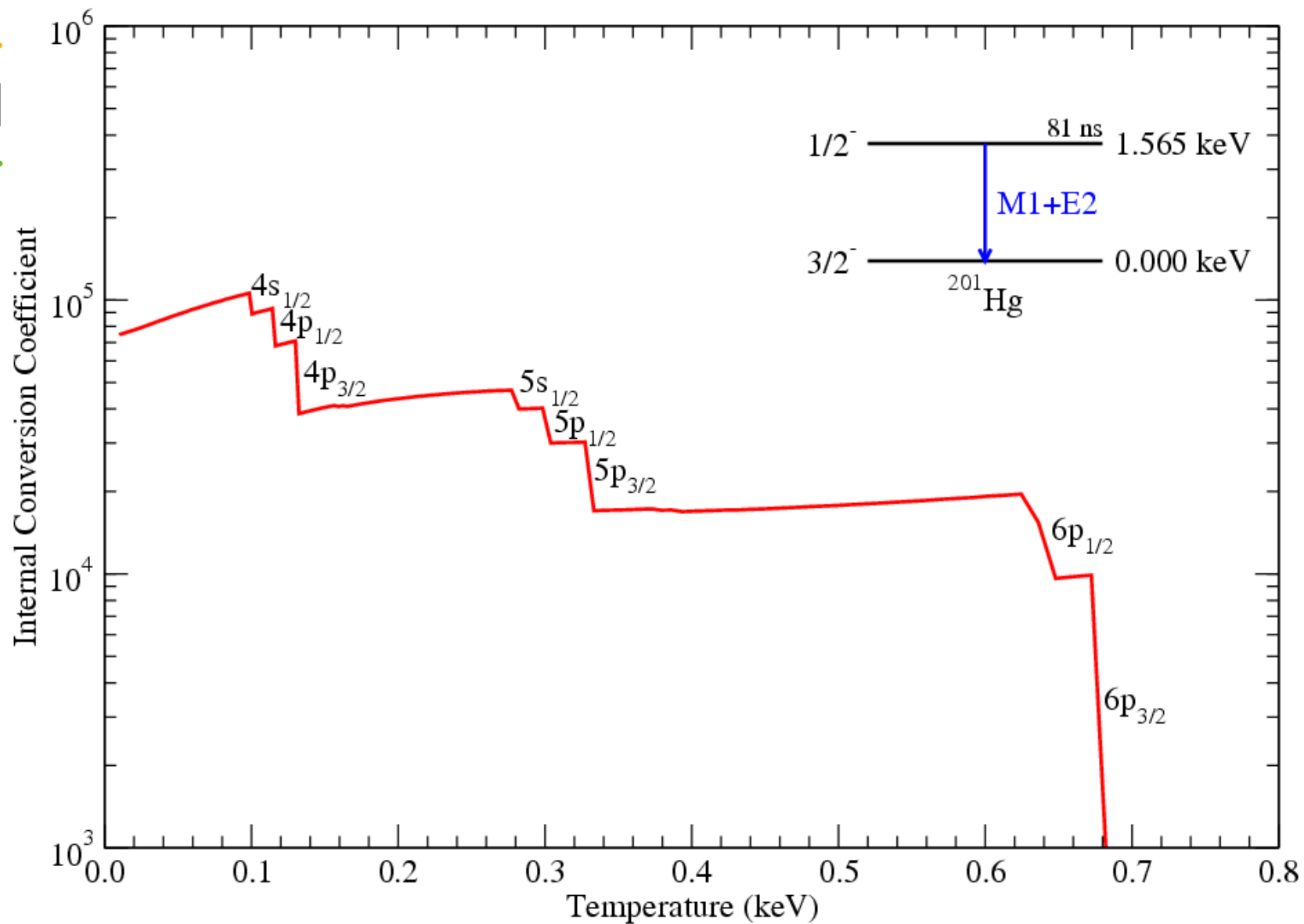
$$\lambda_d = \int_{-\Delta E}^0 \lambda_{\text{IC}}(E) P(E + \Delta E) dE$$

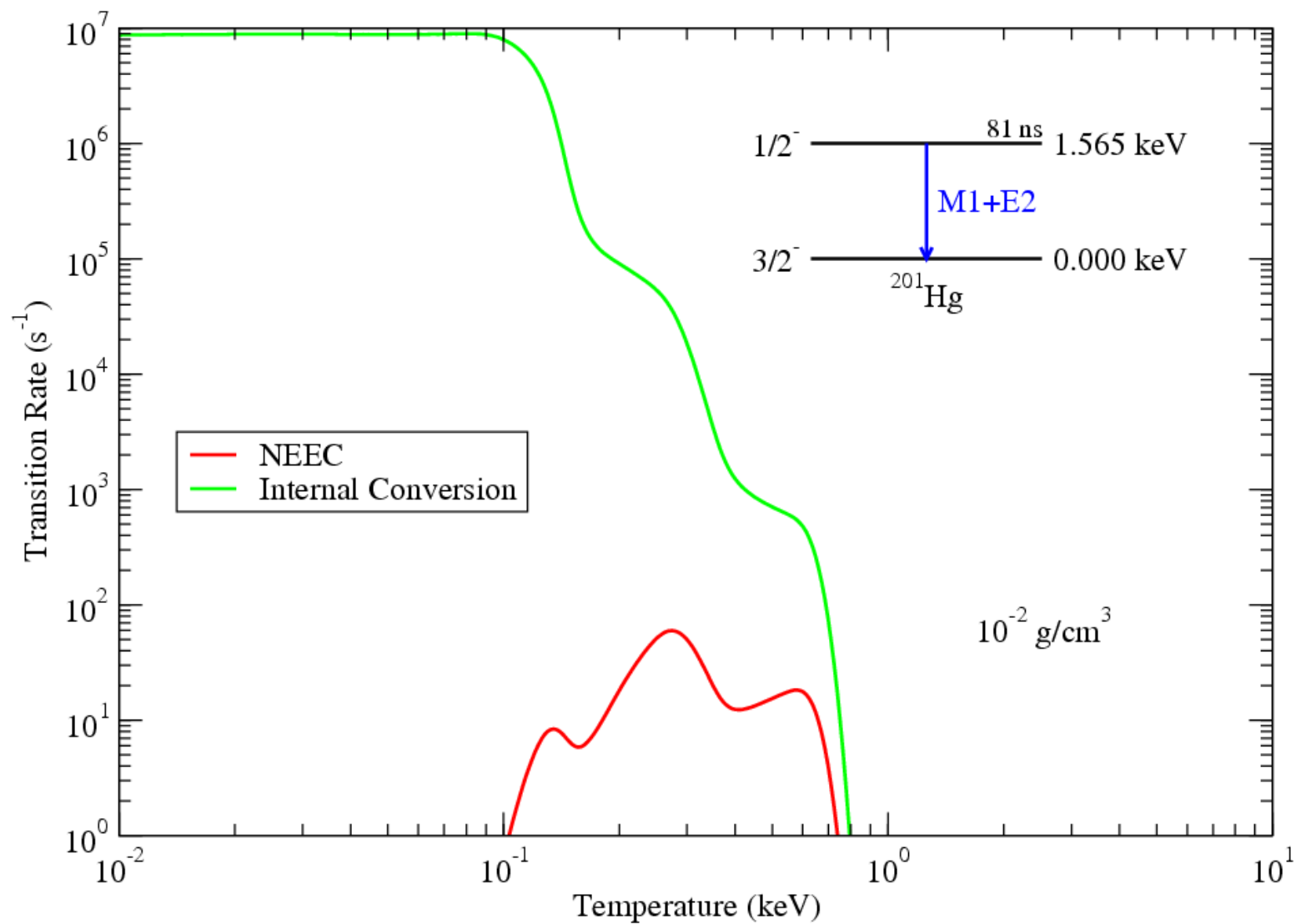
$$\lambda_d = \frac{\alpha(T_e) \text{Log } 2}{T_{J_f \rightarrow J_i}^\gamma} g(E_b) f_{\text{FD}}(E_b) [1 - f_{\text{FD}}(E_r)] \left[\text{erf}\left(\frac{E_r}{\sigma\sqrt{2}}\right) - \text{erf}\left(\frac{E_b}{\sigma\sqrt{2}}\right) \right]$$

Principle of detailed balance fulfilled (Boltzmann)

G. D. Doolen, Phys. Rev. C18 (1978) 2547

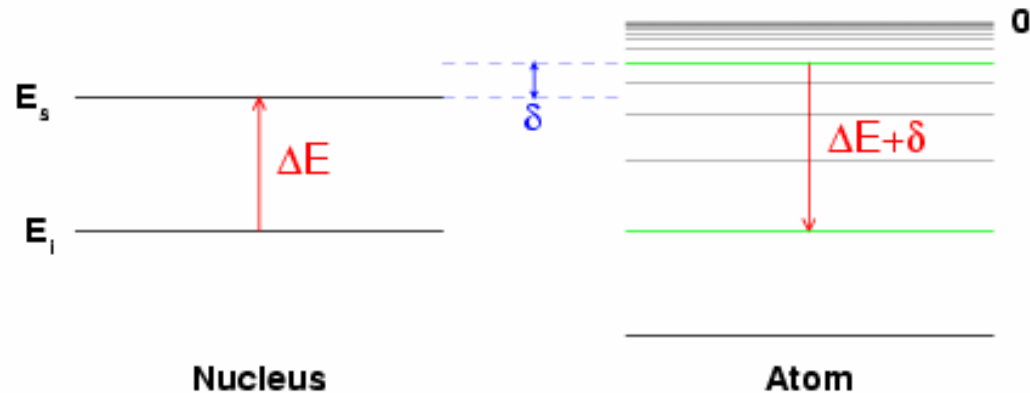
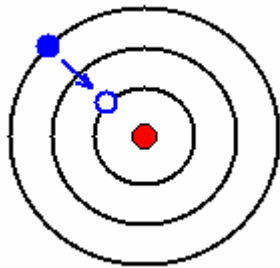






G. Gosselin and P. Morel, Phys. Rev. **C70** (2004) 064603

NEET in plasma



- Resonant phenomenon
 - Energy mismatch : δ
 - Quantum selection rules (spin, parity)

M. Harston, J. F. Chemin, Phys. Rev. **C59** (1999) 2462



Coupling Matrix Element

$$|R_{12}(\delta)|^2 = 4\pi\alpha\omega_N^2 \left(\omega_N + \cancel{\frac{\delta}{2}} \right)^{2L} \frac{2J_2 + 1}{L^2 [(2L + 1)!!]^2} \begin{pmatrix} J_1 & J_2 & L \\ \frac{1}{2} & -\frac{1}{2} & 0 \end{pmatrix}^2 |R_{n_1\kappa_1 n_2\kappa_2}|^2 B(EL)$$

Nuclear Transition Energy

Atomic Matrix Element

Nuclear Matrix Element

Expression nearly identical to Internal Conversion matrix element

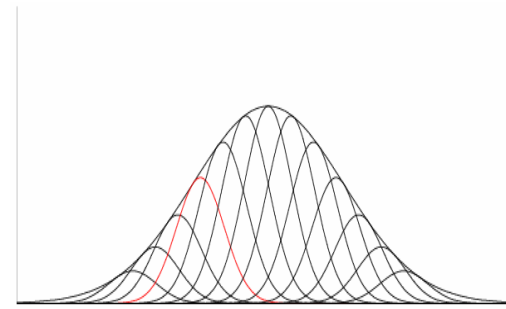
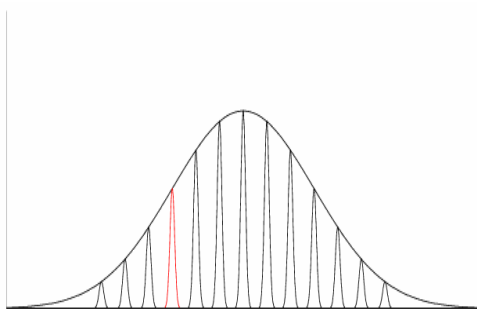
NEET Transition Rate

- Summation over every configuration

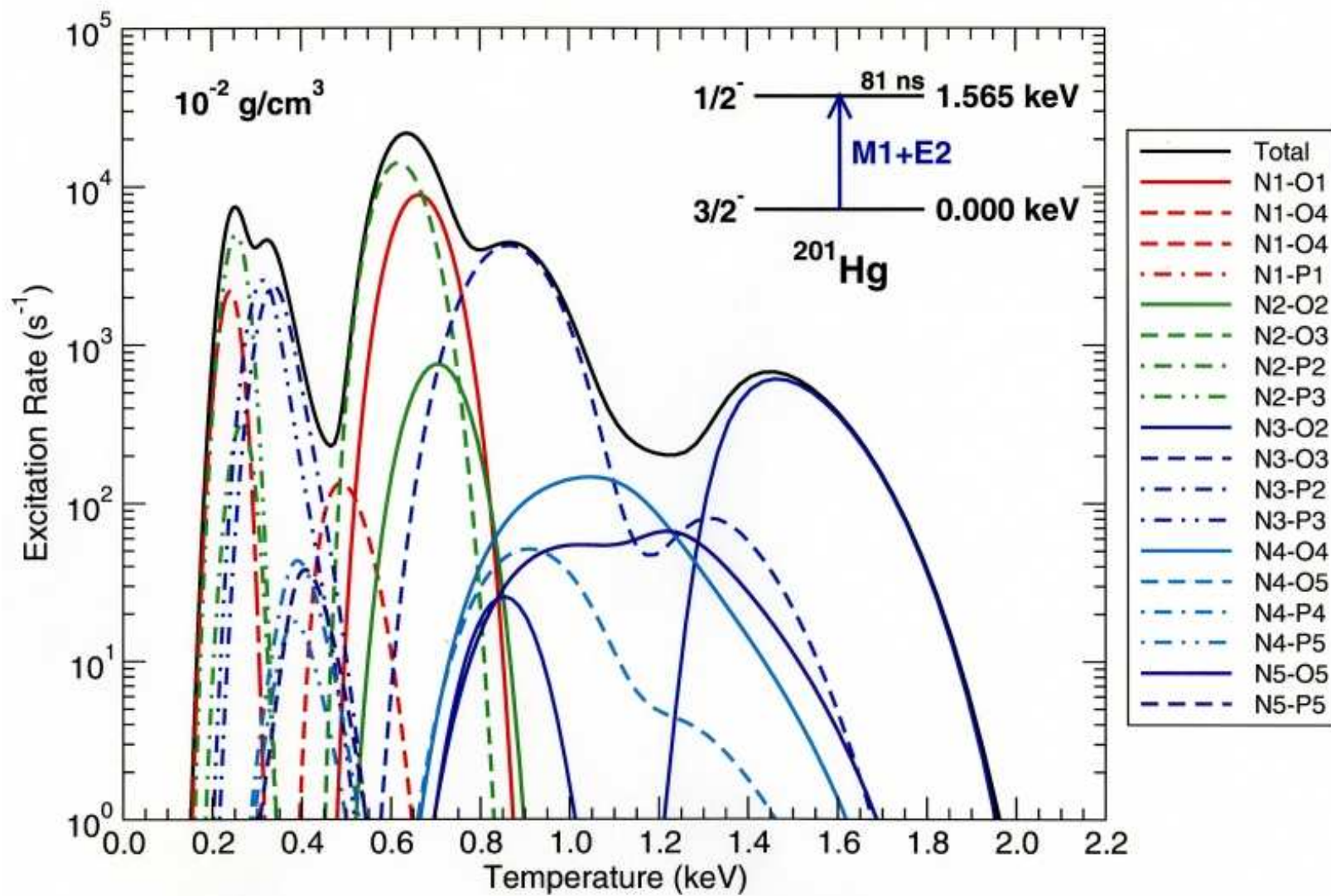
$$\lambda_{\text{NEET}}(\rho, T_e) = \sum_{\alpha, \beta} P_{\alpha}(\rho, T_e) \lambda_{\alpha} P^{\infty}(\delta_{\alpha\beta})$$

- Statistical envelope

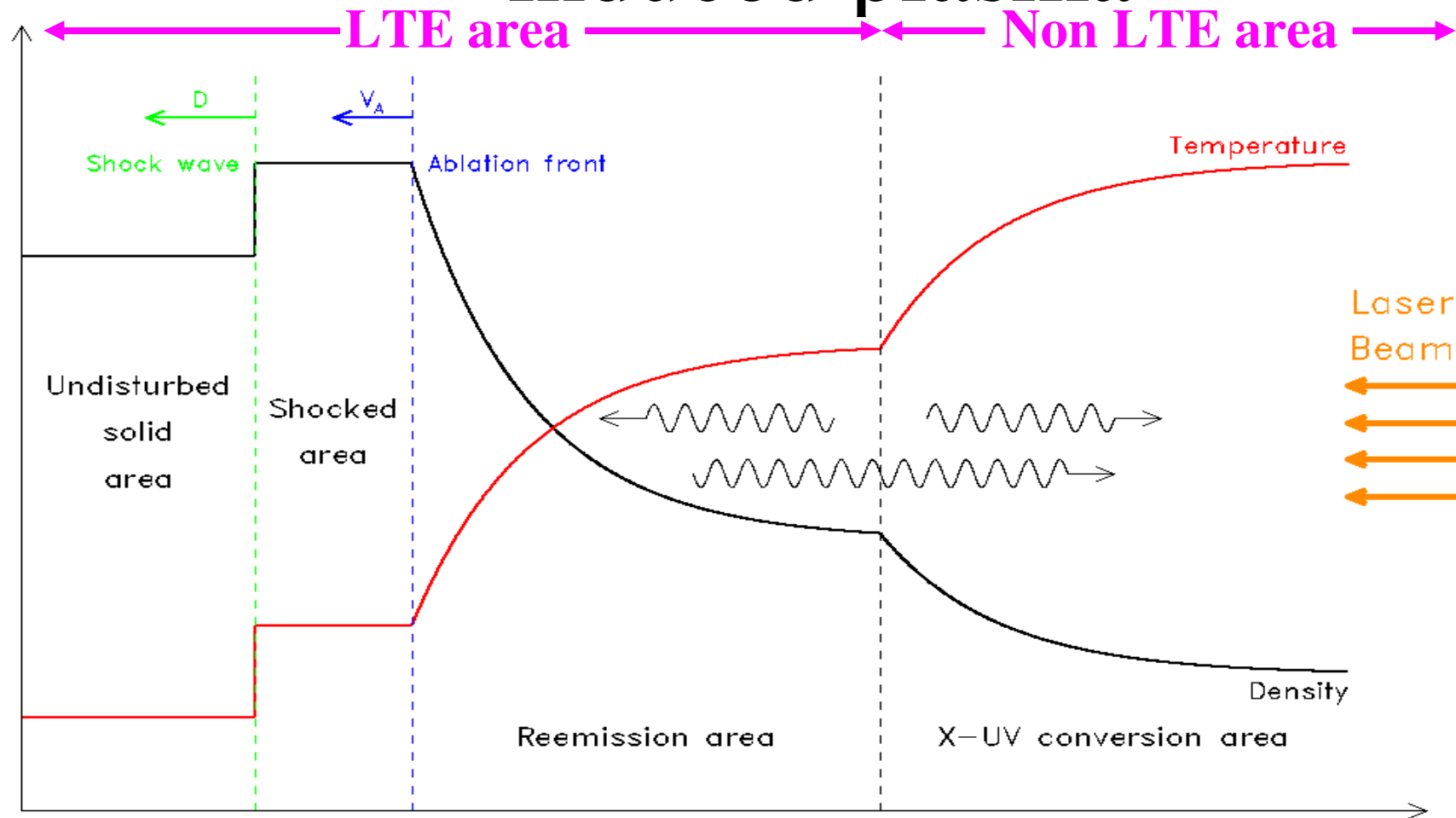
$$\lambda_{\text{NEET}}(\rho, T_e) = \frac{2\pi}{\hbar} D_1 p_1 (1 - p_2) \left| R_{12}(\bar{\delta}) \right|^2 \frac{e^{-\frac{\bar{\delta}^2}{2\sigma^2}}}{\sqrt{2\pi}\sigma}$$



P. Morel, V. Méot, G. Gosselin, D. Gogny, W. Younes, Phys. Rev. A60, 063414 (2004)

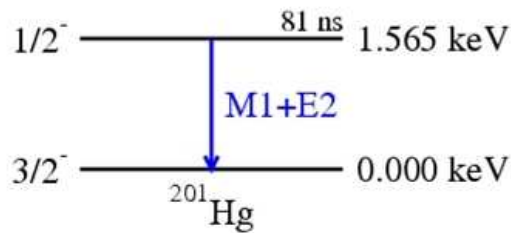


Thermodynamic profile of a laser induced plasma





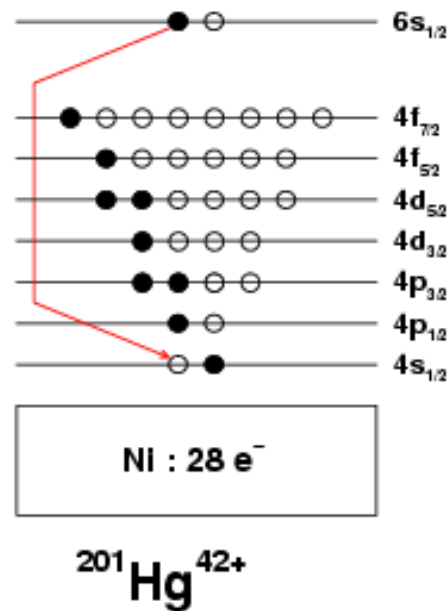
Non LTE NEET rate : ^{201}Hg



- Dirac-Fock model
 - Detailed configurations

3705 NEET-allowing distributions

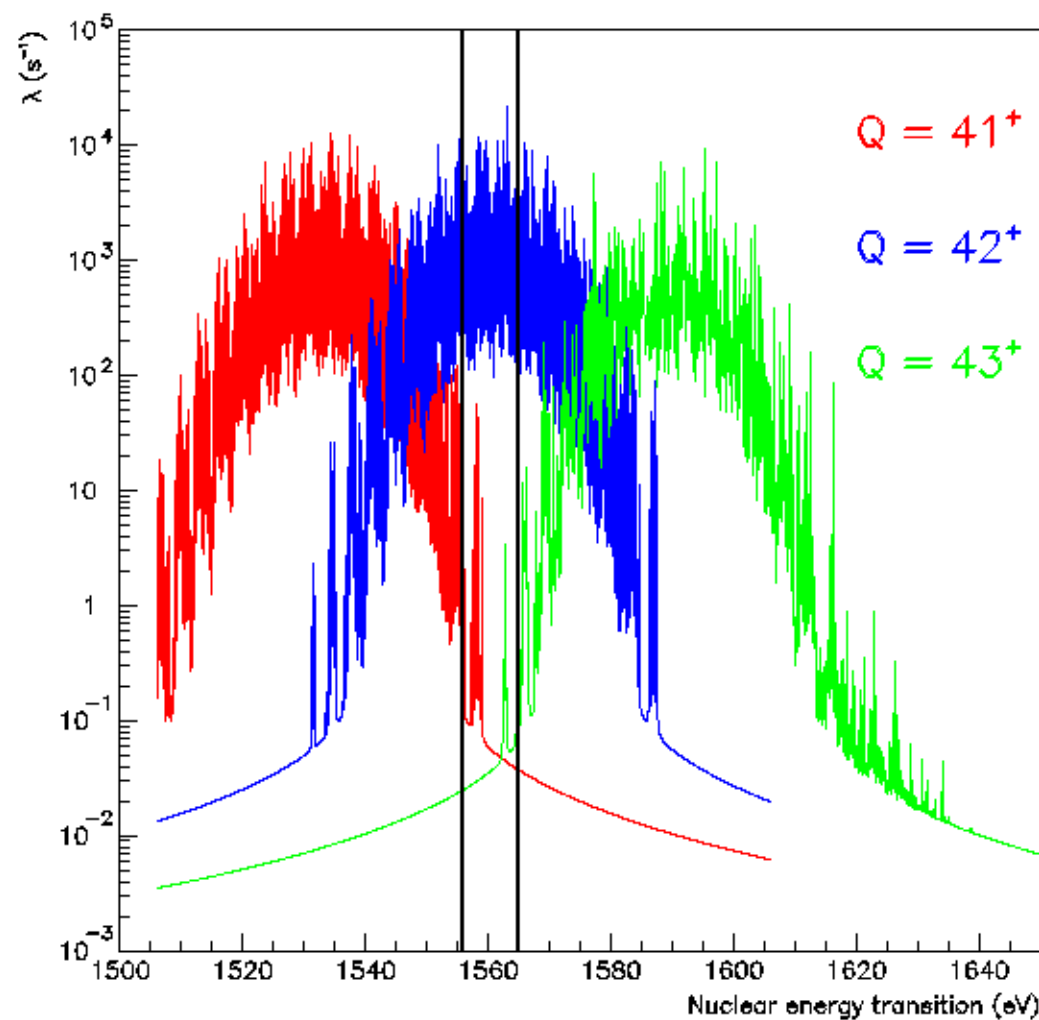
- For each configuration :
 - Total energy by DF calculations
 - Probability by Bernouilli distribution



$$P(N = k) = C_n^k p^k (1 - p)^{n-k}$$



^{201}Hg NEET rate for N1-P1 transition

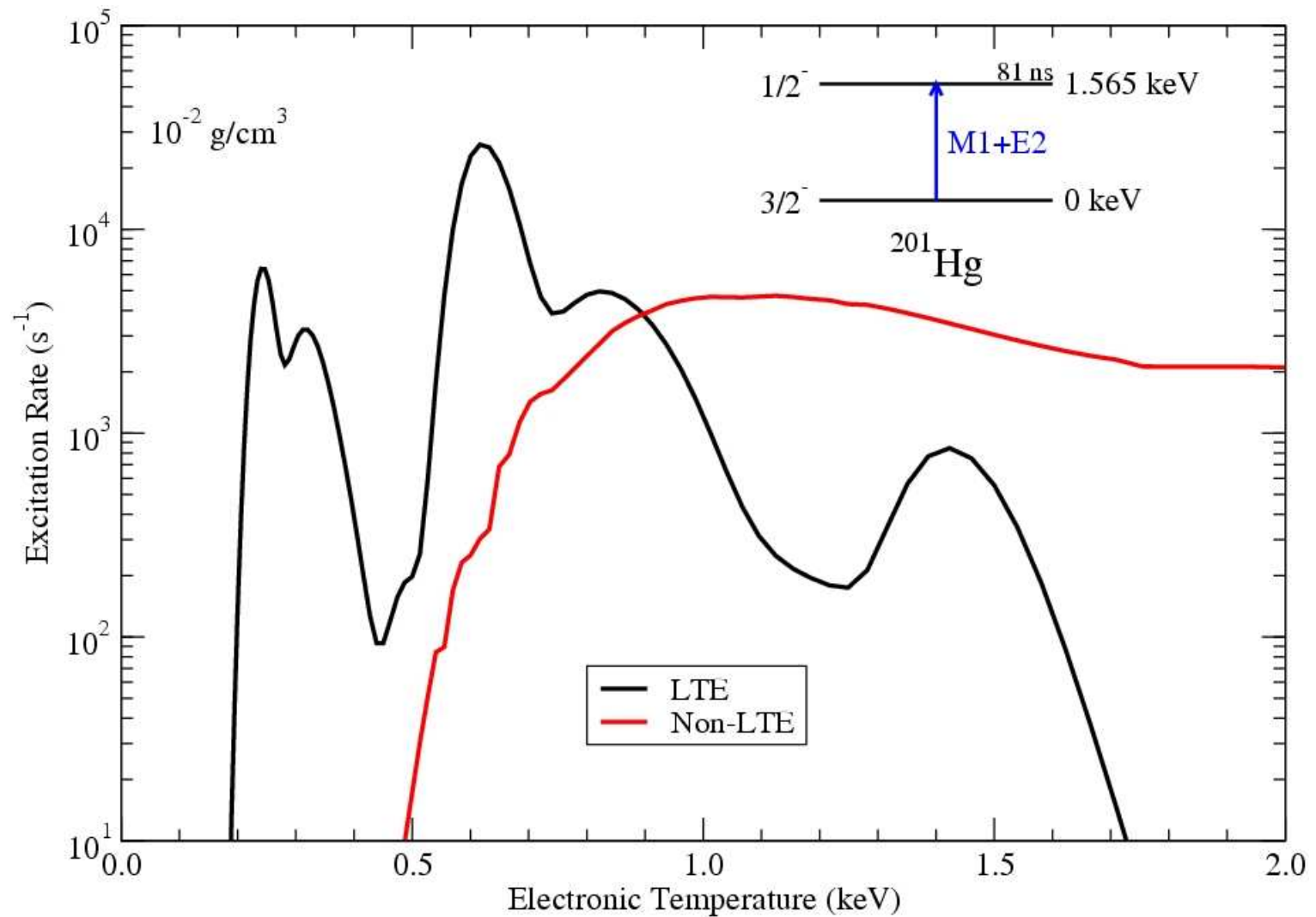




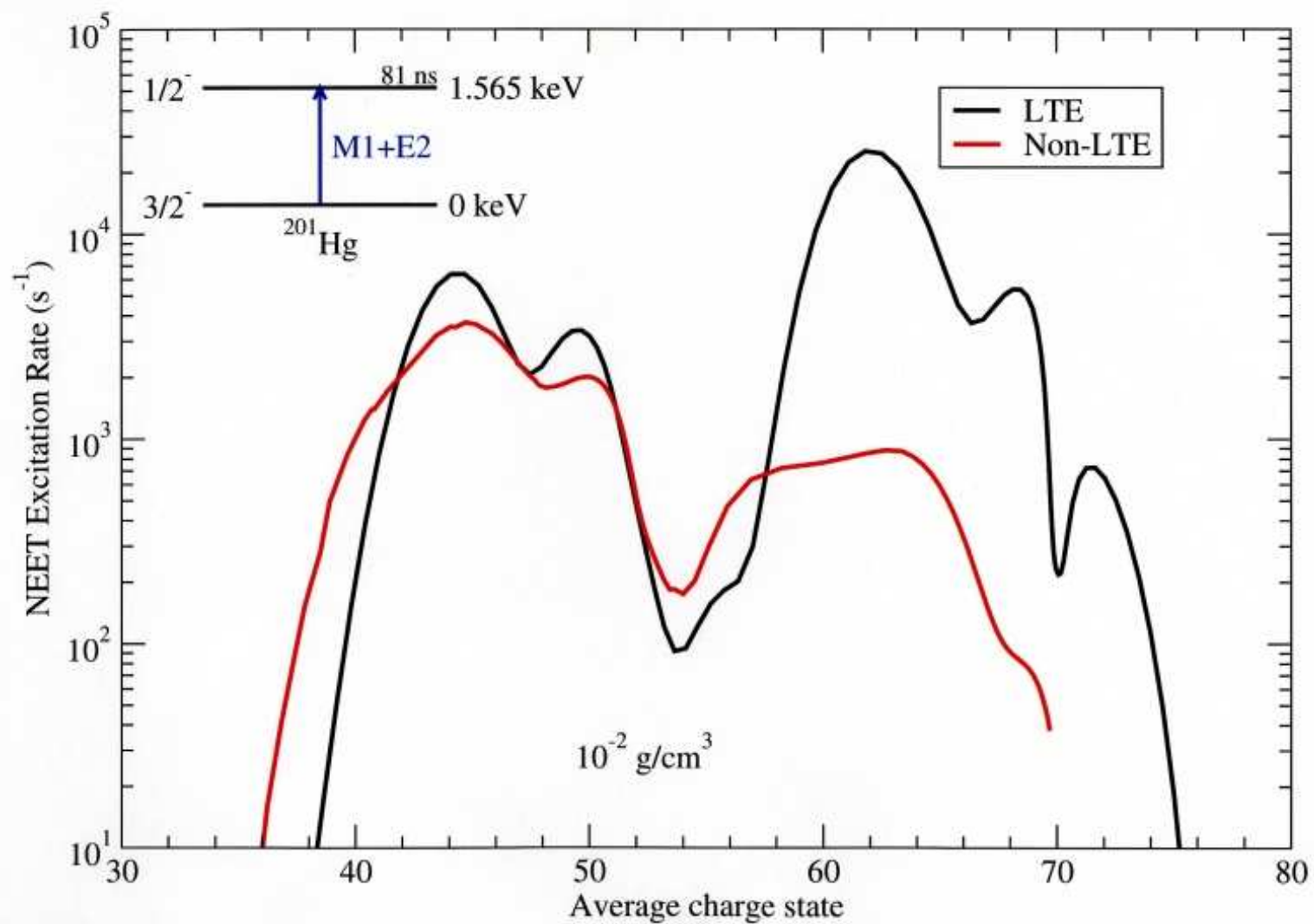
Self consistent average atom model SHAAM

$$\lambda^{\text{NEET}}(\rho, T_e) = \frac{2\pi}{\hbar} D_1 p_1 (1 - p_2) |R_{1,2}|^2 \frac{1}{\sqrt{2\pi\sigma^2}} e^{-\frac{\langle\delta\rangle^2}{2\sigma^2}}$$

- Screened Hydrogenic Average Atom Model
 - LTE / non-LTE atomic physics calculation (ρ , T_e , T_r)
 - Charge state distribution
 - Mismatch variance σ
 - Average electronic shell populations
- Relativistic Average Atom Model
 - Relativistic but only LTE
 - Wave functions
 - Binding energies



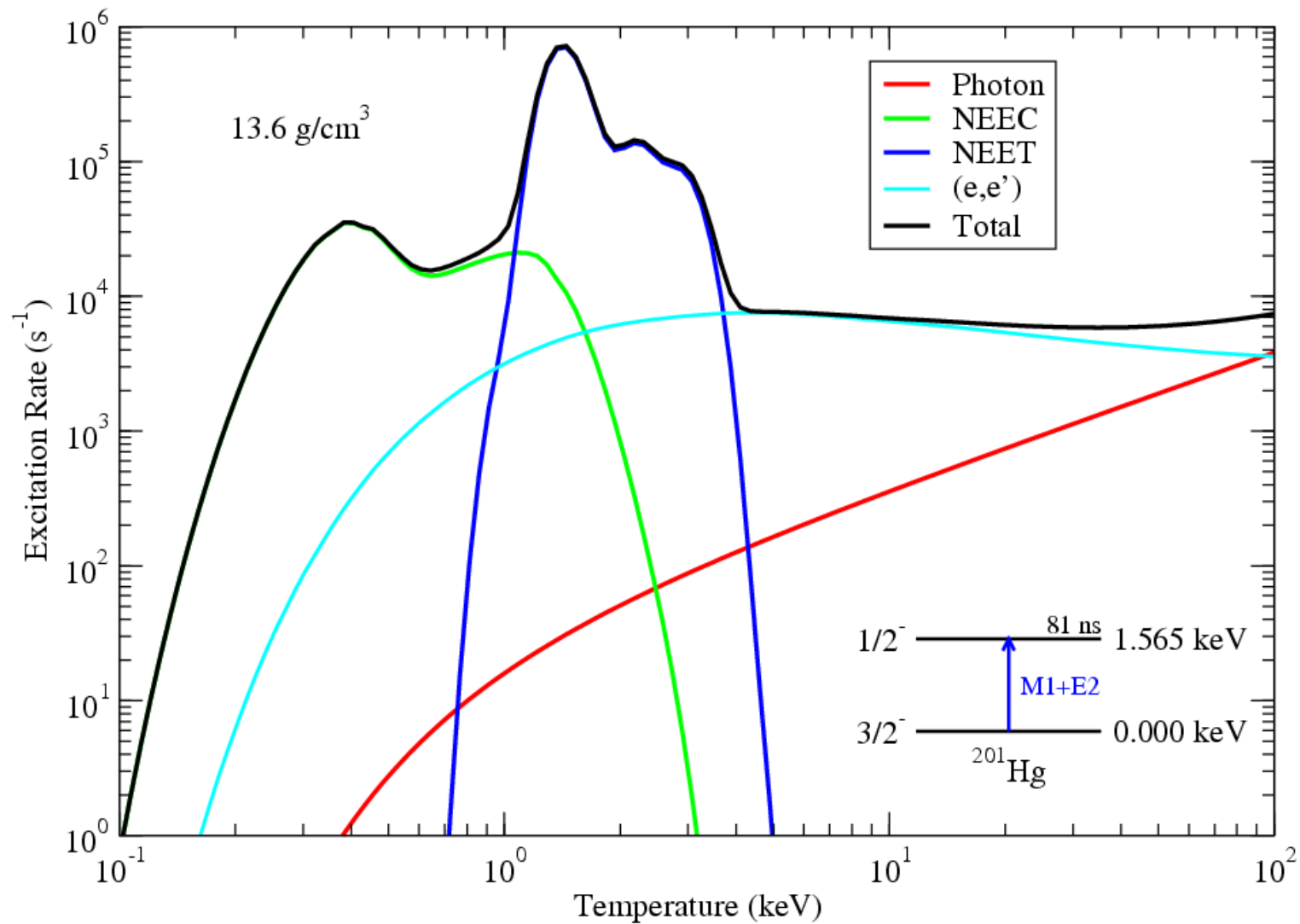
P. Morel, V. Méot, G. Gosselin, G. Faussurier, Ch. Blancard, Phys. Rev. **C81**, 034609 (2010)

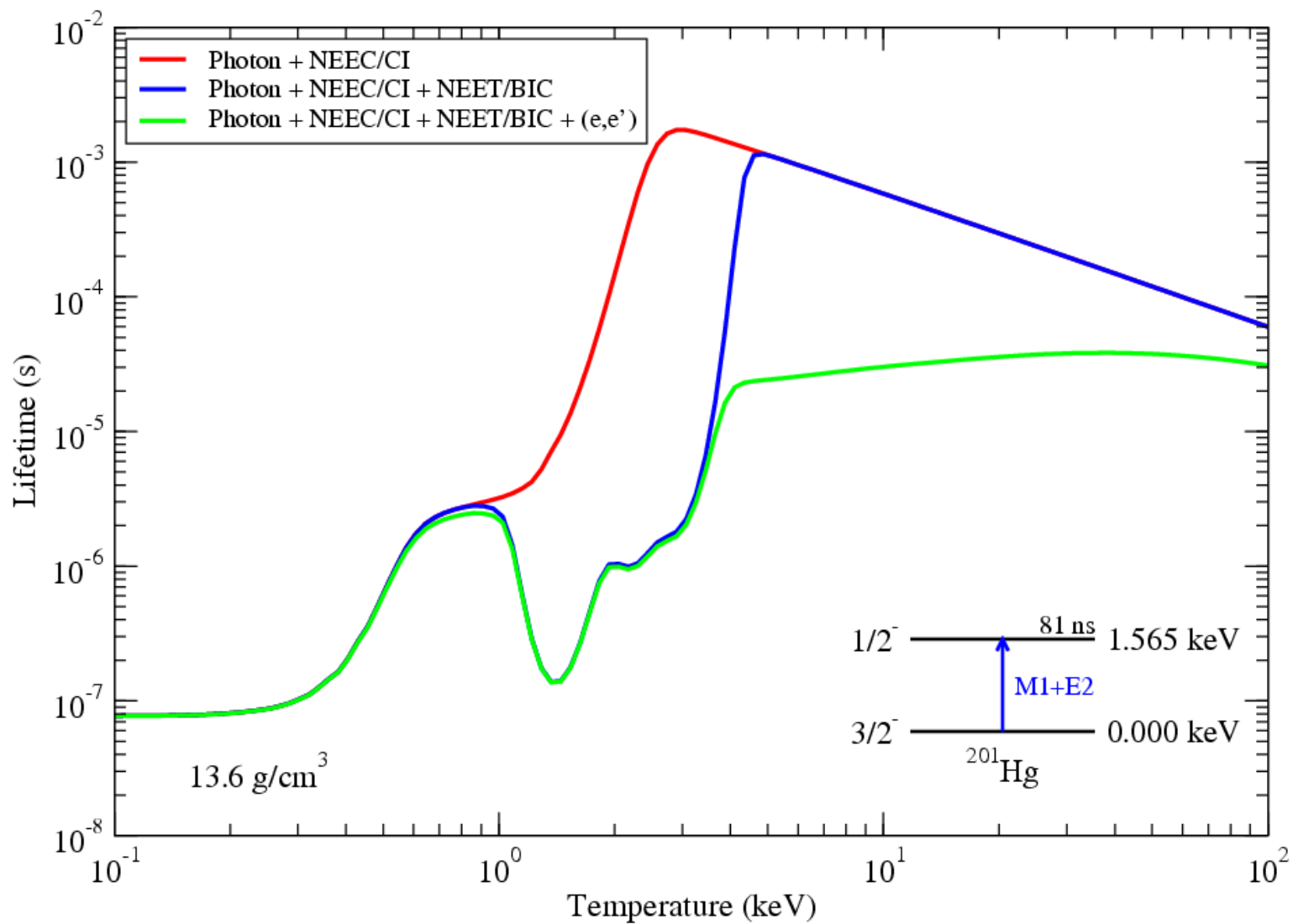




Total Excitation Rate

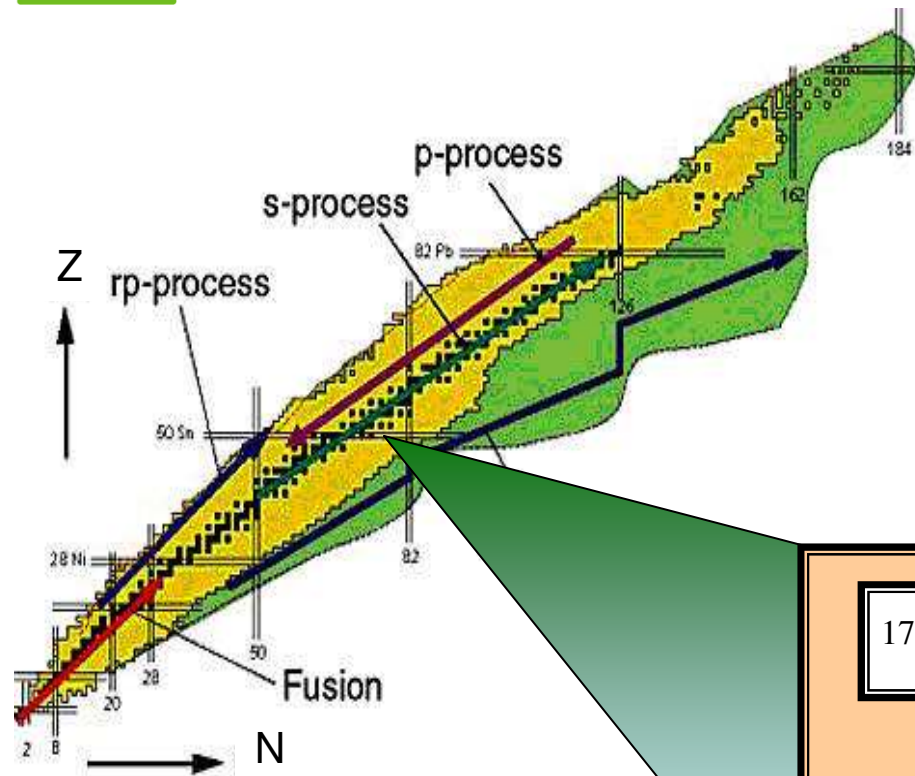
ETL Excitation Rates





G. Gosselin, V. Méot, P. Morel, Phys. Rev. C76, 044611 (2007)

The s-process (slow neutron capture) makes most nuclei with $A > 56$ via (n, γ) in massive stars



S-process conditions

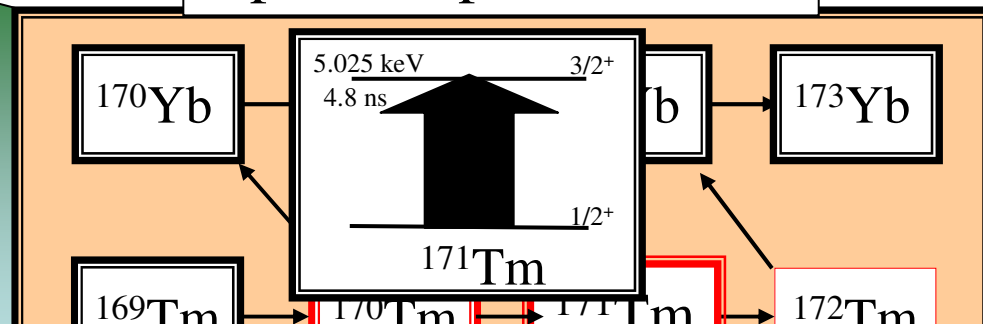
$$k_B T \approx 8, 30 \text{ keV}$$

$$\Phi_n \approx 10^7 / \text{cm}^2 \cdot \text{s}$$

$$\rho \approx 50\text{-}100 \text{ g/cm}^3$$

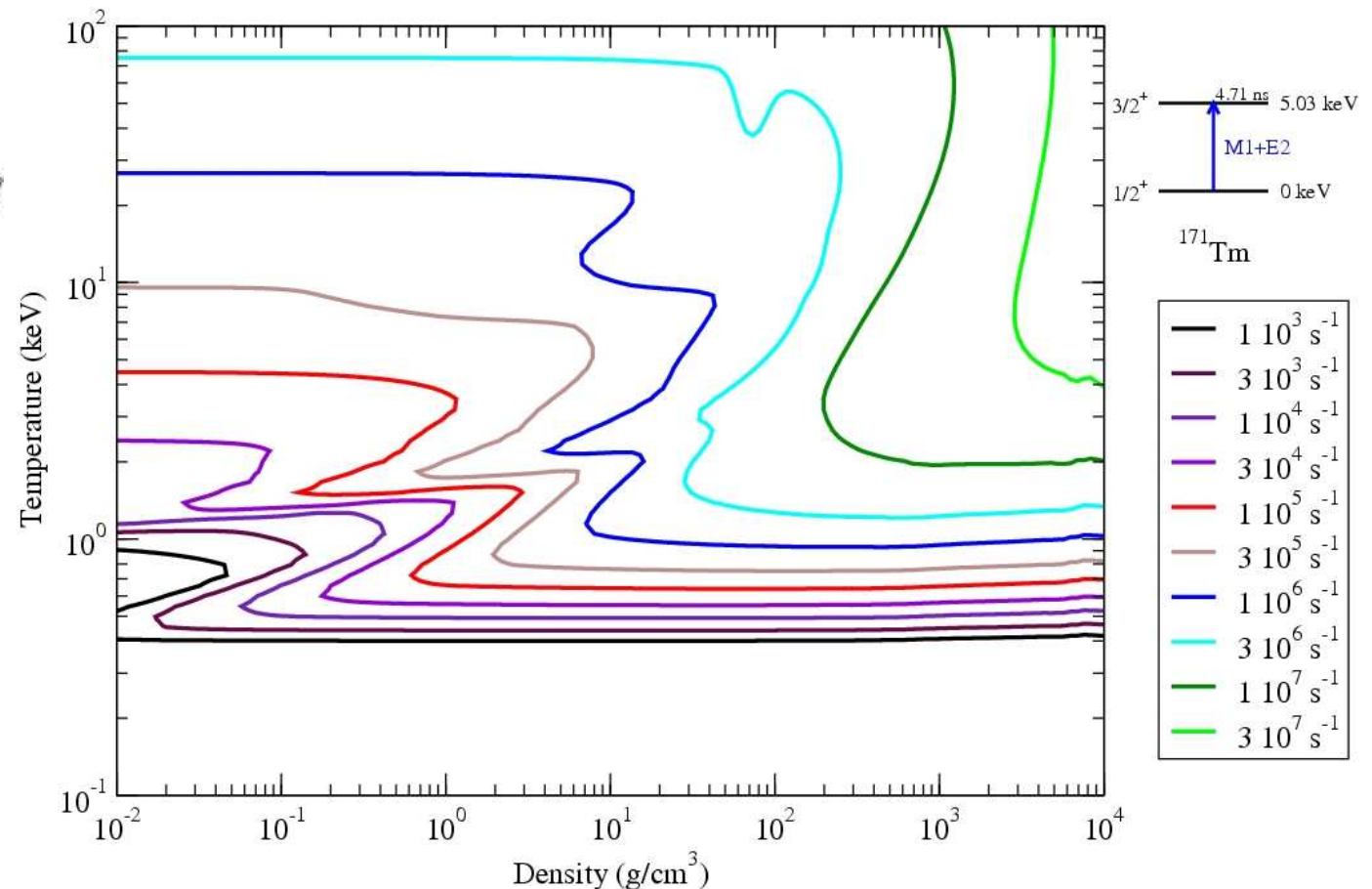
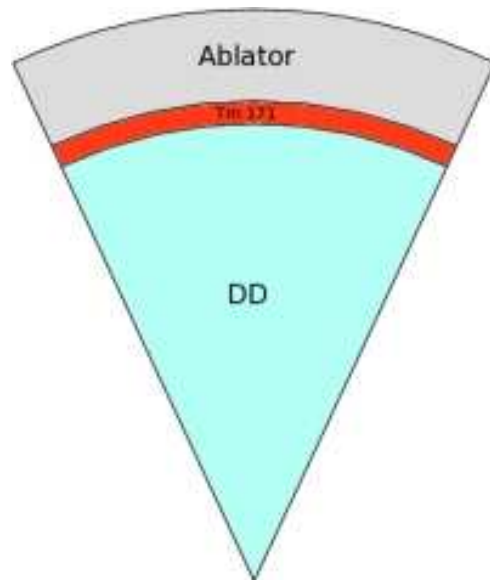
Lee Bernstein, LLNL

s-process path near Tm



S-process branching provides a sensitive measure of stellar conditions

NIF Experiment

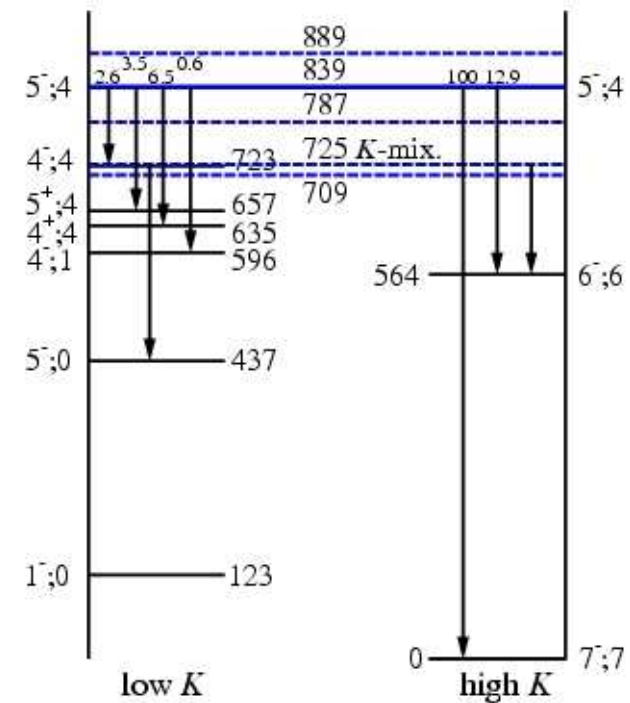
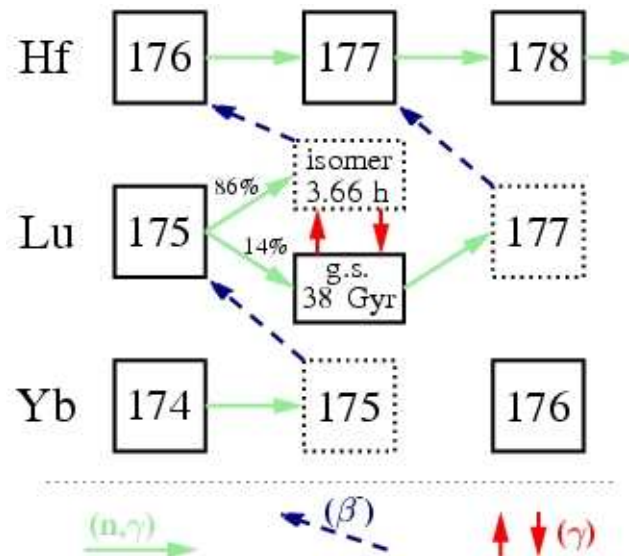


Requires hydrodynamic path optimization



Astrophysics Application

- S-process : ^{176}Lu
 - Low-K and High-K bands



In collaboration with P. Mohr, Diakonie-Klinikum, Schwabish-Hall



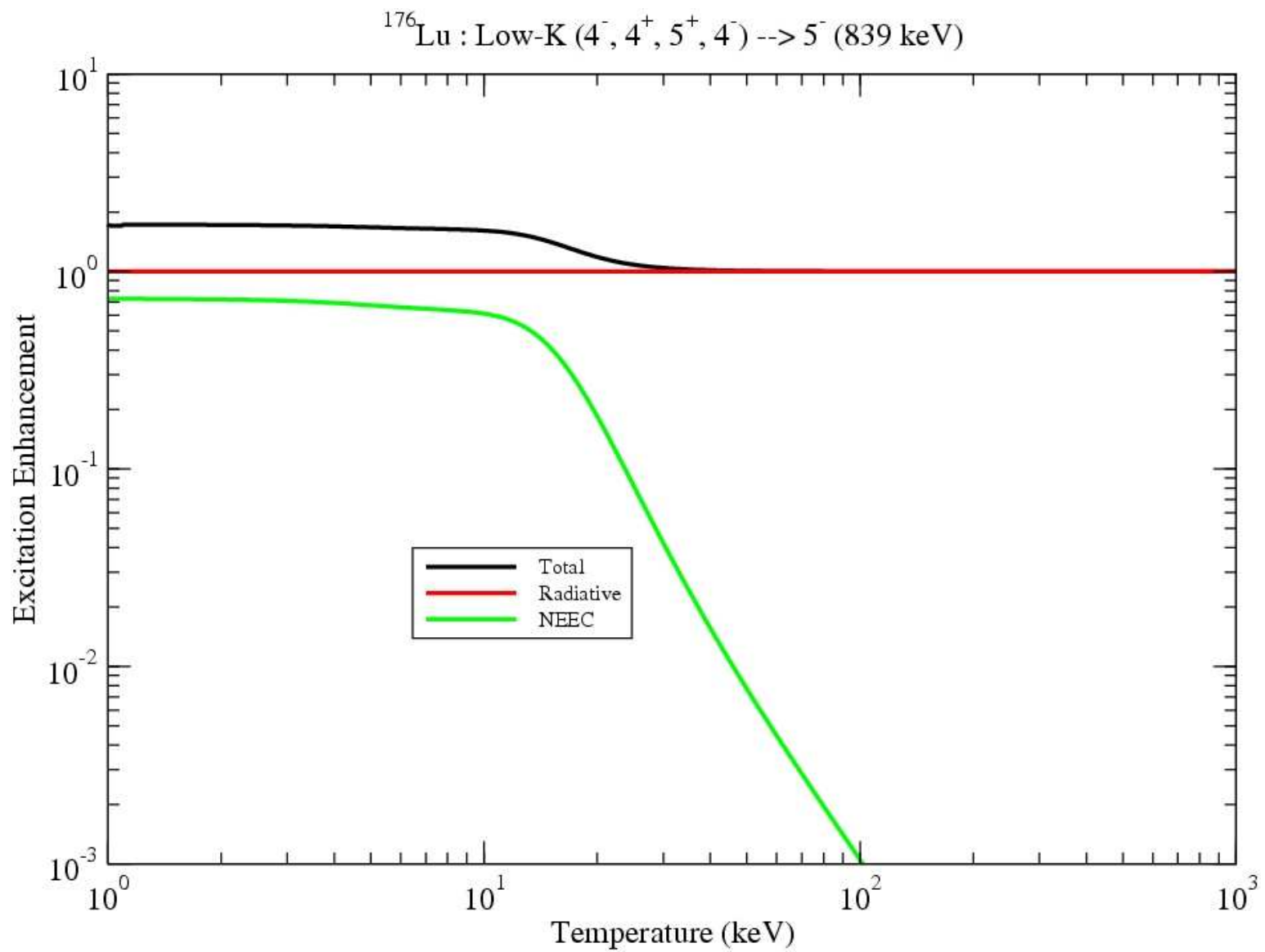
Stellar Transition Rate

- Indirect transition through intermediate state (IMS)

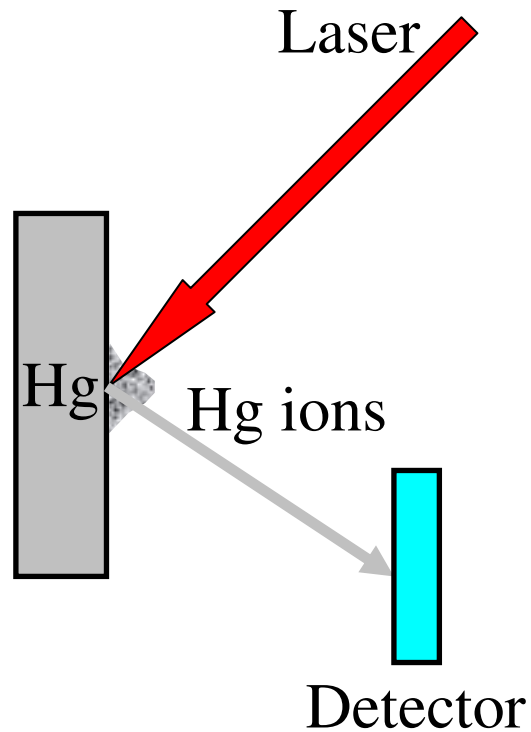
$$\lambda^*(T) = c \sum_i n_\gamma(E_{\text{IMS},i}, T) I_\sigma^*(E_{\text{IMS},i}, T)$$

$$I_\sigma^*(E_{\text{IMS}}, T) = \int \sigma(E) dE = \frac{2J_{\text{IMS}} + 1}{2J_0 + 1} \left(\frac{\pi \hbar c}{E_{\text{IMS}}} \right)^2 \frac{\Gamma_{\text{IMS} \rightarrow \text{lowK}}^* \Gamma_{\text{IMS} \rightarrow \text{highK}}^*}{\Gamma_{\text{IMS} \rightarrow \text{lowK}}^* + \Gamma_{\text{IMS} \rightarrow \text{highK}}^*}$$

- Width enhancement
 - Induced photon emission
 - NEEC/IC
 - Electron Inelastic Scattering

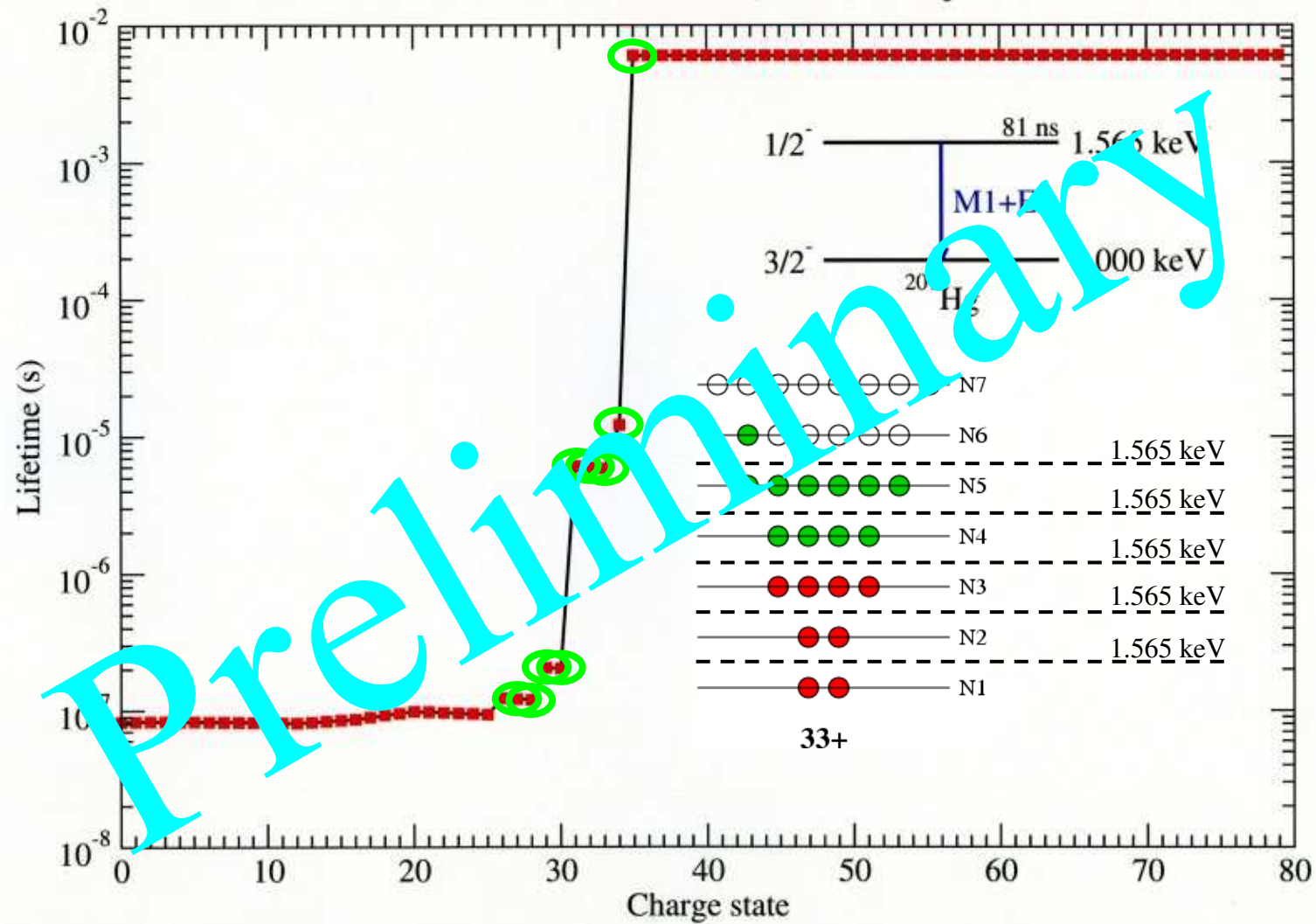


Laser Experiment



- Hg ions in flight
 - Highly ionized
 - Atom in fundamental state
 - Lifetime governed by IC
- How far can we detect them ?
 - MCDF calculations $\propto V(r)$
 - IC coefficients $\propto$ Lifetime

ISOMEX/MCDF Calculation (CEA Bruyeres-le-Chatel)





What we have done

- LTE modelization of excitation processes
 - Radiative
 - NEEC
 - NEET (when atomic transitions overlap)
 - (e,e') (unscreened potential)
- Non-LTE
 - NEET with SHAAM
- Experimental proofs
 - (p,n) reactions in plasma



What we are trying to do

- LTE modelization of excitation processes
 - NEET (when atomic transitions do not overlap)
 - Multi Atom plasmas
 - (e, e') (screened potential)
- Non-LTE
 - NEET with MCDF
- Future experiments
 - Isomer decay in plasma
 - Lifetime variations in plasma
 - Isotopic and/or isomeric ratios modification in Tm



Questions are
welcome

Thank you for
your attention